

INVESTIGATION OF FERRITE PHASE-SHIFTERS  
AT C- AND X-BAND

FINAL REPORT

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INVESTIGATION OF FERRITE PHASE SHIFTERS AT C- AND X-BAND

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July 1965

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This report is subdivided into three parts, each with its own page numbers, references, etc. These parts are:

- I. Phase Shifter Development  
by Robert Seckelmann
- II. Non-Linear Losses in Ferrites  
by Edward J. McKinney
- II. Propagation of TE Modes in Dielectrically Loaded Rectangular Waveguides  
by Robert Seckelmann

Part I. contains data on the experimental work and general acknowledgements.

Part II. and Part III. are theoretical studies.



## PART I. PHASE SHIFTER DEVELOPMENT

### ABSTRACT

The phase shifters discussed are nonreciprocal. They contain circumferentially magnetized ferrite cylinders, which work at remanence. The differential phase shift is controlled by partial magnetization of the ferrite, achieved by controlled flux transfer. Three problems, common to all such phase shifters are dealt with:

1. The efficiency of flux transfer and the possibility to add flux bits. Experimental data and a theory in agreement with these data are presented.

2. The determination of r-f magnetic threshold fields above which the losses of ferrites increase with field strength. The threshold is calculated from the measured threshold power of guides completely or in part filled with ferrite.

3. The influence of waveguide and ferrite dimensions on the frequency dependence of the differential phase shift, on the proportionality between phase shift and transferred flux, and on the peak power handling capability is investigated experimentally. Slabline and waveguide type devices have been used.

<u>Section</u>	<u>Page</u>
I. PROGRAM OUTLINE .....	1
II. CONTROL OF FLUX TRANSFER .....	4
III. NON-LINEAR LOSSES IN FERRITES .....	17
IV. RECIPROCAL TEM MODE PHASE SHIFTER.....	21
V. DIFFERENTIAL PHASE SHIFTERS .....	23
V.1 Waveguide Type Phase Shifters .....	25
V.2 Slabline Type Phase Shifters .....	39
V.3 General Remarks .....	50
VI. DRIVER CIRCUITS .....	51
VII. ACKNOWLEDGEMENTS .....	52
VIII. REFERENCES .....	54

## I. PROJECT OUTLINE

The object of the program reported here was to demonstrate the feasibility of non-reciprocal multibit latched ferrite phase shifters for frequencies in the 4 - 12 GHz range, capable of handling peak powers of 10 KW and controlled by flux transfer techniques.

Tentative goals were set in contract review conferences during the initial phase of the development as follows:

- Investigate slabline and waveguide type devices
- Switch in four bits ( $180^\circ$ ,  $90^\circ$ ,  $45^\circ$ ,  $22.5^\circ$ )
- Achieve at least  $360^\circ/\text{dB}$
- Drive with 500  $\mu\text{joules}$  or less per switching cycle
- Allow for a switching cycle approximately 5  $\mu\text{sec}$ .
- Understand the trade-offs between the various parameters and formulate design rules

The phase shifter requirements could have been fulfilled by using four ferrite bits of different lengths along the axis of the device and switching them between their largest negative remanent magnetization and some positive remanent magnetization smaller than the largest one by means of controlled flux transfer to secure a phase shift independent of temperature.<sup>(1)</sup> In review conferences, however, the investigators were encouraged to pursue their idea to achieve several bits of differential phase shift by using only one bit of ferrite along the axis of the device and to control the phase shift by controlling the state of magnetization of this ferrite. Consequently a device very close to an analogue phase shifter has been developed. This second approach requires a much tighter control of the field distribution in the device and imposes more restrictions on the ferrite and guide shape than the first approach would have done.

Flux transfer circuits have been analyzed theoretically and the theory has been found to be in good agreement with experimental data. Circuits have been designed to add phaseshift bits by adding flux bits.

The phase shifters and the systems using them could be considerably simpler, if it were possible to eliminate some of the switching wires. This should be possible by supplying to the phase shifter voltage pulses of constant height (voltage) and variable length (time). Ideas for such circuits - temporarily called "analogue drivers" - have been conceived during the last phase of the contract, but have not yet been tested. An electronic driver of very small volume and steered by a computer could be incorporated in the phase shifter. Whereas now such a computer has to select - in the case of a four bit phase shifter - any of 16 bit combinations, it would then select any of 16 times, during which it would activate the driver. All phase shifters studied use round ferrite cylinders. The devices have too complex a geometry to be directly dealt with mathematically. Rectangular waveguides with dielectric slabs parallel to the narrow guide walls serve as more or less crude analytical models. The propagation of TE modes in these loaded guides has been studied in detail. Numerical results for a large range of parameters are given in Part III of the report. They will also be published as a G.E. TIS Report. The computer program to obtain these data was financed in part by this contract. The slabline type phase shifter performance could be explained nicely in terms of TE modes in dielectrically loaded rectangular waveguides. The computer results were also used to evaluate tests to determine the threshold field strength for the onset of non-linear losses in a rectangular waveguide containing a ferrite slab.

The differential phase shift versus flux dependence led to a model for magnetization of ferrite toroids, which in turn led to a reciprocal TEM phase shifter controlled by partial magnetization. This reciprocal phase shifter has been tested in one of the high power test structures.

The development of high-power ferrites is covered by other contracts. Here the influence of the device geometry on the on-set of non-linear losses was considered. The requirements were met, and results encouraging the development of phase shifters capable to handle higher peak powers, say 50-100 KW were obtained. An effort has been made to understand and predict the onset of non-linear losses in ferrites of arbitrary shape. A theoretical study is

given as part II of this report. This treatise includes the results obtained by other workers with more restricted models. At the moment, however, it cannot be used to its full extent because not all parameters could be determined in experiment.

Figures of merit of more than  $400^\circ/\text{dB}$  have been obtained throughout the development. Nearly octave bandwidths with only a few percent variation in differential phase shift have been obtained in waveguide type phase shifters. The highest peak power handling capability, the broadest bandwidth and proportionally between differential phase shift and flux transferred have, however, not been obtained simultaneously. Trends have been found for the optimization of each characteristic, but the trade-offs are not fully understood. Nearly all parameters are mutually dependent (e.g. ferrite wall thickness, diameter, waveguide height and width, filling factor).

The project was started with three parallel efforts:

- to investigate flux transfer circuits
- to study the onset of non-linear losses in ferrites
- to develop phase shifters at low powers.

In the second half of the project the flux transfer investigation was replaced by an effort

- to develop driver circuits

Finally, the phase shifters were tested with high peak power pulses. The different efforts will now be described in the order given above.

## II. CONTROL OF FLUX-TRANSFER

A flux transfer circuit is shown in Figure 1. R is a variable resistor in the circuit.  $C_D = C_1, \dots, C_n$  are n "driver" cores (index D) of various sizes, which are switched from negative to positive saturation magnetization or the corresponding remanent magnetization. C is the "microwave" core (no index), which produces the differential phase shift. For both, the driver and the microwave core the simple BH-loops of Figure 1 are used to analyze the transfer circuit. B is the flux density and H is the magnetic field strength.

In an unmagnetized ferrite core the magnetic dipoles are randomly oriented. There exist domains for any orientation. In a core saturated by sending a strong current pulse through the core the ferrite is circumferentially magnetized. Sending now an increasing current in opposite direction through the core, one reverses gradually the sense of circumferential magnetization. Two models to describe the magnetizing process may be used. In a "many shells" model the ferrite may be looked upon as consisting of many successive, very thin, laminar layers. A magnetizing current will either not change the state of magnetization of a particular shell or drive it to saturation magnetization. According to this model, an increasing d.c. current would at first "switch" the inner ferrite rings, then the outer ones. An a.c. current with decreasing amplitudes would produce a sequence of circumferentially magnetized ferrite shells with alternating directions of magnetization. In a "one shell" model the magnetization would be homogeneous throughout the ferrite wall and can vary from complete spin alignment in one direction to random spin orientation to complete alignment in the other circumferential direction, depending on the driving conditions. An a.c. current pulse with decreasing amplitudes sent through a "switching" wire through the axis of the ferrite core would demagnetize the ferrite, i.e. randomly orient the spins. It is felt that often a super position of a "many shells" and a "one shell" model will offer the best explanation of effects. Assuming for the ferrite a simplified BH loop as shown in Figure 1, the shell thickness would be given approximately by  $r_1'/r_2' = H_s/H_c$  where  $r_1', r_2'$  are the radii of the layer switched (with  $r_1 \leq r_1' \leq r_2' \leq r_2$  and  $r_1, r_2$  the inner and outer core radii).

Dynamic permeabilities  $\mu = dB/dH$  are defined by

$$\mu^{(V)} = B_s/(H_s - H_c), \mu_D^{(V)} = B_{Ds}/(H_{Ds} - H_{Dc}) \quad (1a)$$

$$\mu^{(H)} = (B_s - B_r)/H_s, \mu_D^{(H)} = (B_{Ds} - B_{Dr})/H_{Ds} \quad (1b)$$

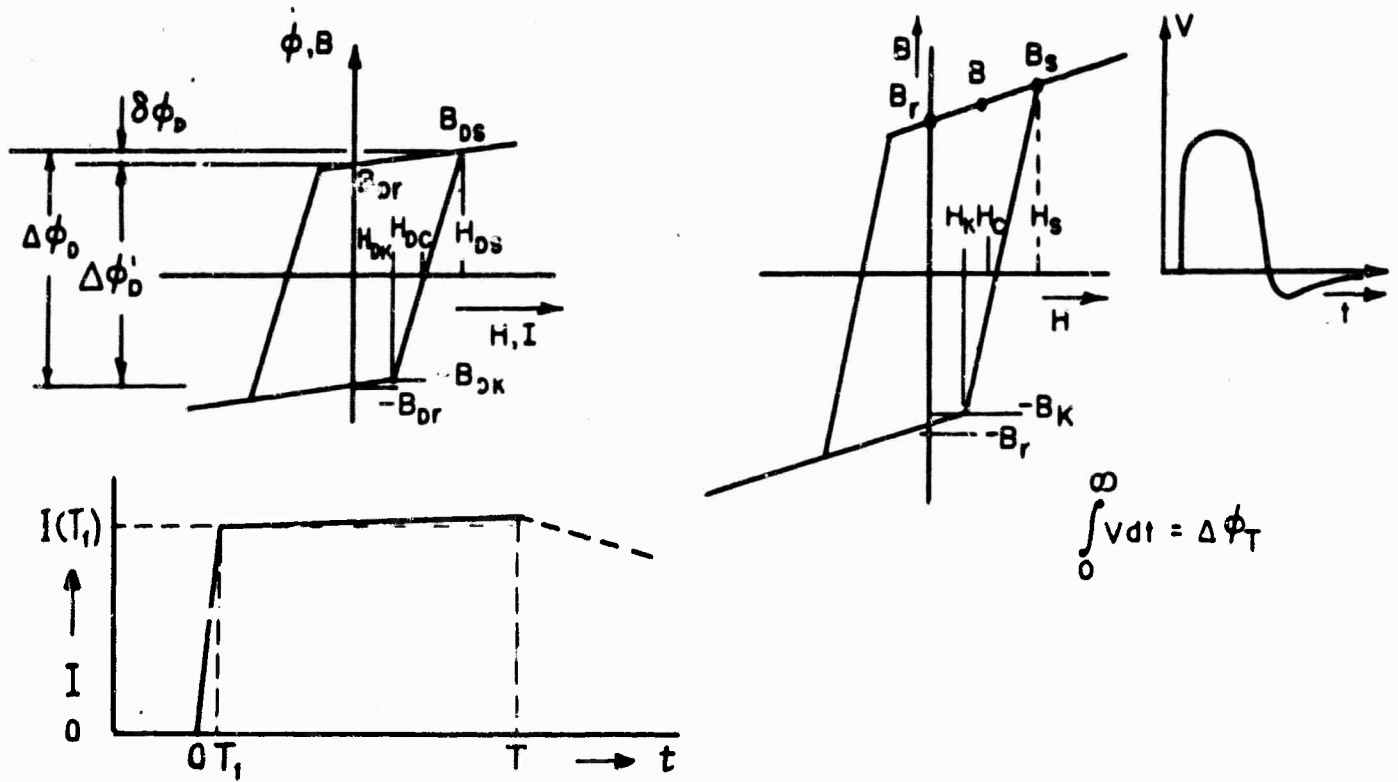
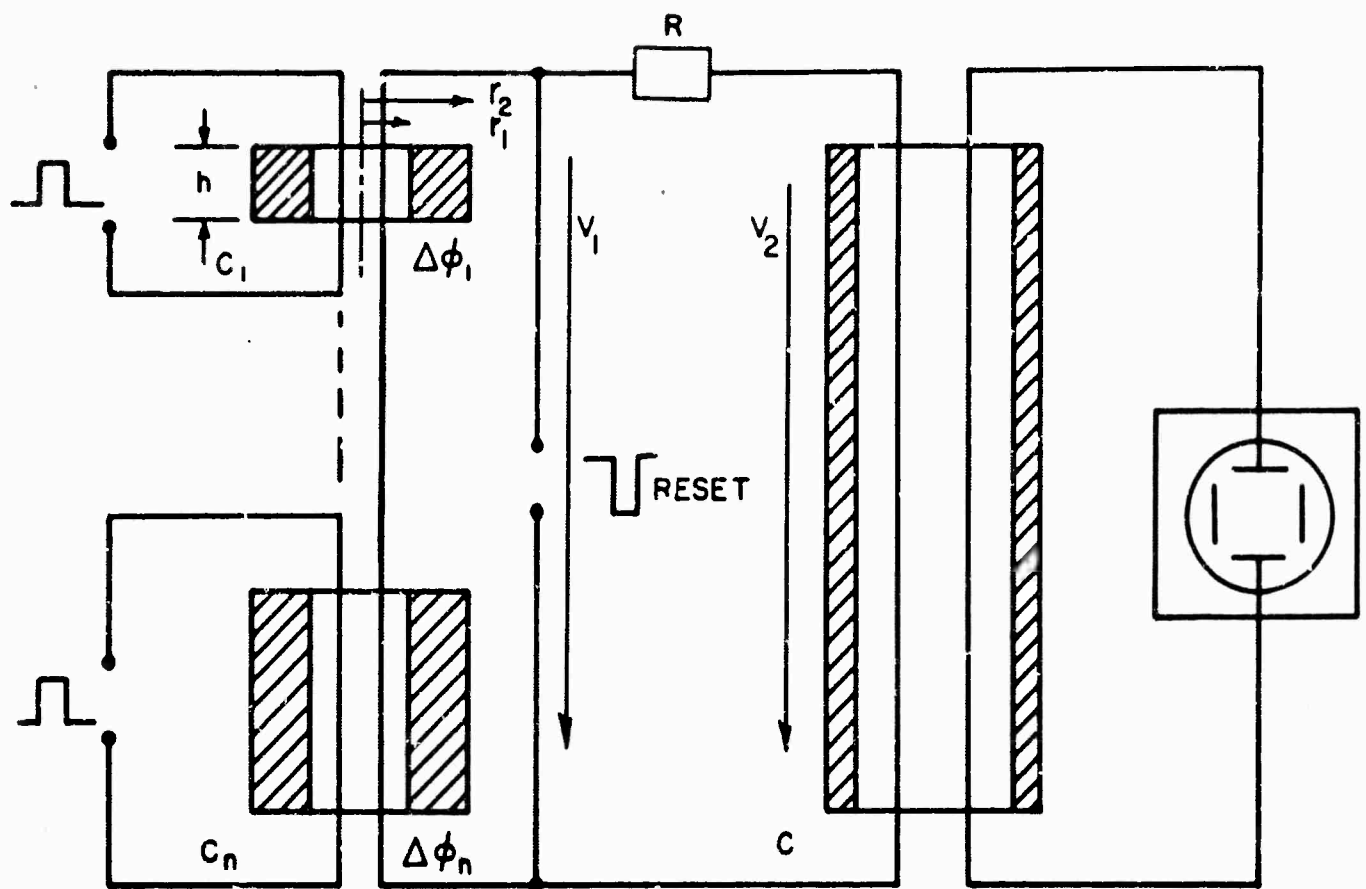
and dynamic inductances with the core height h by

$$L = d\phi/dI = \mu h \ln(r_2'/r_1')/2\pi \quad (2)$$

$A' = h(r_2' - r_1')$  is the area of a particular core, in which the flux it changed by

$$\Delta\phi = \Delta B \cdot A', \text{ with } A'_{\max} = A.$$





FLUX TRANSFER SWITCHING TECHNIQUE

FIGURE 1.

Before a switching cycle begins, a reset pulse drives both the driver and the microwave core to saturation and subsequently to the corresponding remanent state. Let this be the negative state for both cores,  $-B_r$  and  $-B_{Dr}$ , respectively. The driver core is now driven to  $B_{Ds}$  and then falls back to  $B_{Dr} = B_{Ds} - \delta B_D$ , the remanent state. It feeds the transfer circuit with an approximately rectangular voltage pulse of height (voltage)  $V_1$  and length (time)  $T$ , followed by a pulse with a relatively small negative peak voltage.  $T$  is proportional to the length of the driver core. The available flux is

$$\Delta\phi'_D = 2B_{Br} A = \Delta\phi_D - \delta\phi_D, \quad \Delta\phi_D = \int_0^T V_1(t) dt \approx V_1 T \quad (3)$$

Let  $\delta\phi_D \ll \Delta\phi_D$ , (with holding current applied or  $\mu_D^{(H)} \approx 0$ ), so that

$$\Delta\phi'_D \approx \Delta\phi_D \quad (4)$$

The flux transferred is

$$\Delta\phi_t = \int_0^\infty V_2(t) dt = \int_0^T V_2(t) dt - \delta\phi \quad (5)$$

$\delta\phi$  is a relatively small contribution caused by the falling back of the microwave core from a value  $B_1$  to  $B_1 - \delta B_1$ . With  $\mu_D^{(H)} \approx 0$ , the time constant for this fall back process is given by  $L^{(H)}/R$ . Let the microwave core be driven to a value  $B < B_s$ . The transferred flux is for a lossy circuit:

$$\Delta\phi_t \approx \int_0^\infty V_1(t) dt - \int_0^\infty i(t) R dt \quad (6)$$

$$\Delta\phi_t + \delta\phi = \int_0^T V_1(t) dt - \int_0^T i(t) R dt = \int_0^T L(V) \frac{di(t)}{dt} dt \quad (7)$$

so that

$$i(t) = (1 - \exp(-tR/L(t))) V_1/R \approx tV_1/L(t) \text{ with } t \leq T \ll L(t)/R$$

Initially, during  $0 \leq t \leq T_1 \ll T$ , the inductivity of the microwave core is  $L^{(H)}$ . The current in the transfer circuit rises fast as

$$I(t) = (1 - \exp(-tR/L^{(H)})) V_1/R \approx tV_1/L^{(H)} \text{ with } t \leq T_1 \ll T \quad (8a)$$

At  $t = T_1$  the inductivity becomes  $L^{(V)} \gg L^{(H)}$ , and the transfer current increases only slowly as

$$\begin{aligned} I(t) &= I(T_1) + (1 - \exp(-(t - T_1)R/L^{(V)})) (V_1 - I(T_1)R)/R \\ &\approx I(T_1) + (t - T_1)(V_1 - I(T_1)R)/L^{(V)} \approx I(T_1) \end{aligned} \quad (8b)$$

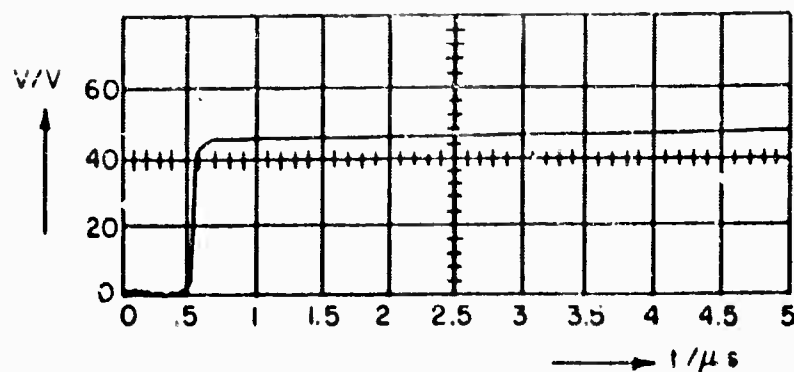
$$\text{with } T_1 \leq t \leq T \ll L^{(V)}/R$$

A HP 214A pulse generator was used to saturate the driver core. Figure 2 shows the output voltage of the unloaded source, the output voltage of the driver-core-loaded source, the output current of the driver-core-loaded source, the voltage across the secondary side of the driver core (open transfer circuit), the voltage across the secondary side of the driver core (with microwave core applied). One recognizes the fast increase in current (from the source) due to the small inductivity between  $B_{Dr}$  and  $B_{Dk}$ , then the slow increase of current due to the large inductivity between  $B_{Dr}$  and  $B_{Ds}$  and the fast increase between  $B_{Ds}$  and  $B > B_{Ds}$ . Figure 3 shows the current in the transfer circuit ( $R = 0$ ) and the voltage across the microwave core. The decaying current trail at the end of the current pulse produces a negative voltage across the microwave core, as indicated in Figure 1. In Figure 3 it cannot be seen. It follows the positive voltage pulse and has a very low peak value. Figure 4 shows the voltage  $V_1$  across the driver core for several driver core sizes (with holding current). The driver core supplies an approximately rectangular voltage pulse of height (voltage)  $V_1$  and a width (time interval)  $T$ . In Figure 4  $V_1$  is approximately constant and  $T$  proportional to the size of the driver core. It is not necessary to produce the time voltage area  $V_1 T$  via a driver ferrite. It can also be done electronically. As indicated in Figure 1, the flux loss via ohmic losses during the current rise time  $T_1$  is very small compared with the loss during  $T_1 \leq t \leq T$ . With this approximation and  $T - T_1 \approx T$ , the transfer efficiency becomes

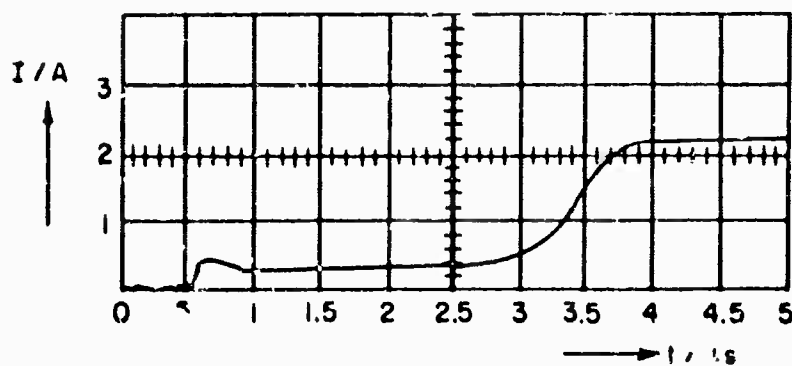
$$\eta = (\Delta\phi_t / \Delta\phi_D) \approx 1 - (\delta\phi / \Delta\phi_D) - (RT_1 / L^{(V)}) \quad (9a)$$

As long as the microwave core is not saturated  $L^{(V)} \sim \Delta\phi_D$ . In the "many shells" model this means that only a certain fraction of the total number of laminar layers are switched; in the "one shell" model  $L^{(V)}$  is a function of an effective permeability, which increases with driving current.

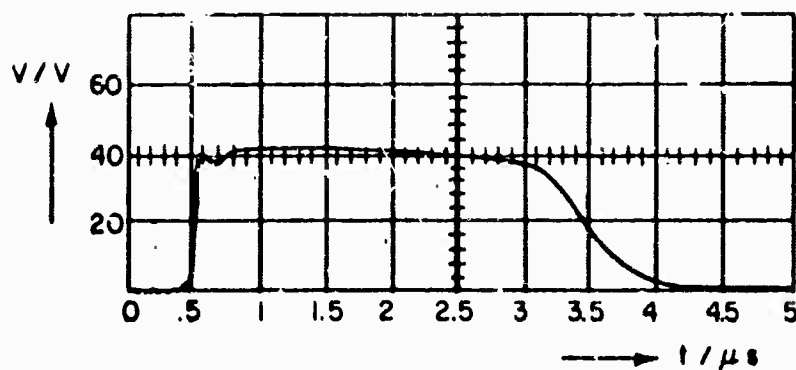
In Figure 5 the driver and microwave core sizes have been varied by using various numbers ( $n_D$  and  $n_M$ ) of equal ferrite rings for both, the



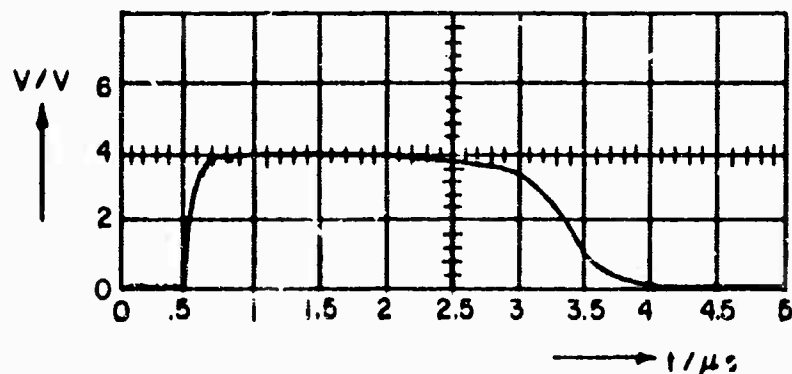
VOLTAGE OF UNLOADED  
PULSE GENERATOR HP-214A



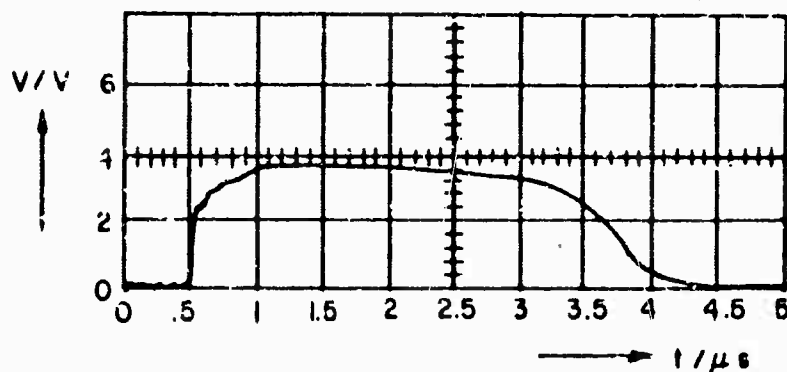
CURRENT FROM HP-214A  
TO UNLOADED DRIVER CORE



VOLTAGE ACROSS  
UNLOADED DRIVER CORE  
(PRIMARY SIDE)

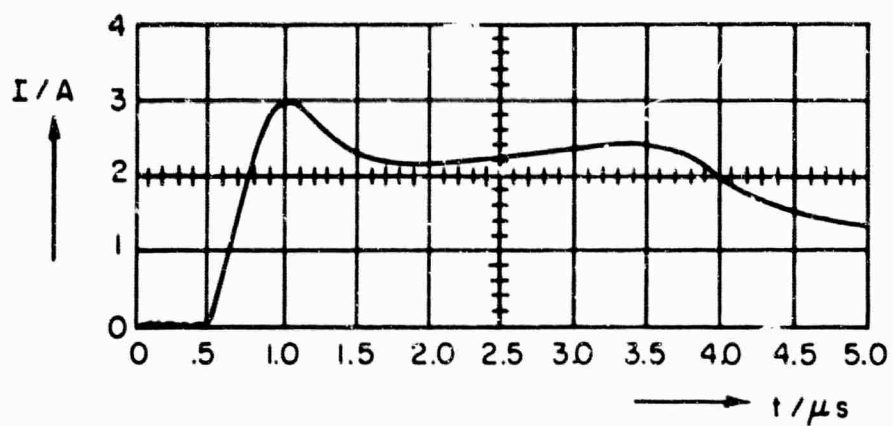


VOLTAGE ACROSS  
UNLOADED DRIVER CORE  
(SECONDARY SIDE)

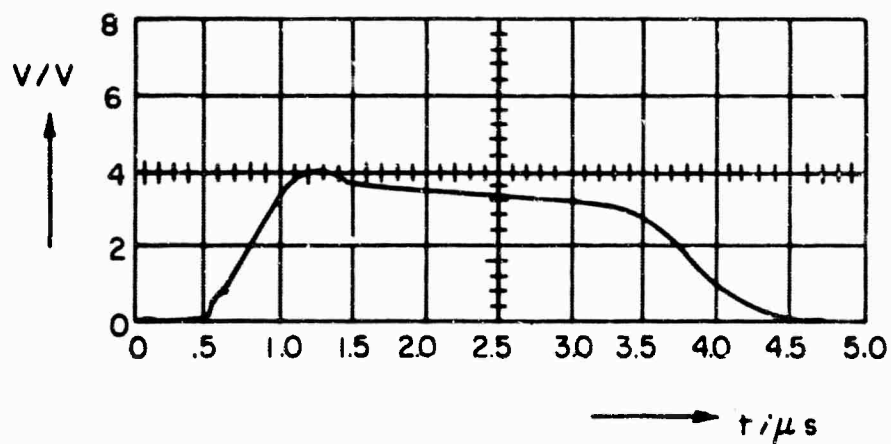


VOLTAGE ACROSS  
LOADED DRIVER CORE  
(SECONDARY SIDE)

FIG. 2 VOLTAGE AND CURRENT OF DRIVER CORE



CURRENT IN  
TRANSFER CIRCUIT

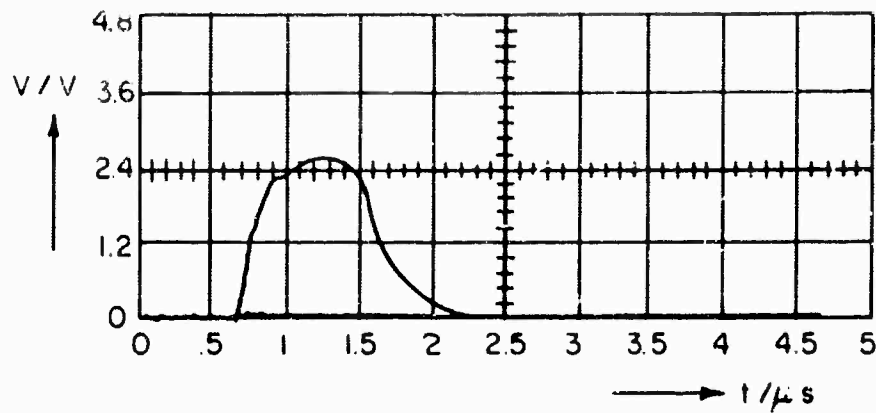


VOLTAGE ACROSS  
MICROWAVE CORE

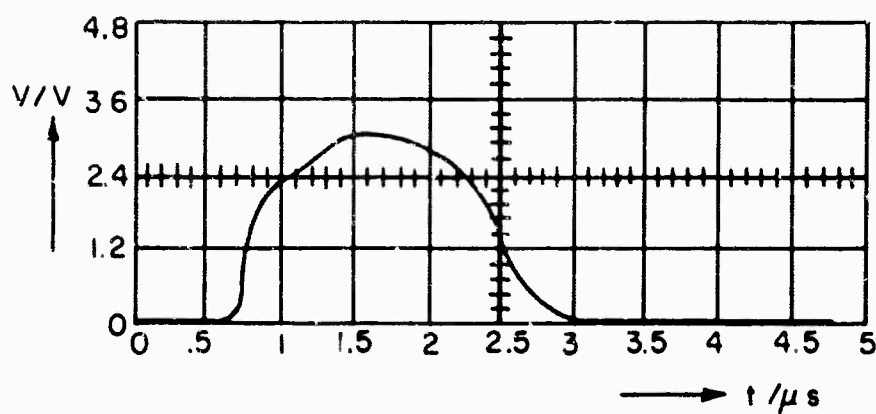
VOLTAGE AND CURRENT OF MICROWAVE CORE

FIGURE 3

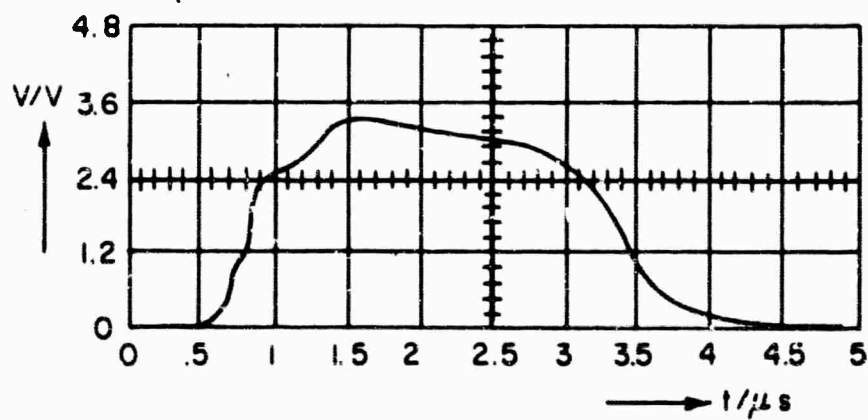




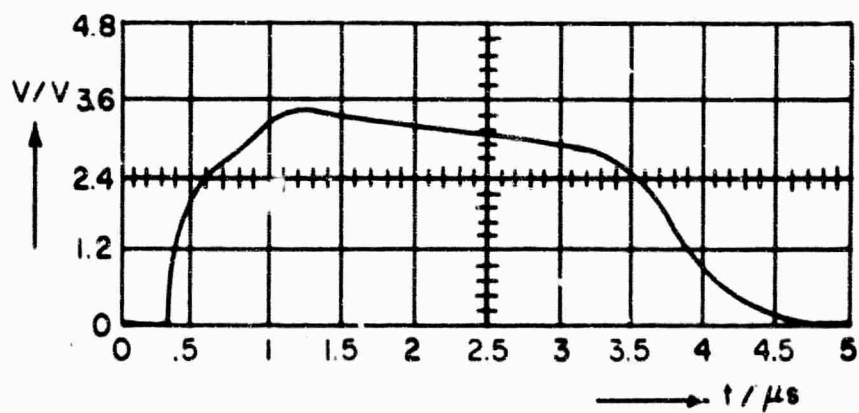
$$\int V dt = 1.96 \mu Vs$$



$$\int V dt = 4.2 \mu Vs$$



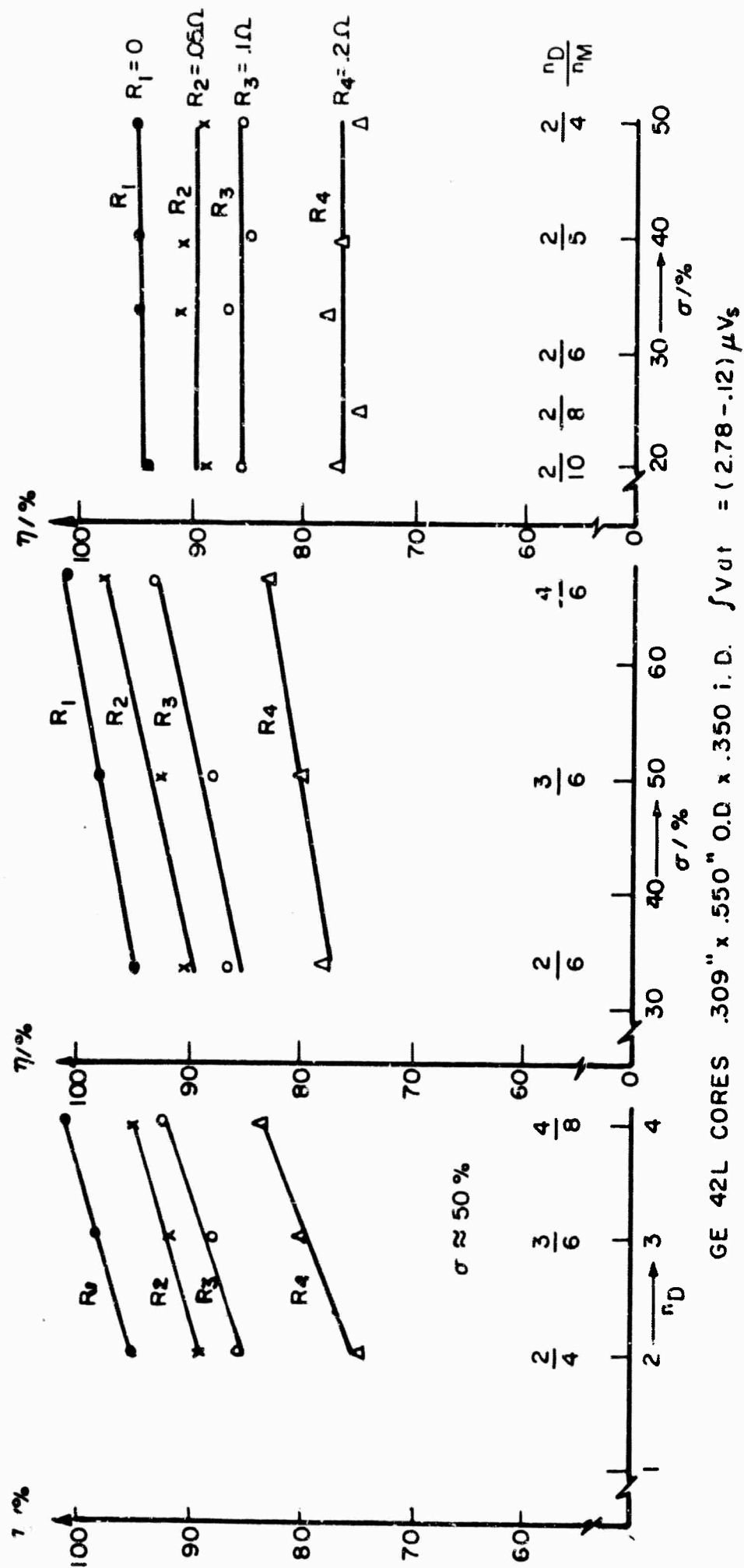
$$\int V dt = 6.52 \mu Vs$$



$$\int V dt = 8.74 \mu Vs$$

#### DRIVER CORE AS PULSE GENERATOR

FIGURE 4



EFFICIENCY OF FLUX-TRANSFER UNDER VARIOUS CONDITIONS

FIGURE 5

driver and the ferrite core. Then  $\Delta\phi_1 = n_D \cdot \frac{1}{2} \cdot \frac{1}{n_M} \cdot \delta\phi = n_D \cdot \delta\phi / n_M$ ,  $P_1 = n_D/n_M$  (assumed). With proportionality constants  $K_1, K_2$

$$\eta = 1 - K_1 \frac{n_M}{n_D} - K_2 \frac{n_M/n_D}{n_D} \quad (9b)$$

Let the ratio "available flux/flux necessary to drive microwave core from  $-M_s$  to  $M_s$ " be the saturation

$$\sigma = \Delta\phi_{1r}/\Delta\phi_r \quad \text{with} \quad \Delta\phi_r = \frac{1}{2} K_2 A \quad (10)$$

In Figure 5a  $\sigma = n/n_M$  is kept constant.  $\eta$  increases with  $n_D$ . In Figure 5b  $n_M$  is kept constant.  $\eta$  increases with  $n_D$ . In Figure 5c  $n_D$  is kept constant.  $\eta$  decreases with increasing  $n_M$ . These results agree with equation (9b).

It is possible to "step-up" the BH loop of the microwave core by adding a constant bit of flux, to whatever state of magnetization has been reached, as long as the core is unsaturated. The transfer efficiency is nearly constant. Figure 6 shows for different driver core sizes and saturation values that the transfer efficiency decreases somewhat with increasing step number. In Figure 7 the fall back flux per step is plotted versus the total flux transferred. The results of Figure 6 indicate that the vertical branch of the BH loop is not parallel to the B-axis. As seen later, this effect can be compensated by an increase in phase shift for a constant flux transferred when stepping up the BH-loop. In a first approximation however, the transfer efficiency can be looked upon as constant until the microwave core is nearly saturated. The transferred flux is then  $\Delta\phi_D = \delta\phi$ . If two bits  $\Delta\phi_1$  and  $\Delta\phi_2$  are switched one after the other, the total flux transferred is

$$\Delta\phi_t^{(1+2)} = \Delta\phi_t^{(1)} + \Delta\phi_t^{(2)} = \Delta\phi_1 + \Delta\phi_2 - 2\delta\phi \quad (11)$$

If the two bits are switched at a time, the transferred flux is

$$\Delta\phi_t^{(1,2)} = \Delta\phi_1 + \Delta\phi_2 - \delta\phi \quad (12)$$

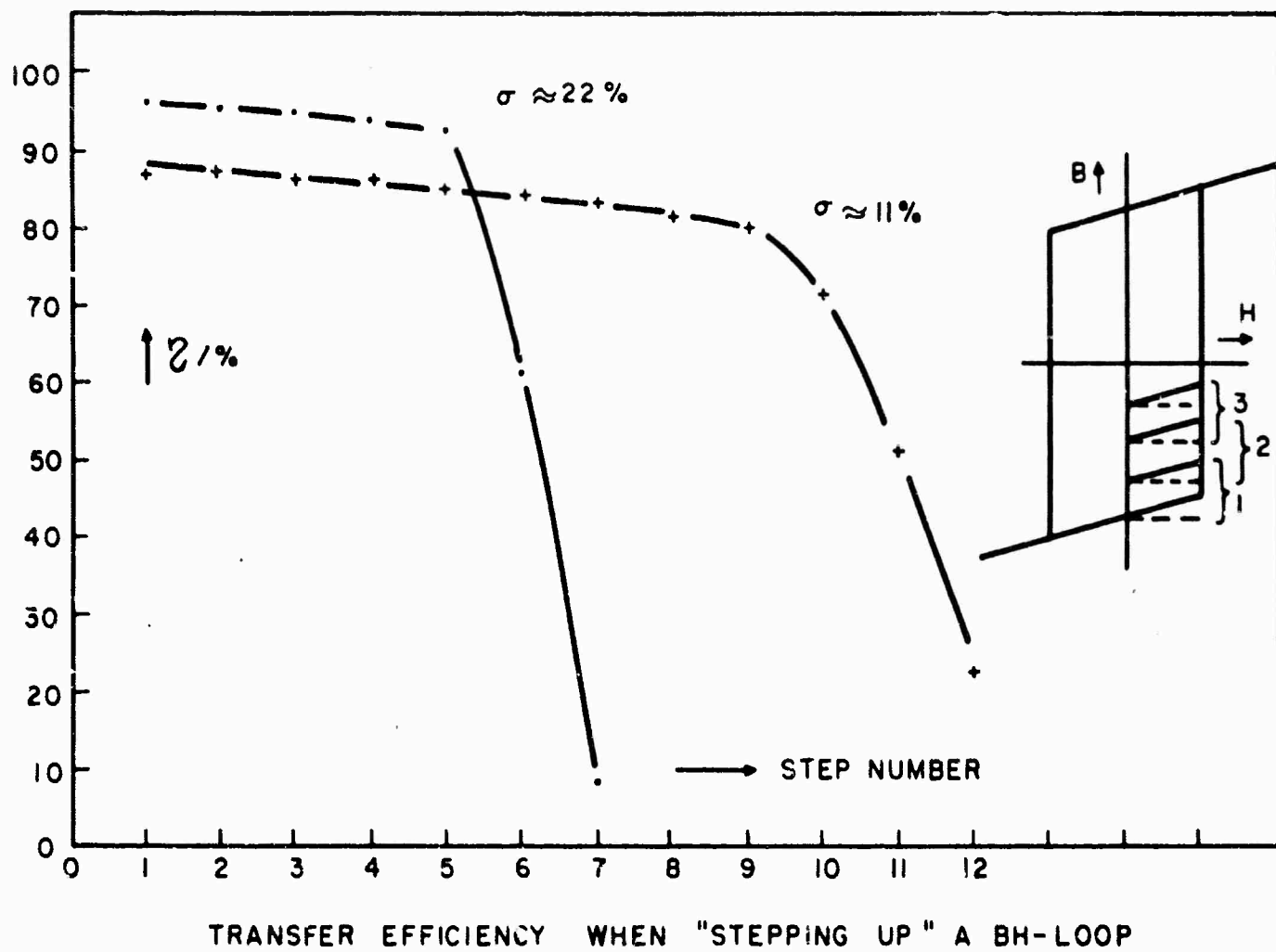
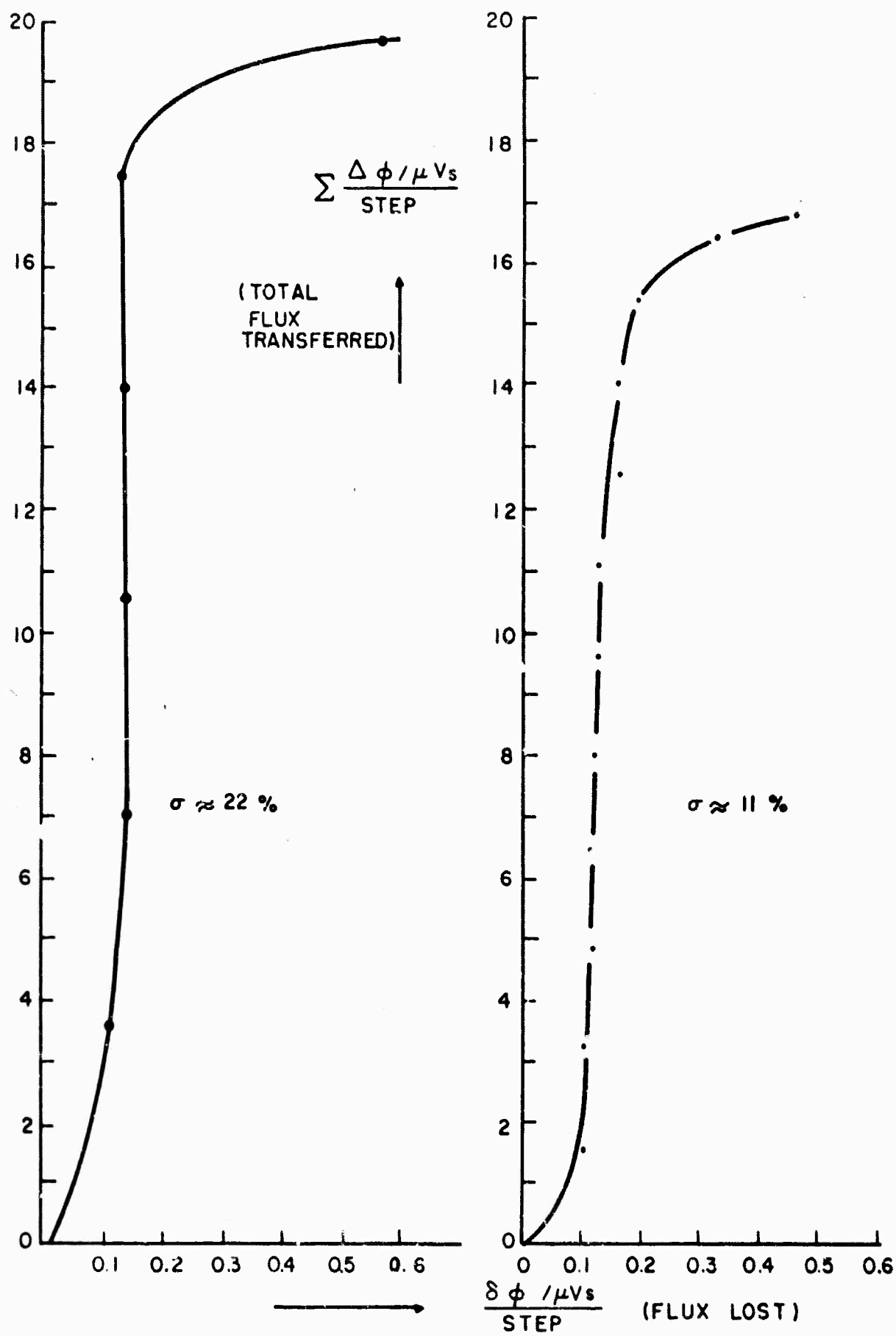


FIGURE 6



FLUX LOST PER STEP WHEN STEPPING UP A BH-LOOP



The difference can be overcome by adding to the  $n$  different bits a compensation flux bit  $\Delta\phi_c = \delta\phi$ . Whenever bits are switched, individually or several at a time, the extra bit is switched at the same time. The transferred flux is then

$$\Delta\phi_t^{(n)} = (\Delta\phi_n + \Delta\phi_c) - \delta\phi = \Delta\phi_n \quad (13)$$

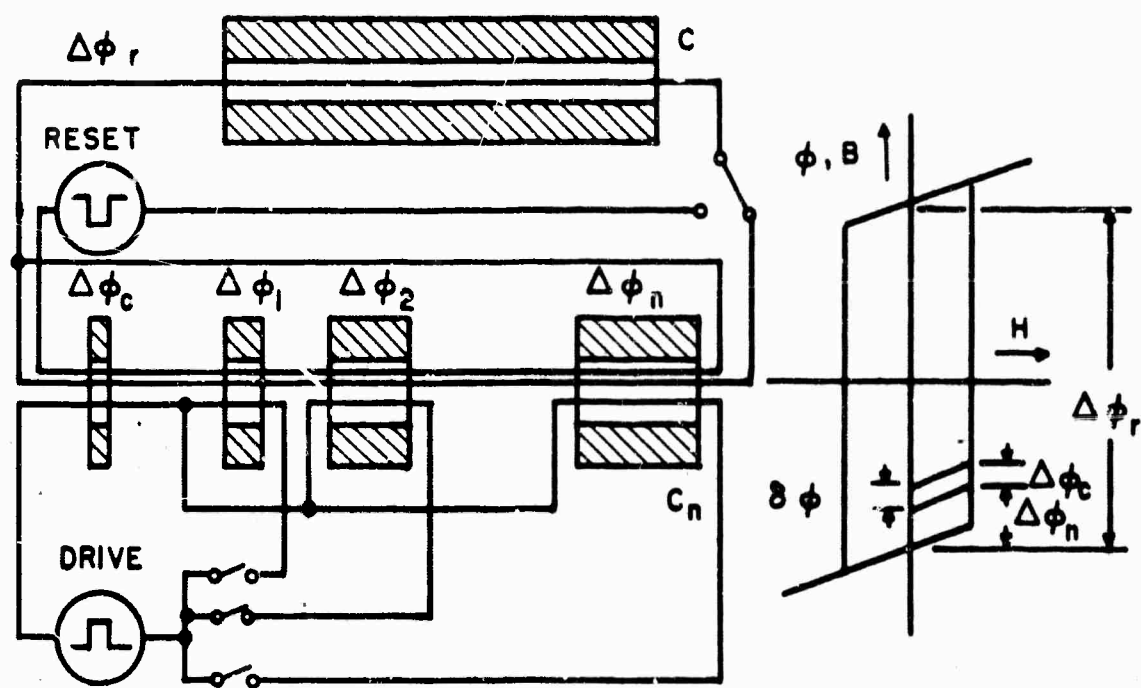
$$\Delta\phi_t^{(1+2)} = (\Delta\phi_1 + \Delta\phi_c) - \delta\phi + (\Delta\phi_2 + \Delta\phi_c) - \delta\phi = \Delta\phi_1 + \Delta\phi_2 \quad (14)$$

$$\Delta\phi_t^{(1,2)} = (\Delta\phi_1 + \Delta\phi_2 + \Delta\phi_c) - \delta\phi = \Delta\phi_1 + \Delta\phi_2 \quad (15)$$

The transfer circuits of Figure 1 and Figure 8 contain the microwave core, the active driver core to achieve a certain flux transfer, and some idle driver cores. The influence of the idle core on the transfer efficiency has been found to be negligible. In a setup as in Figure 1 different numbers of equal ferrite toroids have been used for microwave, driver and idle cores to establish this result. The idle cores are, of course, reset to negative saturation just as the active driver before flux is transferred. Thus, their inductivity is close to zero. The results are summarized in Table 1. The flux transferred to both the microwave and the idle cores has been measured. The low efficiency for  $\sigma = 20\%$  is due to the fact that rather small cores have been used, so that ohmic losses in the transfer circuit represented a considerable fraction of the total losses.

Active cores	$\Phi$ available	Idle cores	Flux Transf.	Microw. cores	Flux transf.	$\sigma$	$\eta$
2	5.32 $\mu$ Vs	0	0 $\mu$ Vs	5	4.46 $\mu$ Vs	40%	83.7%
2	5.32 "	3	<0.02 "	5	4.45 "	40%	83.6%
1	2.66 "	0	0	5	1.66 "	20%	62.5%
1	2.66 "	4	<0.03 "	5	1.66 "	20%	62.5%

TABLE 1. Influence of Idle Cores on Transfer Efficiency



TRANSFER CIRCUIT FOR ADDING FLUX BITS

FIGURE 8

### III. NON-LINEAR LOSSES IN FERRITES

Non-linear losses occur when the rf - magnetic field in the waveguide surpasses a threshold value. Power is then dissipated into the spin wave spectrum. The dominant mechanism is subharmonic generation, where spin-waves of half the frequency of the driving rf-field are excited. This is a "first-order" effect because the spin wave amplitude is proportional to the rf-field amplitude. Via a second mechanism spin-waves can be excited at the frequency of the driving rf field. This effect is of "second order" because the amplitude of these spin waves is proportional to the square of the rf-field amplitude. Whenever the subharmonic generation takes place it occurs at such low rf-field amplitudes that the second order effect may be neglected. Assuming only subharmonic generation, E. McKinney has derived threshold field strengths for ferrites of arbitrary shape at remanence. This treatise is given fully as Part II of this report. It includes the results of other workers. McKinney has to normalize the threshold field strength to a spinwave line width or relaxation time. Were this line width known, it would be possible to predict the threshold field strength for ferrites as a function of frequency, and for devices where the field distribution is known one could also predict the threshold power. So far it was not possible to measure the relaxation time accurately at one frequency only, even less to determine whether the relaxation time is a function of frequency, ferrite shape, driving field strength or other parameters.

By testing ferrites of equal size and shape in equal test devices at 6.5 GHz, relative power handling capabilities could be determined. Two test series were performed: In the first one a coaxial line was completely filled with ferrite. The diameters were chosen so that only TEM modes should propagate. This test device was matched into 50 ohm lines via step transformers at both ends. The VSWR was below 1.5 in all cases, corresponding to a reflection loss  $< 0.17$  dB or a power reflection of less than 4%. In the second series a ferrite slab was placed in the center plane of a rectangular

waveguide parallel to the narrow wall and extending over the full height of the guide. The filling factor was five percent. The ferrite was tapered at both ends to secure a VSWR < 1.5. The rf magnetic field amplitude was calculated under the assumption that the relative ferrite permeability  $\mu_f = 1$  and the relative ferrite dielectric constant  $\epsilon_f = 11, \dots 16$ , as given by the manufacturer. In the coaxial line the strongest rf-magnetic field occurs at the inner conductor. With

$$P = ZI^2 = Z\pi^2 D^2 H^2 \quad (16)$$

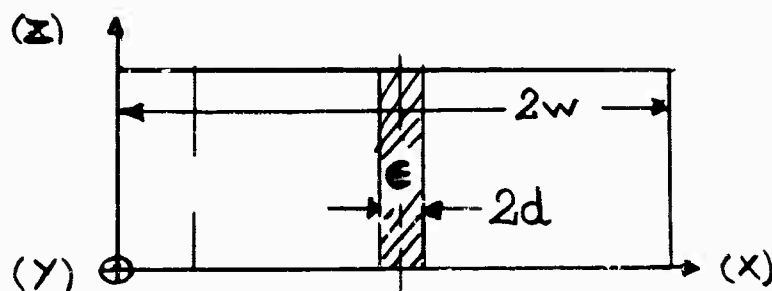
where  $P$  = peak power,  $I$  = current in inner conductor,  $D$  = diameter of inner conductor,  $Z$  = line impedance,  $H$  = rf-magnetic field. One obtains with  $\pi^2 \approx 10$

$$\frac{H}{A/cm} = 10 \frac{D}{cm} \sqrt{\frac{P/KW}{Z/ohm}} \quad (17)$$

In the rectangular waveguide (WR 137) at the test frequency of 6.5 GHz the field distribution outside the ferrite slab follows a hyperbolic function and one obtains with the dimensions sketched below

$$P_y = \oint (E_z \times H_x) dA = 2b \int_0^w (E_z \times H_x) dx \quad (18)$$

$$\equiv 2wb \frac{\omega\mu_0}{k} H^2 \left[ \int_0^{1-\delta} \frac{\sinh^2 P\varphi}{D^2} d\varphi + \int_{1-\delta}^1 \sin^2 (Q\varphi + \theta) d\varphi \right]$$



where

$$\begin{aligned}
 k &\hat{=} \text{longitudinal propagation constant} \\
 P/w &\hat{=} \text{transverse propagation constant in empty part of guide} \\
 Q/w &\hat{=} \text{transverse propagation in dielectric} \\
 l/D &\hat{=} \text{relative amplitude in empty part of guide} \\
 l &\hat{=} \text{relative amplitude in dielectric} \\
 \delta &= d/w \\
 \varphi &= x/w
 \end{aligned}$$

The full theory is given as Part III of this report. For the given frequency and waveguide ( $2w = 1.372$  inches,  $b = .622$  inches) one obtains finally

$$\frac{H}{A/cm} = N \sqrt{\frac{P}{kW}} \quad (19)$$

where  $N$  is a function of the ferrite dielectric constant

$\epsilon_f$	1	8	9	12	14	16
$N$	.925	1.055	1.108	1.130	1.162	1.180

The test results are summarized in Figure 9. The coaxial line test leads to considerably higher values for the threshold field than the wave guide test does. The reason is not yet understood. The spin wave spectra excited in both test series may not be equal. Other possible explanations are for the coaxial line test:

The impedance is other than calculated with  $\epsilon_f$

Several modes exist (partly due to air gaps between ferrite and conductors)

The rf-magnetic field strength varies considerably over the ferrite wall thickness.



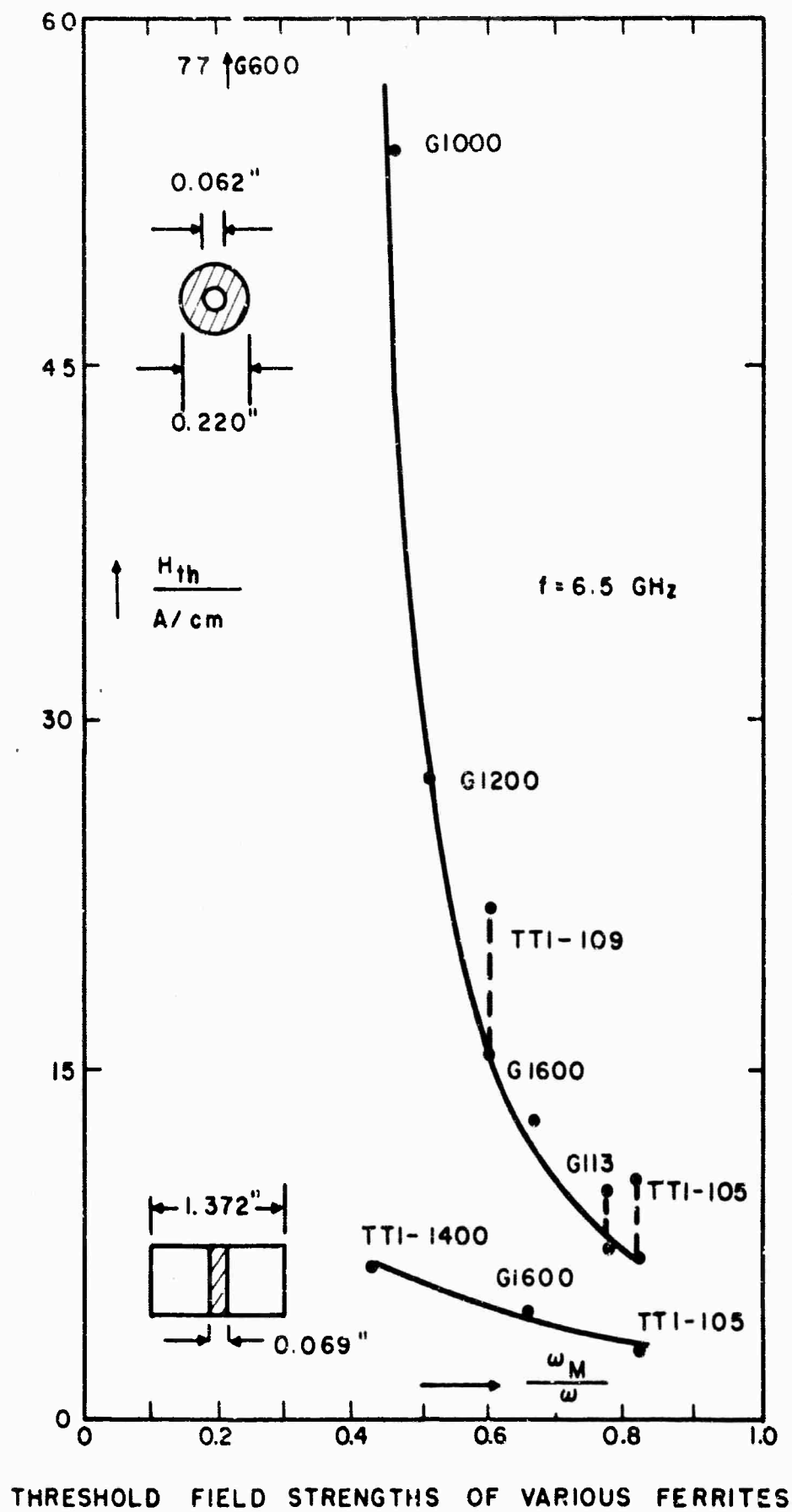
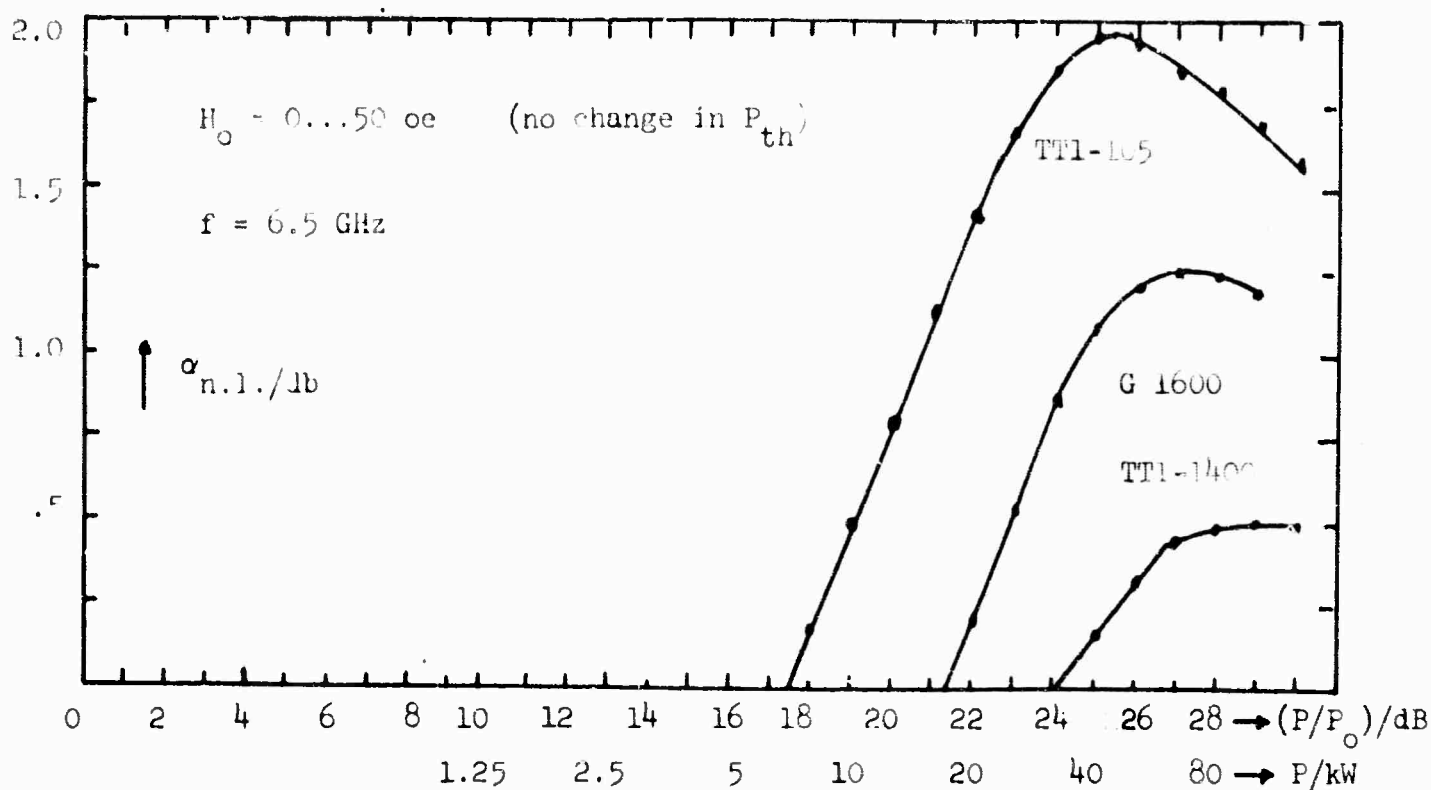
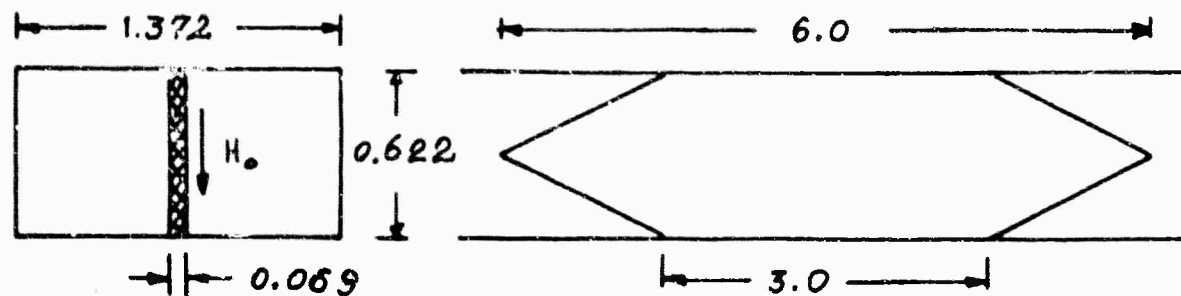
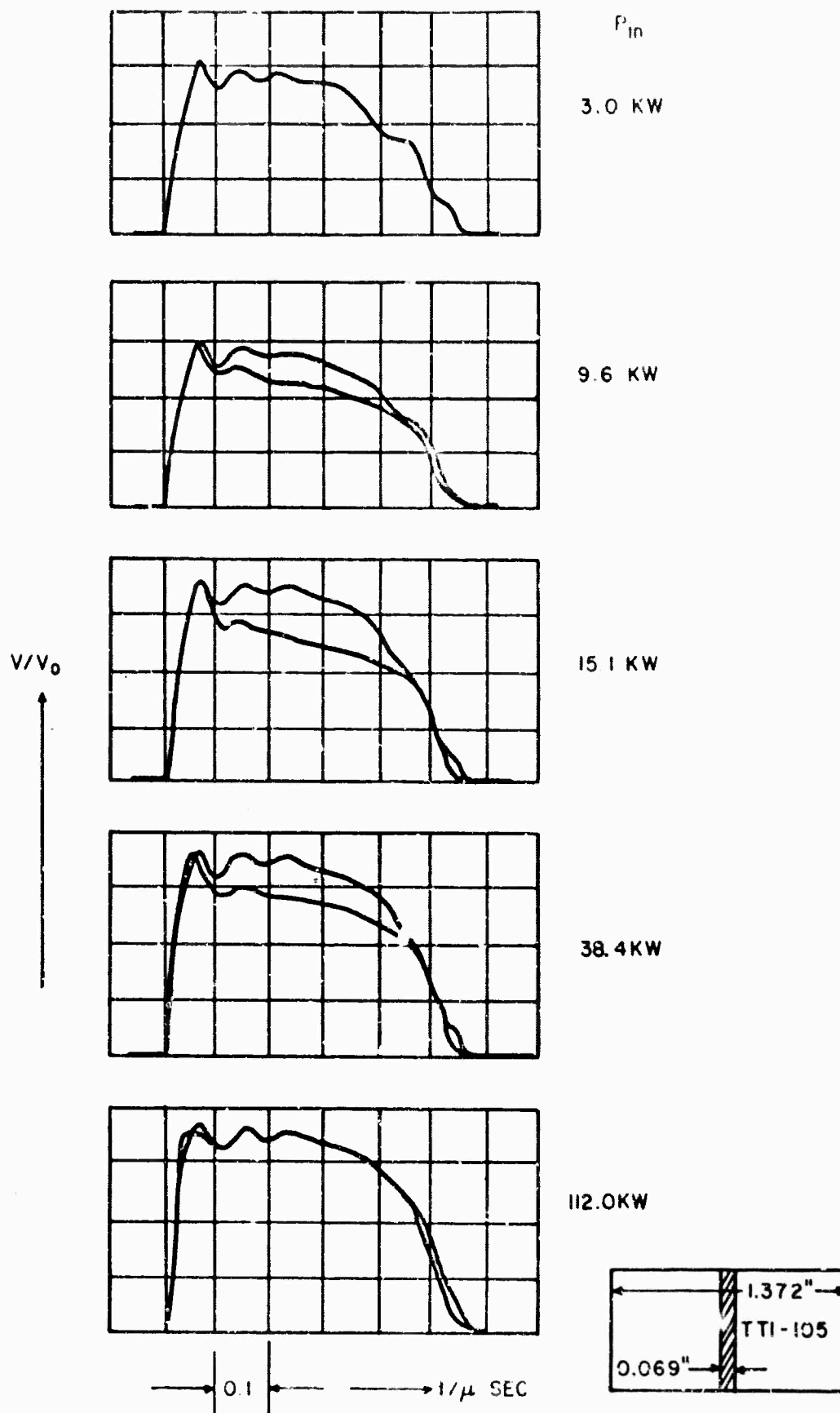


FIGURE 9

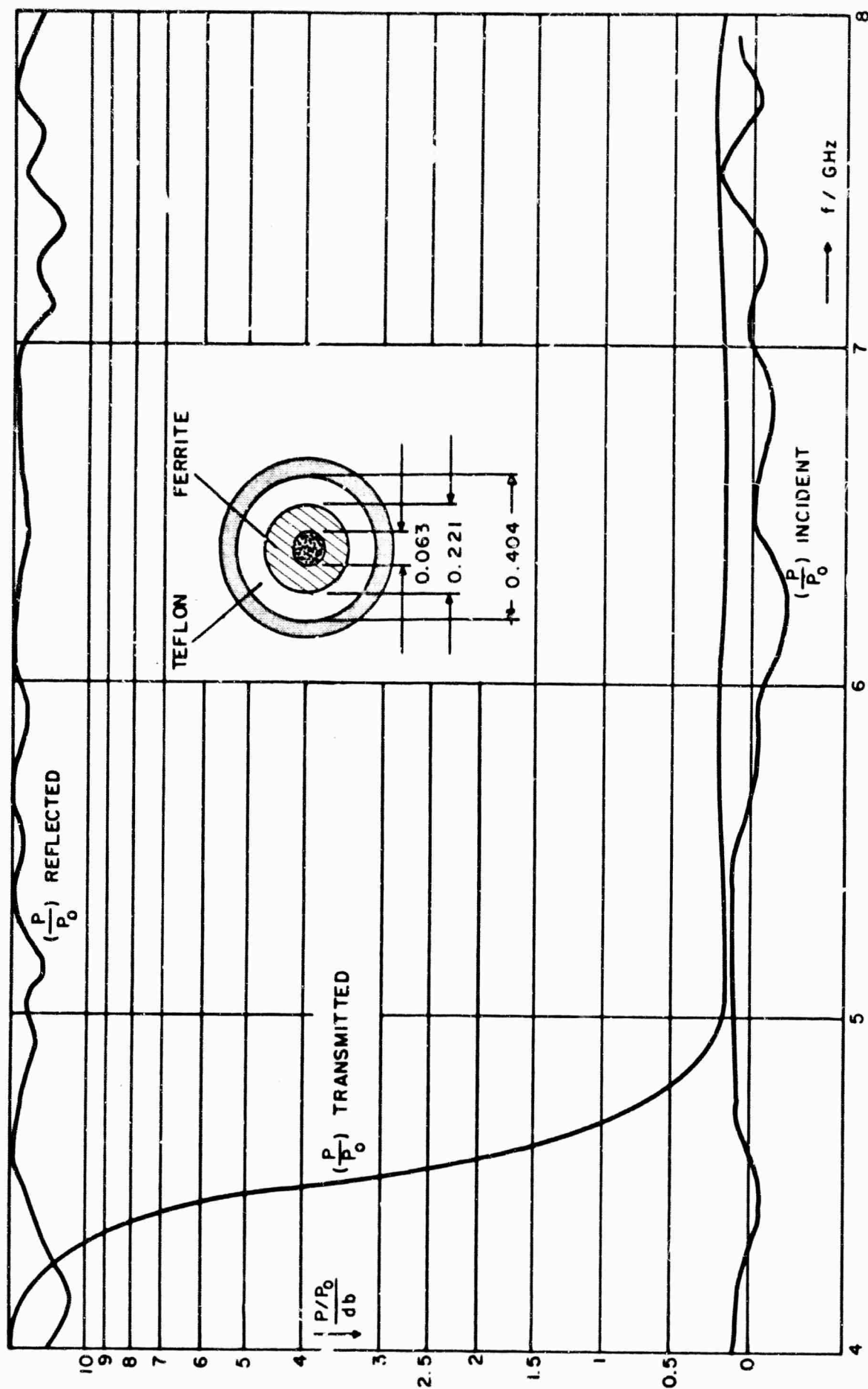


Ferrite (TransTech)	Listed $M_s/G$	$f_m/f$	$P_{th}/kW$	
G 600	515	0.22	>125	Voltage breakdown
TT1-414	660	0.28	>125	
TT1-900	(700)	(0.3)	>125	(...) not listed, assumed
TT1-109	860	0.37	>125	
TT1-1400	(1000)	(0.43)	32	(...) not listed, assumed
G1600	1165	0.5	18	
TT1-105	1380	0.58	7	

FIGURE 10. Ferrites at High Peak Powers



DEFORMATION OF RF-PULS BY NON-LINEAR ATTENUATION AT 6.5 GHz  
 UPPER TRACE; INPUT PULS, LOWER TRACE; OUTPUT PULS



MAIN RESONANCE LOSS IN 50 ohms COAXIAL LINE PARTIALLY LOADED WITH FERRITE (TRANS TECH. G 113)

The average threshold power in the waveguide test could be determined accurately, because the non-linear loss/power characteristic is a straight line in a double logarithmic scale as shown in Figure 10. If the ferrite would not warm up, this line would end abruptly at a maximum power-independent loss. Such a maximum loss has been maintained over a large power range when the ferrite was kept cool. Figure 10 shows a peak in loss due to the warming up of the ferrite and the subsequent change in saturation magnetization. The onset of non-linear losses in the coaxial line tests was not as well defined. The losses increased to a small value over a large power range before they followed the straight line. Besides the results were not as reproducible as in the waveguide tests.

Figure 11 shows the input pulse and the deformed output pulse at various power levels. The reduction in loss at high powers shows up clearly.

To determine the threshold power it has been assumed that the input pulse is rectangular and 0.42  $\mu$ sec wide. Assuming an exponential growth of the spin waves it should be possible to find the relaxation time from the decay of the high peak of the output pulse.

With the poor input pulse shape, however, this does not yield accurate results. The spin wave line width appears to be between 0.5 to 2 A/cm.

The resonance frequency has been taken from the manufacturers data or from the onset of losses at lower frequencies in a coaxial line filled with ferrite, as illustrated in Figure 12, or from measuring the flux available from a toroid. In any case the accuracy may be  $\pm 5\%$ . Plotting the threshold fields versus  $\omega_m/\omega$ , as done in Figure 9 and assuming a linewidth of 1 A/cm, one finds that the waveguide test data come closer to the theoretical results than the coaxial line test data do.

In any case the power handling capability increases as the saturation magnetization decreases.

#### IV. RECIPROCAL TEM MODE PHASE SHIFTER

The "one shell" and the "many shells" model for the magnetization of ferrite toroids - mentioned under "Control of Flux Transfer" - leads to an interesting consequence. The "many shells" model forbids, while the "one shell" model allows a reciprocal, latched ferrite phase shifter using the TEM modes in a coaxial line. According to the "many shells" model an alternating current with decreasing amplitudes applied to the inner conductor of the coaxial line would produce a sequence of circumferentially magnetized ferrite shells with alternating directions of magnetization. In any such sequence no microwave interaction between the ferrite and rf magnetic field of a TEM mode would occur. The effective relative permeability for this mode would always be one. In a "one shell" model the magnetization would be homogeneous throughout the ferrite wall thickness and could vary from complete alignment in one direction to random spin orientation to complete alignment in the other circumferential direction, depending on the driving conditions. This model allows a reciprocal latched coaxial phase shifter. The electrical length of the device depends on the degree of dipole alignment in the ferrite, not on the circumferential direction. In a coaxial line filled with ferrite (TTL-109,  $L = 4$  inches,  $D_o = .221$  inches,  $D_i = .063$  inches), a differential phase shift has been measured between a) the demagnetized state, achieved by sending an a.c. pulse with decreasing amplitude through the center conductor and, b) the state of partially circumferential magnetization (either direction) achieved by sending a d.c. pulse through the inner conductor.

The reciprocal phase shift could be controlled by the amplitude of the d.c. current pulse or by controlled flux transfer. At about 5 GHz the reciprocal phase shift could be stepped up to approximately 20, 40, 60, and 80 degrees. The coaxial line used in this experiment was built for ferrite tests at high peak powers at 6.5 GHz. The phase shift experiment was only a feasibility demonstration. The demagnetization was done at 60 Hz. The switching properties of nonreciprocal phase shifters - investigated under the present contract - indicate that demagnetization at a few kHz should be possible.

Stepping up the B-H loop from one state of circumferential magnetization to the demagnetized state to the opposite state of magnetization gave a smaller maximum reciprocal phase shift than starting from the a.c. demagnetized state.

## V. DIFFERENTIAL PHASE SHIFTERS

Three types of nonreciprocal phase shifters are shown in Figure 13. All of them use as the nonreciprocal element ferrite toroids, which are circumferentially magnetized and work at remanence. In all of them the phase shift is controlled by partial magnetization, achieved through controlled flux transfer.

The "normal" helix type phase shifter has been developed and described by H. A. Hair.<sup>(1)</sup> The "inverted" helix type phase shifter has been found less effective<sup>(2)</sup> than the normal one.

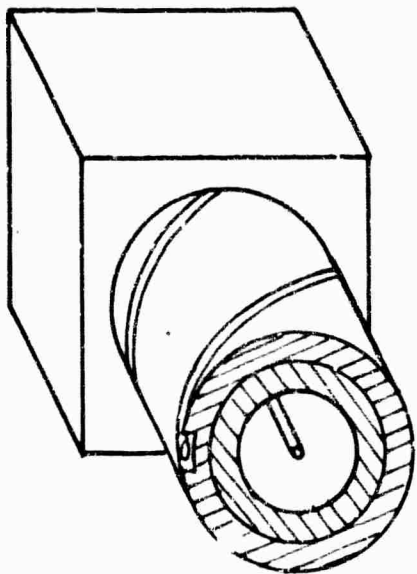
The slabline type phase shifter and the waveguide type phase shifter with ferrite toroids were used in the experiments described further down. Rectangular ferrite tubes were used in the first waveguide type phase shifter by Levey and Silber<sup>(3)</sup> and the improved versions by Blevins,<sup>(4)</sup> and Taft and Sweeney.<sup>(5)</sup>

The direction of magnetization in the phase shifters considered is in a plane transverse to the direction of propagation. Differential phase shift can be obtained when circularly polarized rf fields exist in a plane transverse to the direction of magnetization. Hybrid modes, which do not occur in empty guides, and TE modes, which have counterparts in empty guides, fulfill this condition. The tensor permeability of the ferrite makes the exact computation of the propagation constants in these guides so difficult, that so far no exact solution exists. An analysis has to be based on more or less crude models. In simple configurations exact solutions may be found for dielectrically loaded guides. Perturbation calculations can then take care of thin magnetized ferrites.<sup>(6)</sup> But even in the complex configurations investigated, the dielectrically loaded waveguide model analyzed in Part III of this report offers good explanations.

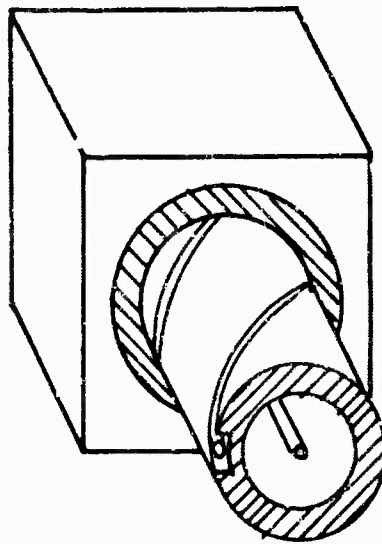
In the phase shifters considered here it has been attempted to achieve a proportionality between the differential phase shift and the flux transferred to the phase shifter. The ferrite was driven between its maximum negative remanent state of magnetization and various other remanent states between this one and the maximum positive remanence. All tests were performed with Trans-Tech ferrites.



"normal"

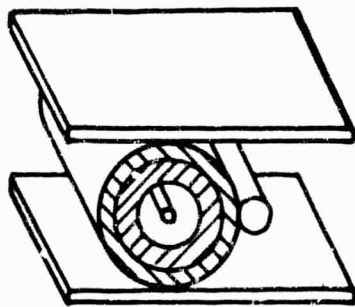


"inverted"

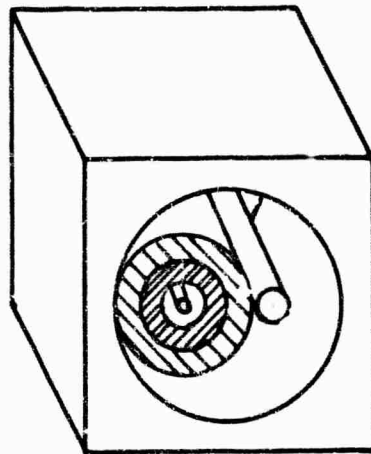


helix type

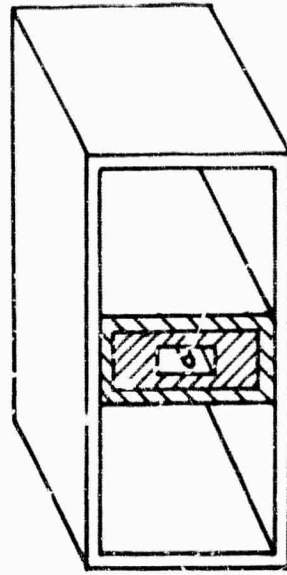
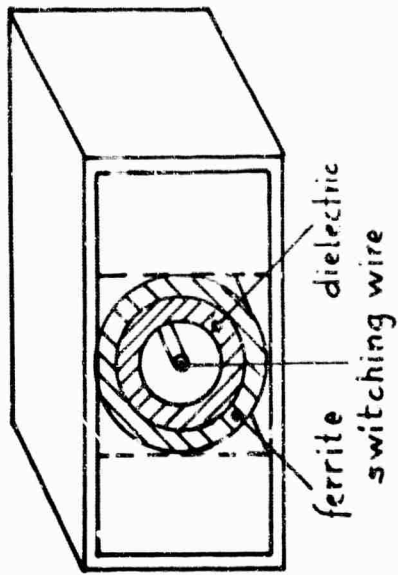
"open"



"shielded"



slabline type



waveguide type

LATCHED PHASE SHIFTERS

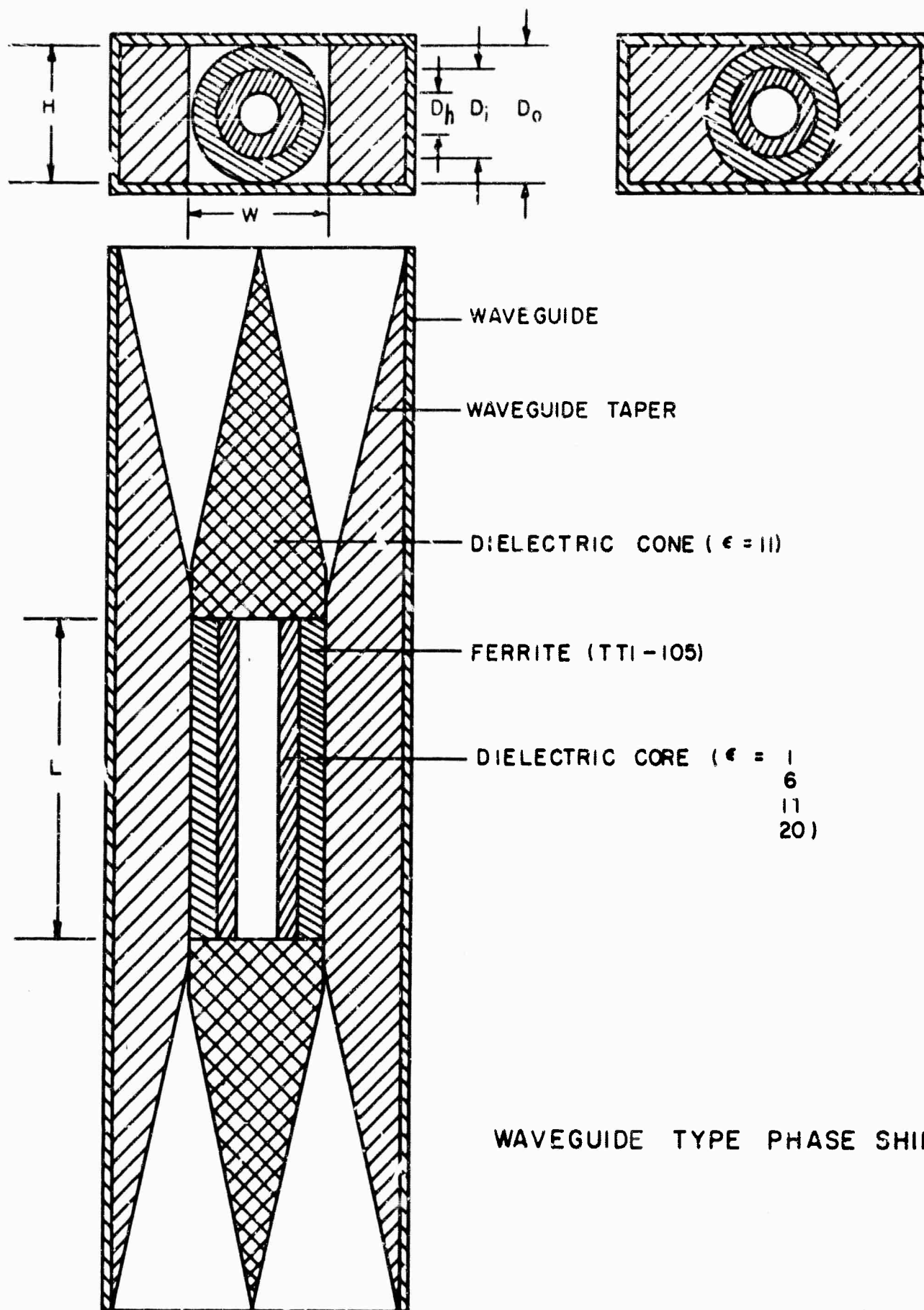
FIGURE 13

## V.1 Waveguide Type Phase Shifters

All the phase shifters investigated use round ferrite cylinders. Previous work done in this laboratory indicated, that they should perform at least as well - as far as figure of merit and bandwidth are concerned - as phase shifters using rectangular ferrite cylinders.

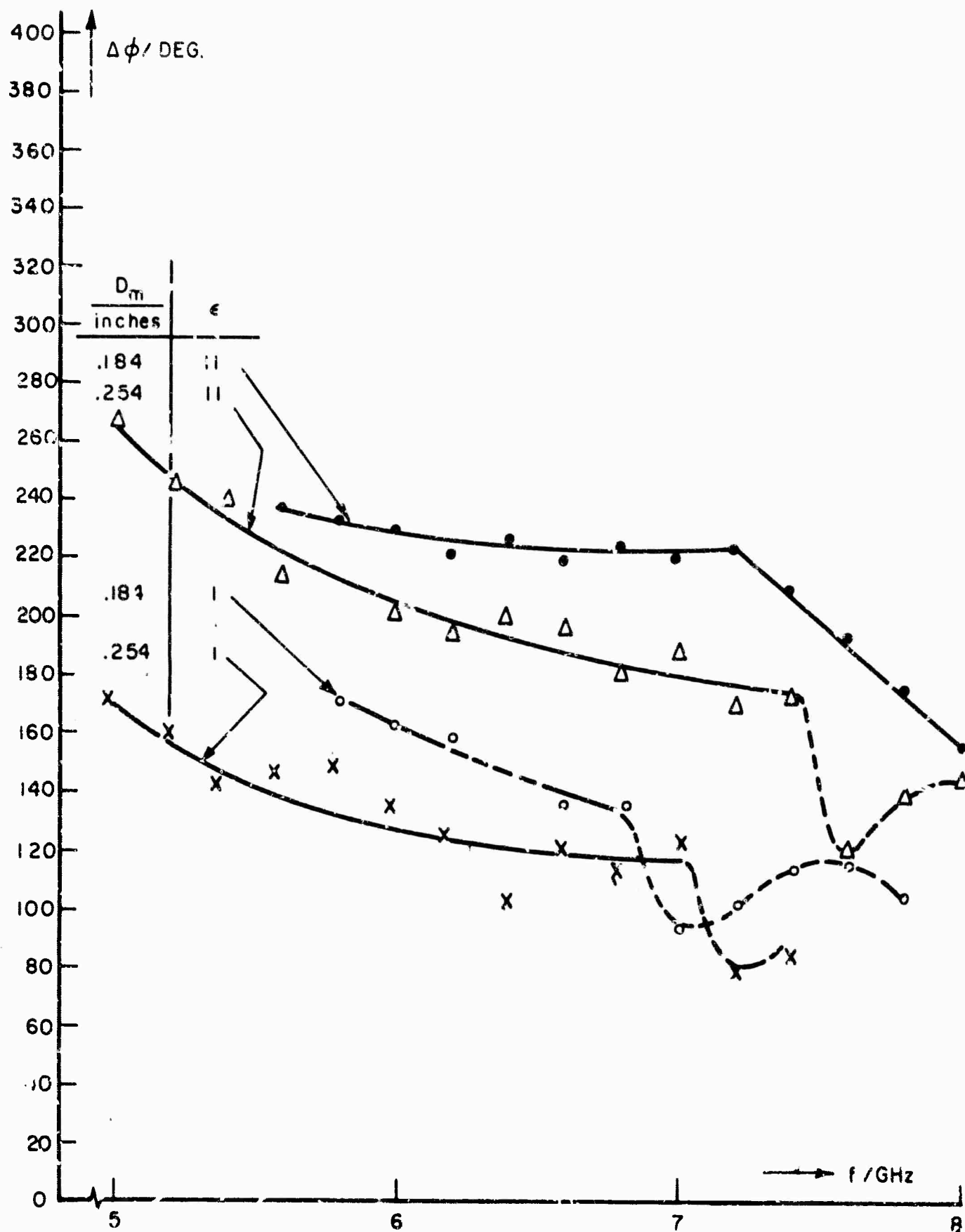
The first experiments were done with a TTL-105 ferrite toroid ( $M_s \approx 17 \mu\text{Vs}/\text{cm}^2$ ) with .662" outer diameter in a standard WR 137 waveguide. With tapered slabs the guide width was reduced. The final step was to go to a circular guide of the ferrite diameter. The device is shown in Figure 14. Other parameters varied were the ferrite wall thickness and the core dielectric constant. The results are shown in Figures 15, 16, 17, (phase shift versus frequency) and Figure 18 (phase shift when stepping up the B-H loop). The results show that thin wall ferrites with proper dielectric cores can produce phase shift nearly as large as achieved with thick wall ferrites and that the phase shift versus frequency curve is the flatter the wider the guide. The "phase shift per step" characteristic for stepping up a BH-loop could be improved when the guide cross-section was changed as shown in Figure 19. The guide height is reduced to less than the ferrite toroid diameter and the toroid rests in grooves in the top and bottom wall of the guide. The threshold power with this device was .8 KW.

Since the threshold field strengths for ferrites generally increases with decreasing magnetization the ferrite TTL-1400 ( $M_s \approx 10 \mu\text{Vs}/\text{cm}^2$ ) was chosen for further experiments. Figure 20 shows first results with this ferrite. Reducing the ferrite diameter reduced the slope of the phase shift versus frequency characteristic. Essentially, constant differential phase shift over a bandwidth of 2.5 GHz, centered at 6.2 GHz was achieved with a filling factor of about 31%. With increasing filling factor the slope becomes negative, as seen in the illustrations of this report; with decreasing filling factor the slope becomes positive, as shown for guides with rectangular ferrite toroids by Taft and Sweeney.<sup>(5)</sup> The optimum filling factor for constant phase shift over a broad band seems to be in the neighborhood of 35%. In Figure 22 it is seen that the guide height has no great



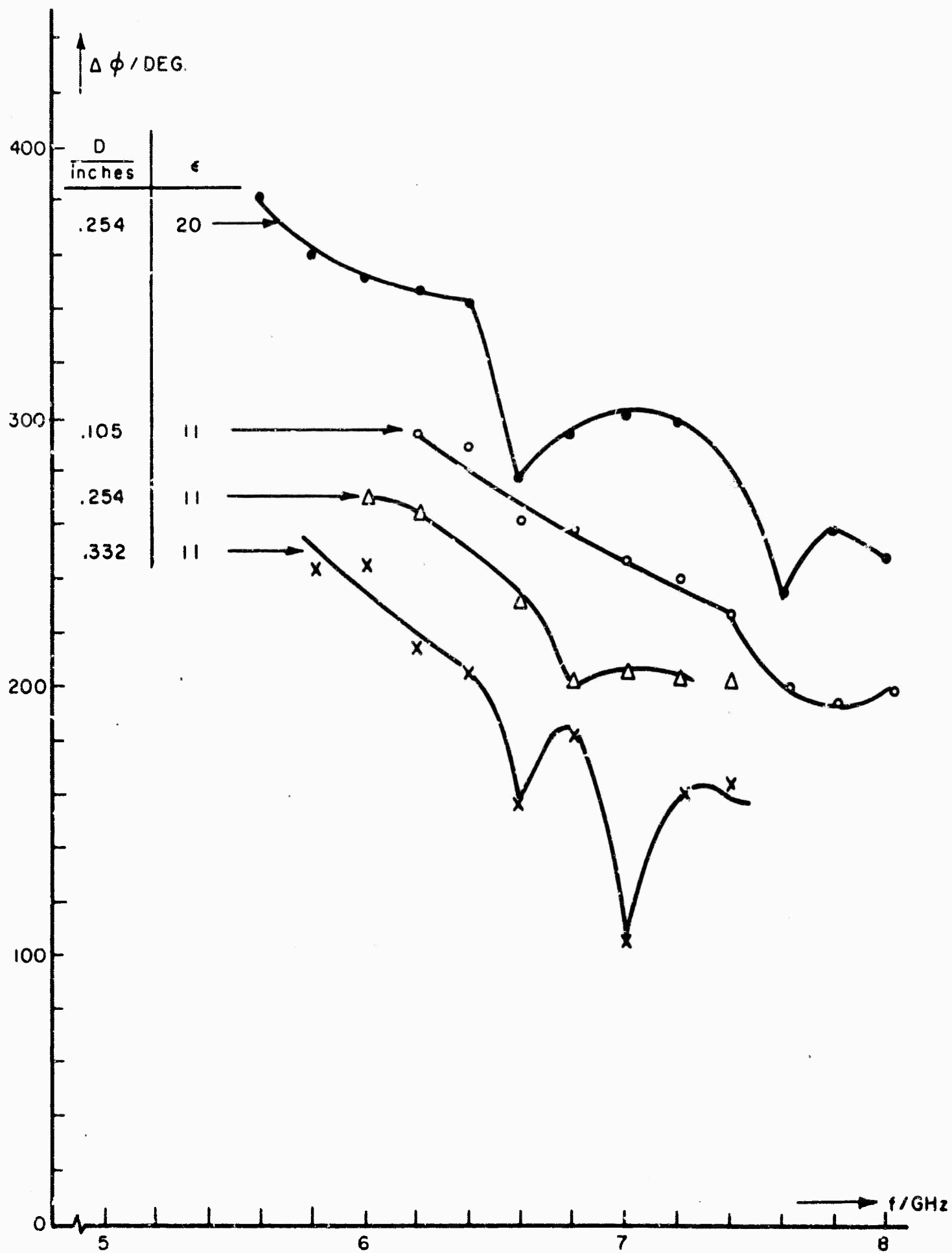
WAVEGUIDE TYPE PHASE SHIFTER

FIGURE 14



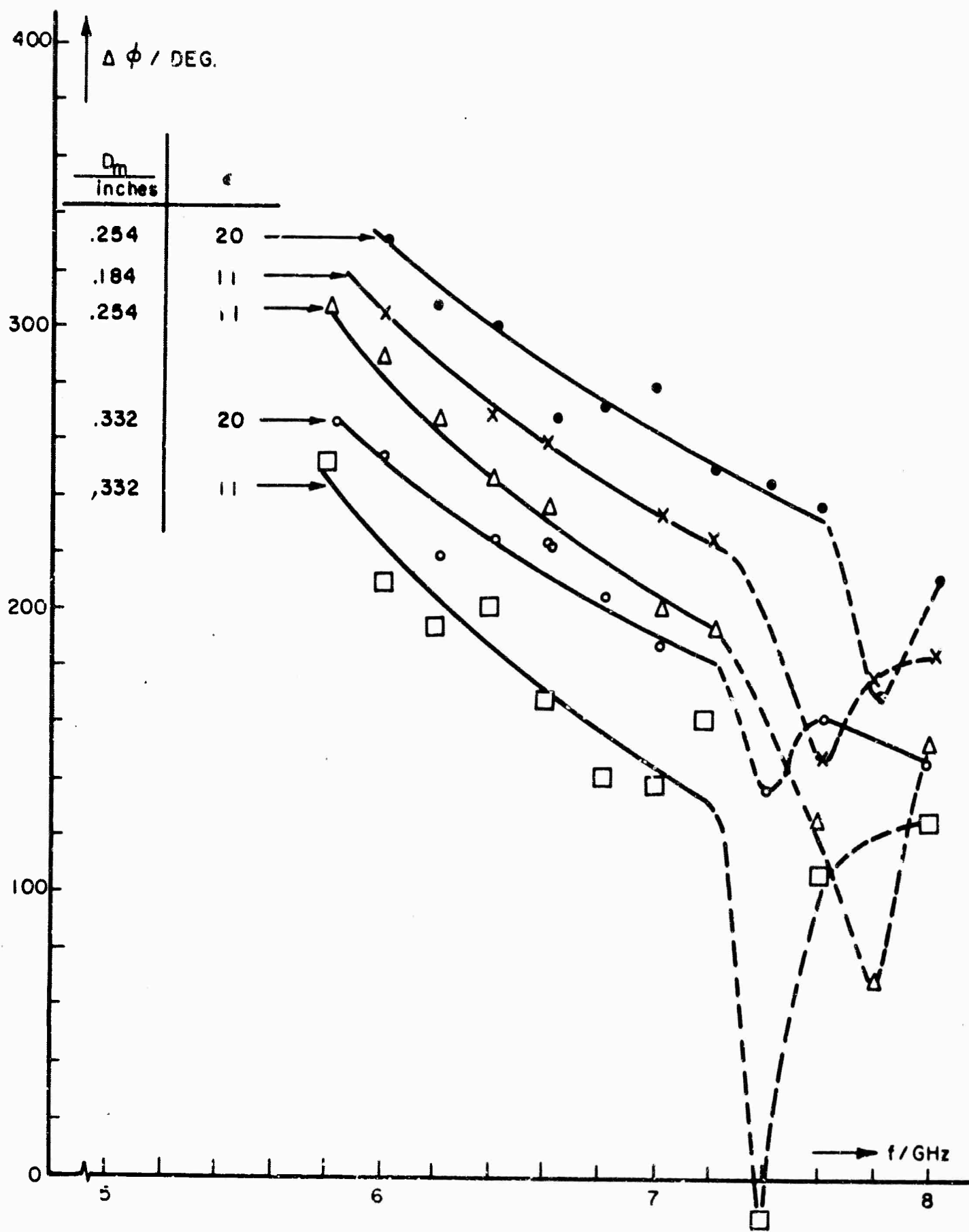
PHASESHIFT IN WAVEGUIDE WITH REDUCED WIDTH ( $W = .782''$ )

FIGURE 15



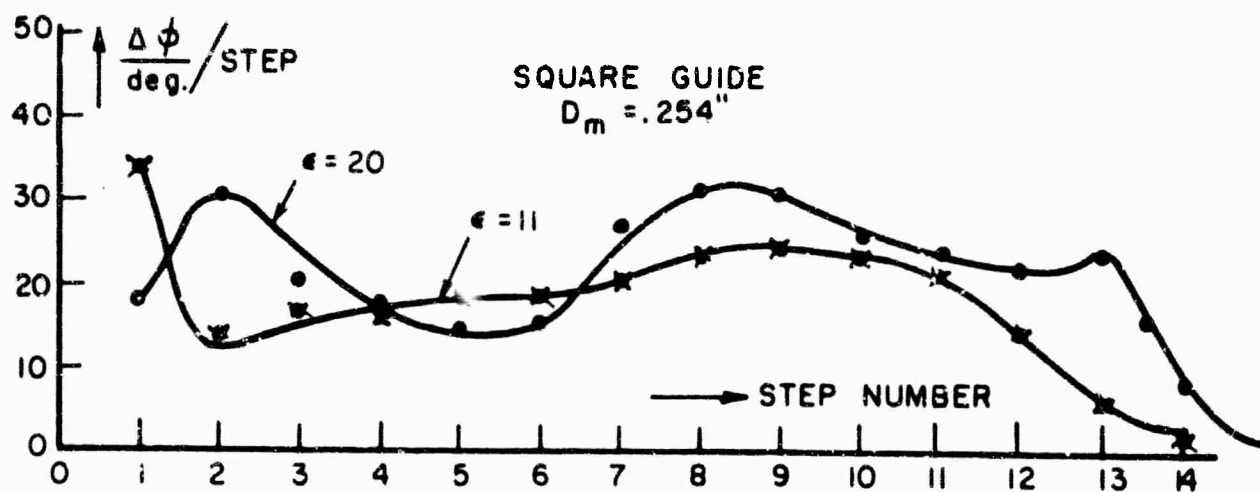
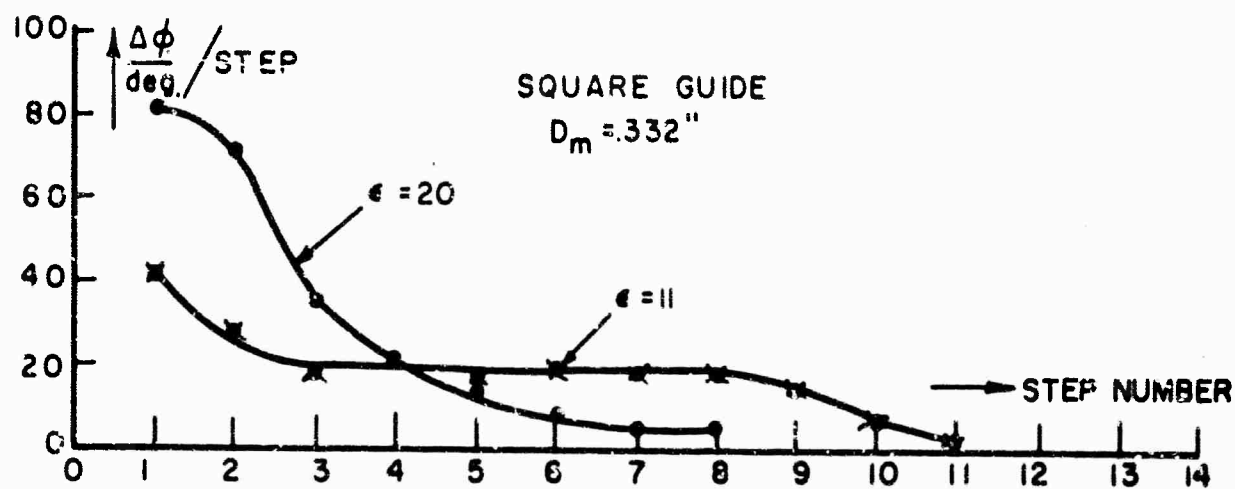
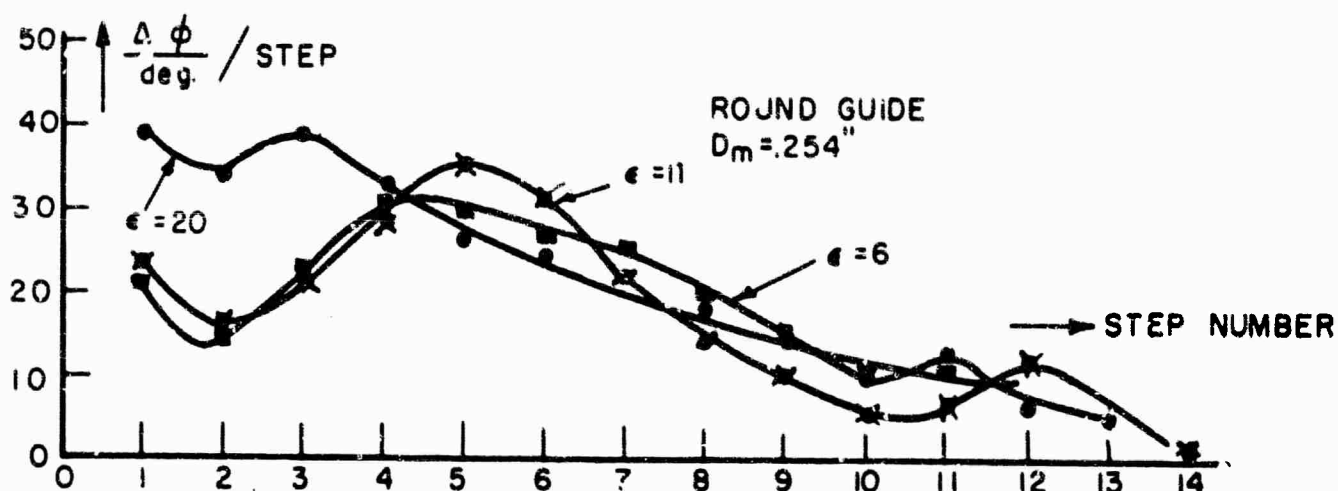
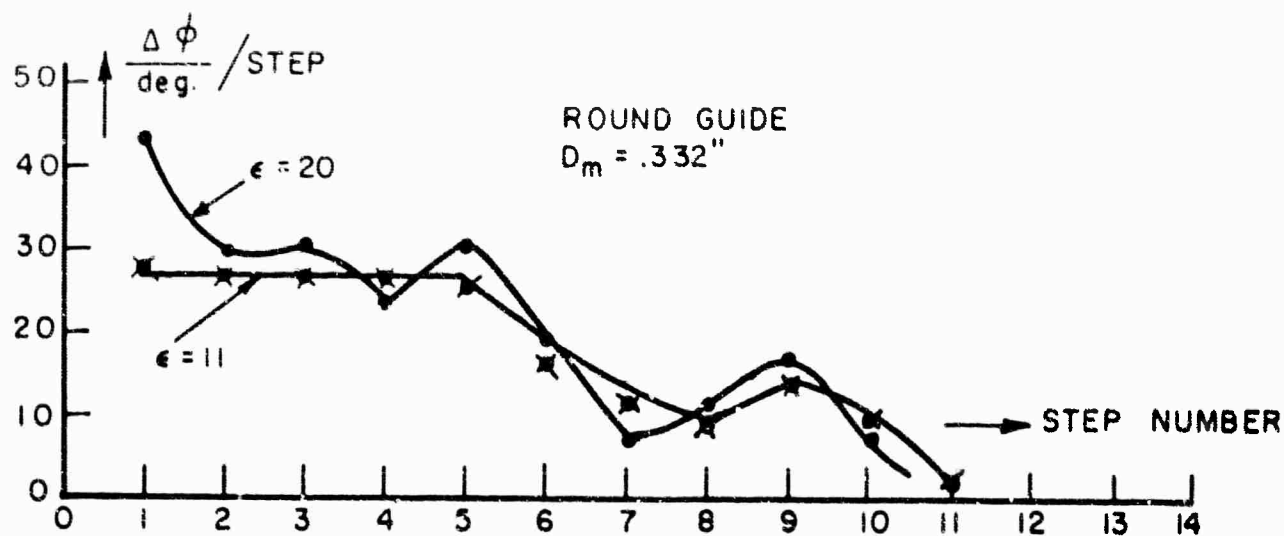
PHASESHIFT IN SQUARE WAVEGUIDE ( $W = .622''$ )

FIGURE 16



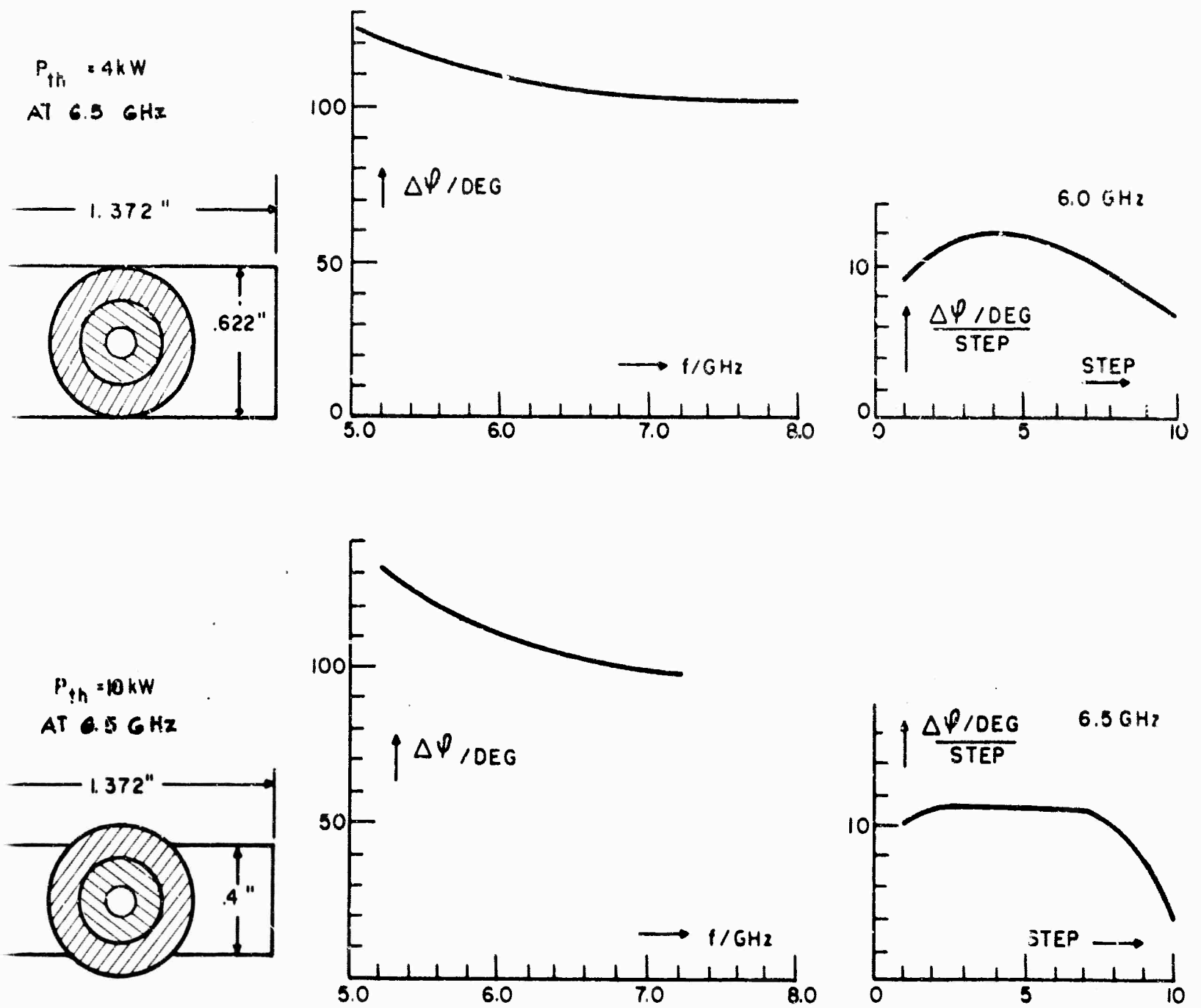
PHASE SHIFT IN ROUND WAVEGUIDE ( $D = .622''$ )

FIGURE 17



PHASESHIFT PER STEP AT 6.4 GHz

FIGURE 18

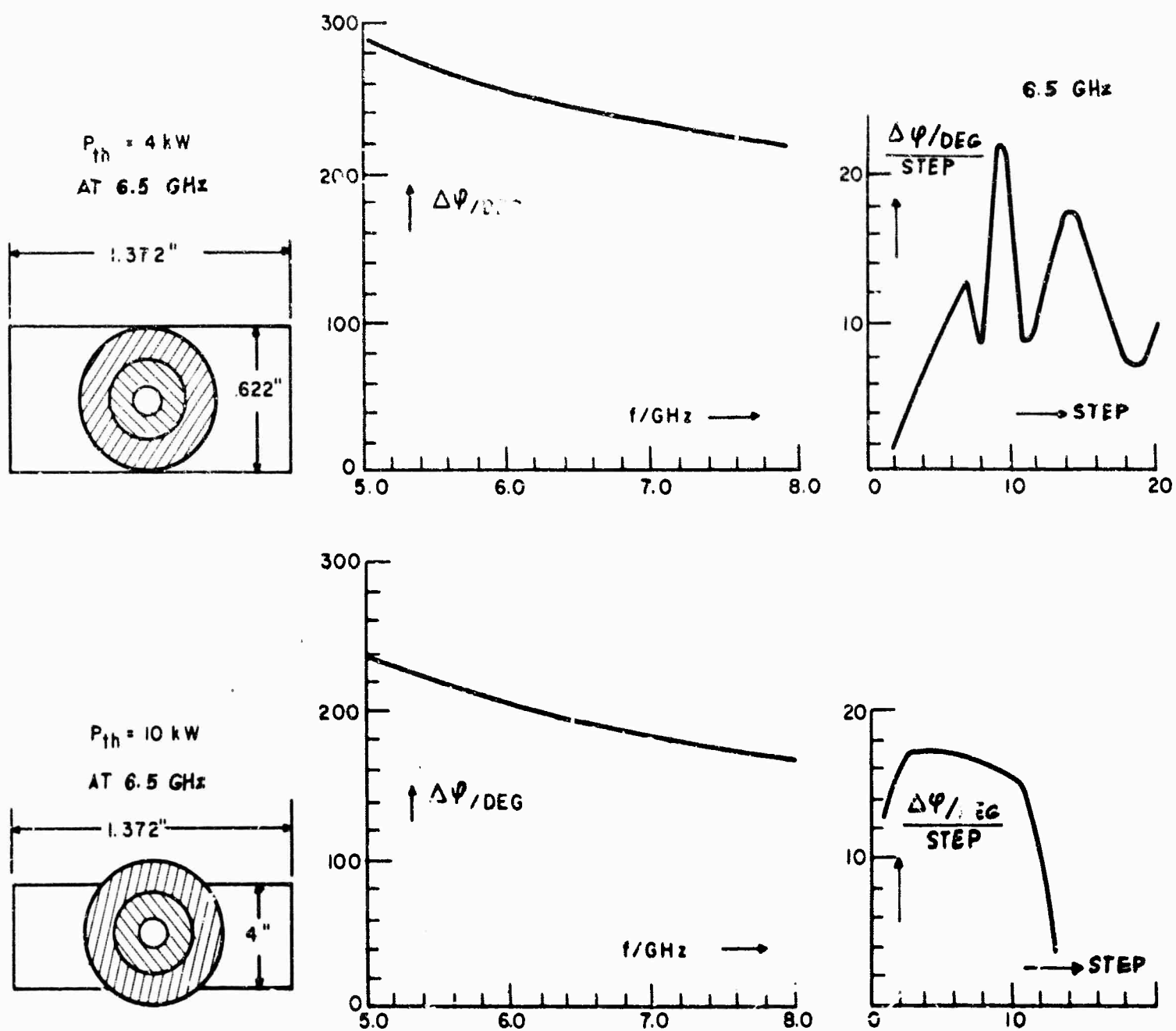


FERRITE: TTI-105,  $D_o = 0.622"$ ,  $D_i = 0.437"$ ,  $l = 1.5"$  CORE: E-11

### WAVEGUIDE TYPE PHASE SHIFTER

FIGURE 19

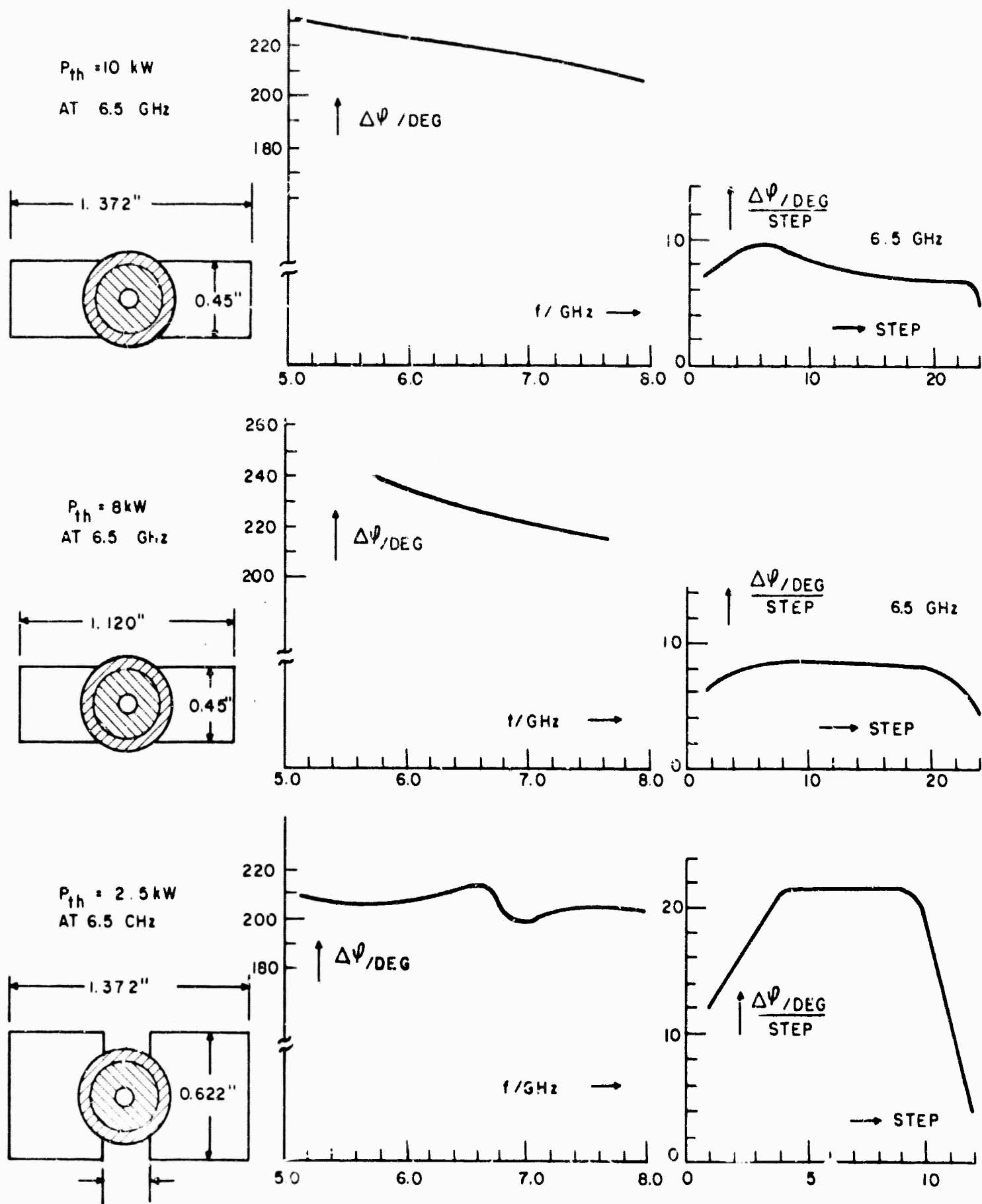




FERRITE: TTI-1400,  $D_o = 0.622"$ ,  $D_i = 0.380"$ ,  $l = 3.0"$ , CORE:  $\epsilon = 11$

### WAVEGUIDE TYPE PHASE SHIFTER

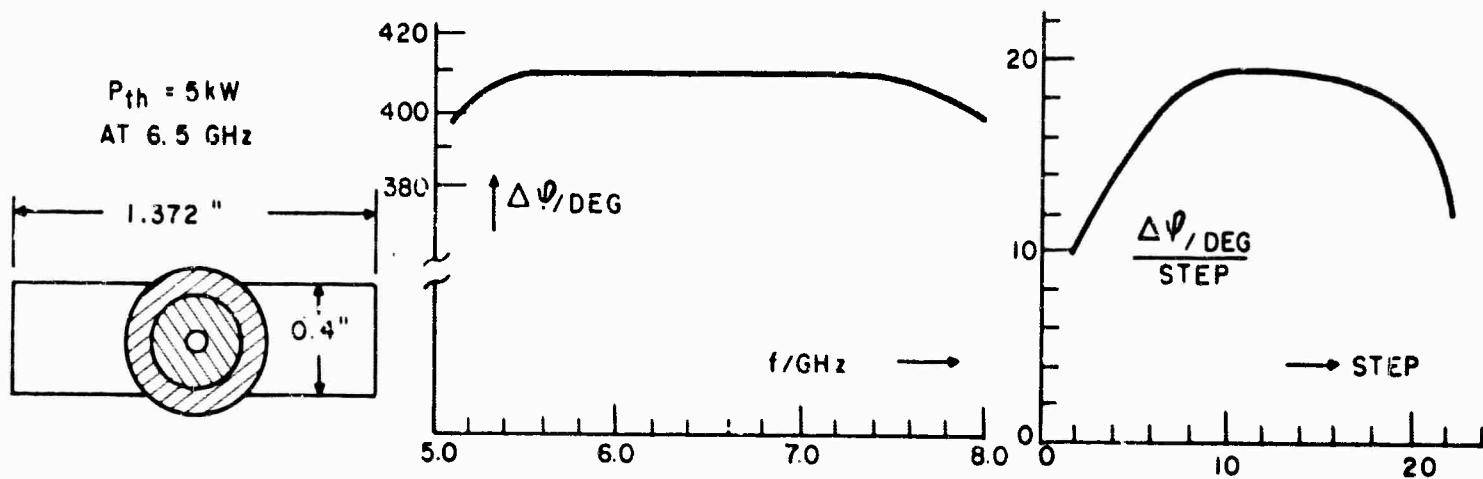
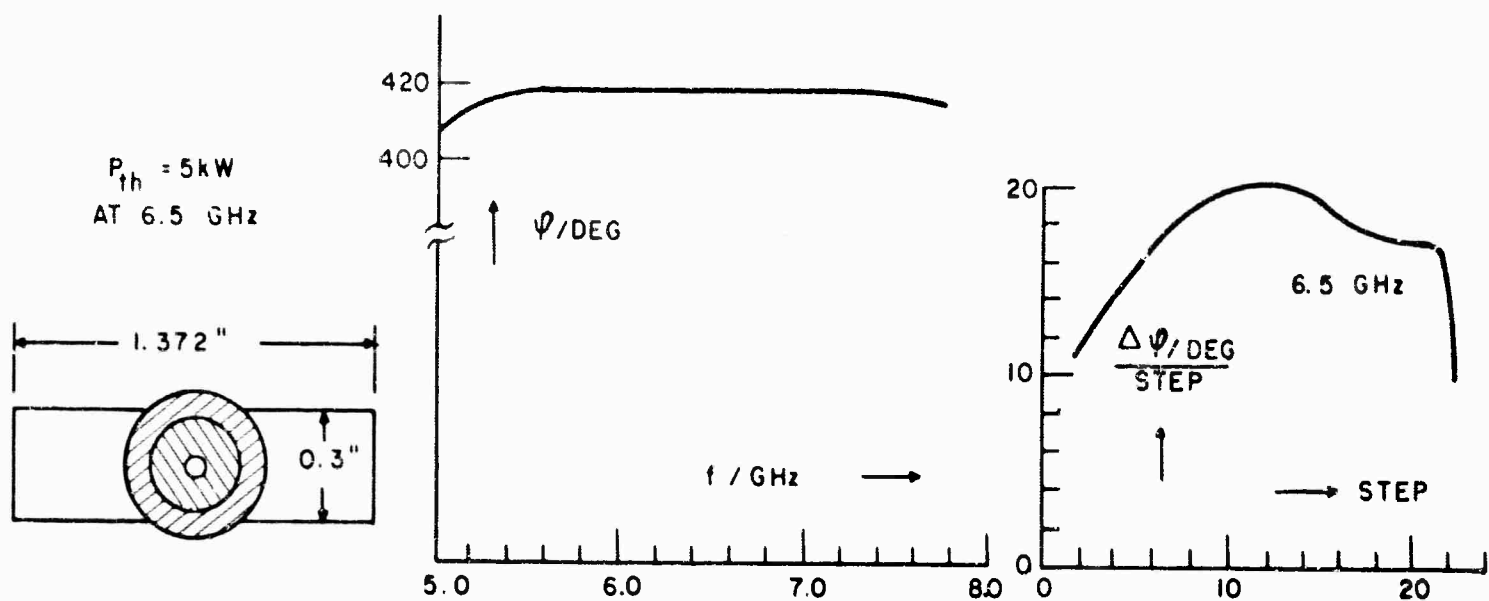
FIGURE 20



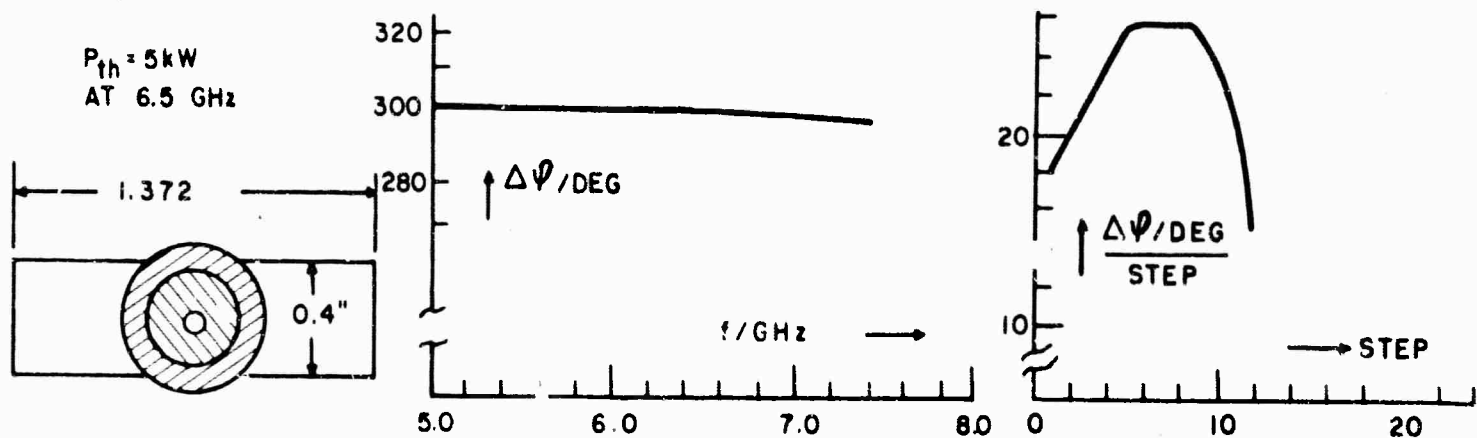
FERRITE : TTI-1400  $D_o = 0.500"$ ,  $D_i = 0.380"$ ,  $l = 5.0"$ , CORE: E = 11

### WAVEGUIDE TYPE PHASE SHIFTER

FIGURE 21



FERRITE TTI-1400,  $D_o = 0.437"$ ,  $D_i = 0.200"$ ,  $l = 5.0"$ ,  $\epsilon = 11$



FERRITE TTI-1400,  $D_o = 0.437"$ ,  $D_i = 0.300"$ ,  $l = 5.0"$ , CORE:  $\epsilon = 11$

WAVE GUIDE TYPE PHASE SHIFTER

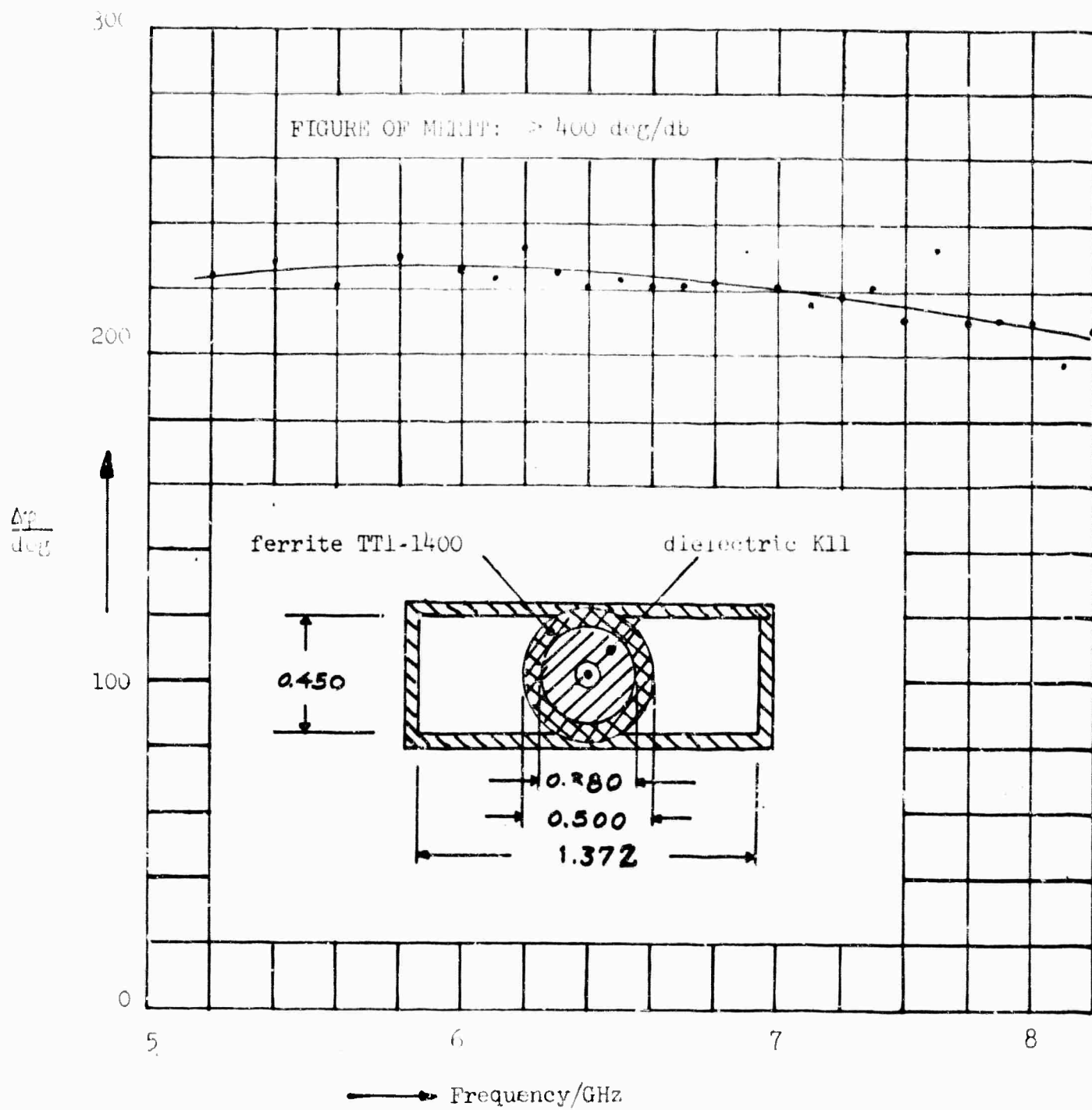
FIGURE 22

influence on the phase shift. Increasing the wall thickness by 80% increased the phase shift by 40%. The threshold power was independent of guide height and ferrite wall thickness.

The flattest phase shift versus transferred flux characteristic at 6.5 GHz was achieved with a filling factor of about 36%, as shown in the center row of Figure 21. The performance of this phase shifter is illustrated in more detail in Figures 23, 24, and 25. The phase shift/flux curve (Fig. 24) indicates that a "one driver bit" phase shifter could be built, if enough time is available to step up the BH-loop. The driver bit would be charged and discharged to the phase shifter as often as necessary to achieve a certain phase shift. The high power properties shown in Figure 25 are typical for all the phase shifters investigated. The ferrite was driven between its maximum remanent states. The differential phase shift stayed constant up to a power 1 to 3 dB above the threshold power. Then it decreased. This decrease is mainly due to the change in insertion length in one state of magnetization, the length in the other state stays approximately constant.

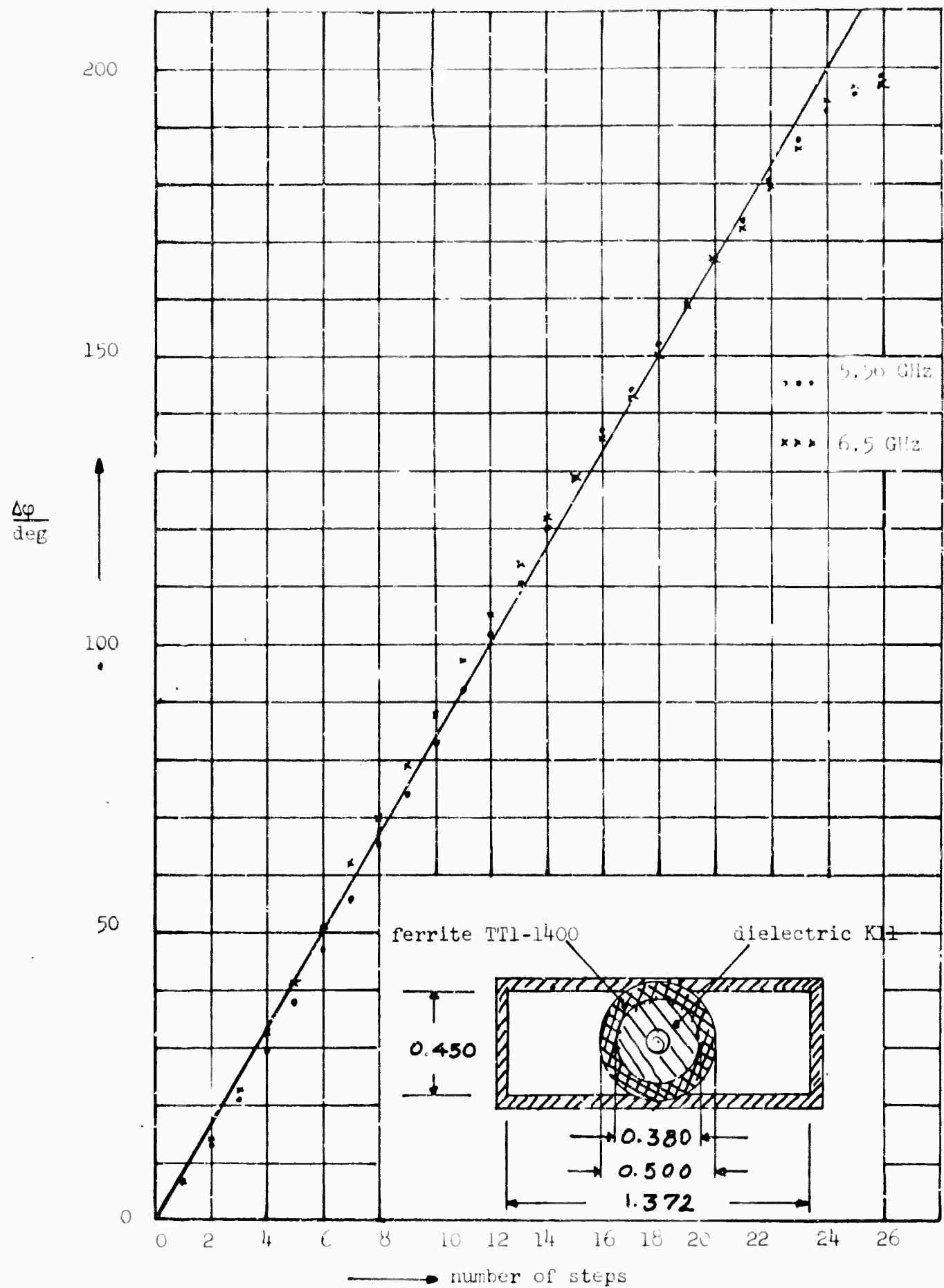
If the ferrite warms up even without non-linear losses, the insertion length and differential phase may vary even then. Methods to prevent this change are a) the use of temperature compensated garnets, b) an effective cooling method, and c) to drive the ferrite between its maximum negative remanent magnetization and a positive magnetization smaller than the maximum. This last method has been demonstrated first by Hair.<sup>(1)</sup> It is inherent in the flux transfer controlled phase shifters investigated here.

While this report was being written a final test was made with four inches of TT1-900 ferrite ( $M_s \approx 7 \mu\text{Vs}/\text{cm}^2$ ). The ferrite was selected because it was found to possess a very high threshold field strength (see Section III). The maximum differential phase shift was  $104^\circ$ . The bandwidth and flux control of phase shift were comparable to the result achieved with TT1-1400 with the same device and ferrite dimensions. The onset of non-linear losses occurred at above 100 KW. At 150 KW the non-linear losses were  $< 0.5$  dB.



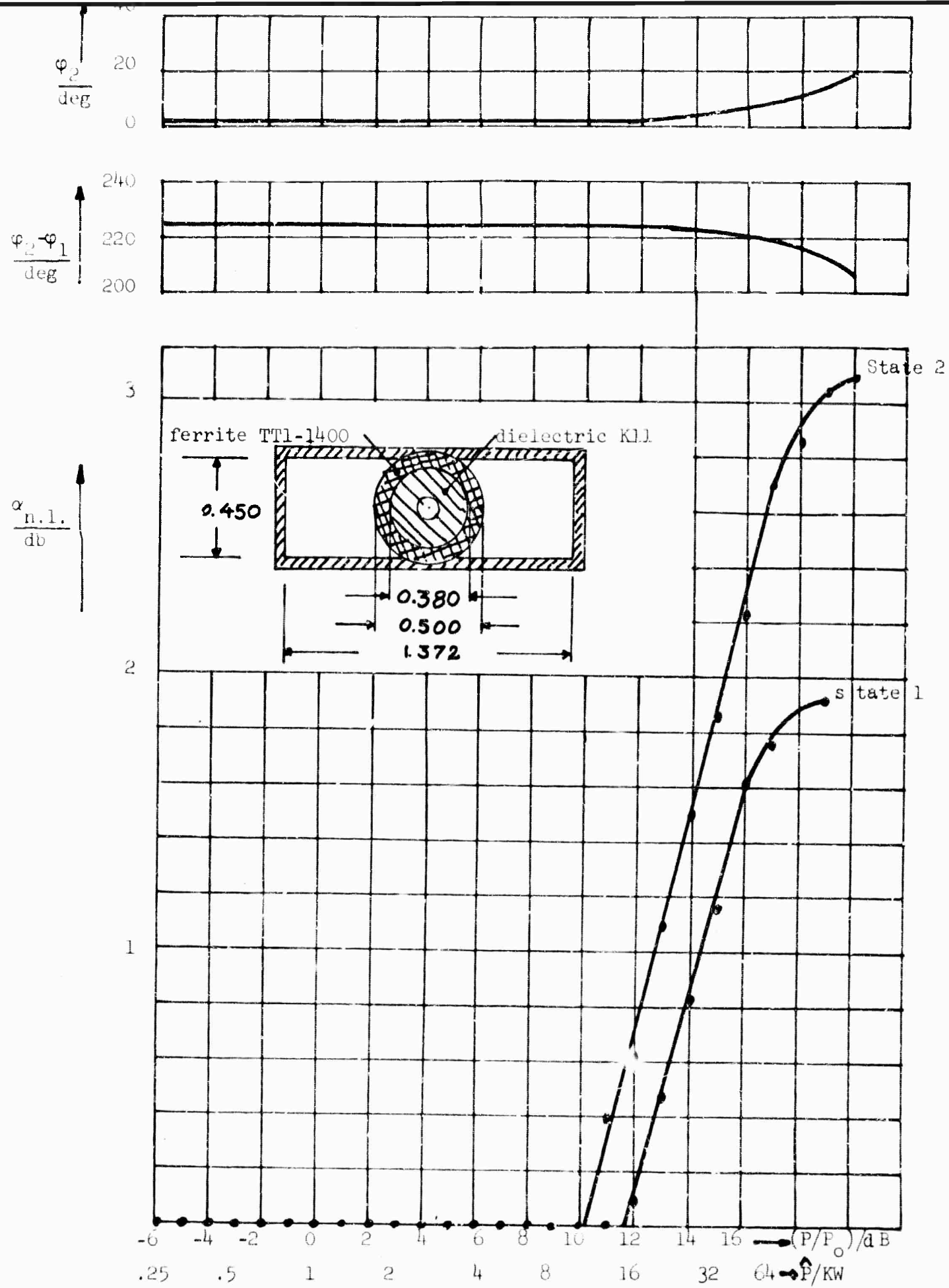
Differential Phase Shift Between States of Maximum Remanent Magnetization

FIGURE 23



Differential Phase Shift when "Stepping Up" a B-H Loop

FIGURE 24



High Power Performance of Differential Phase Shifter at 6.5 GHz

FIGURE 25

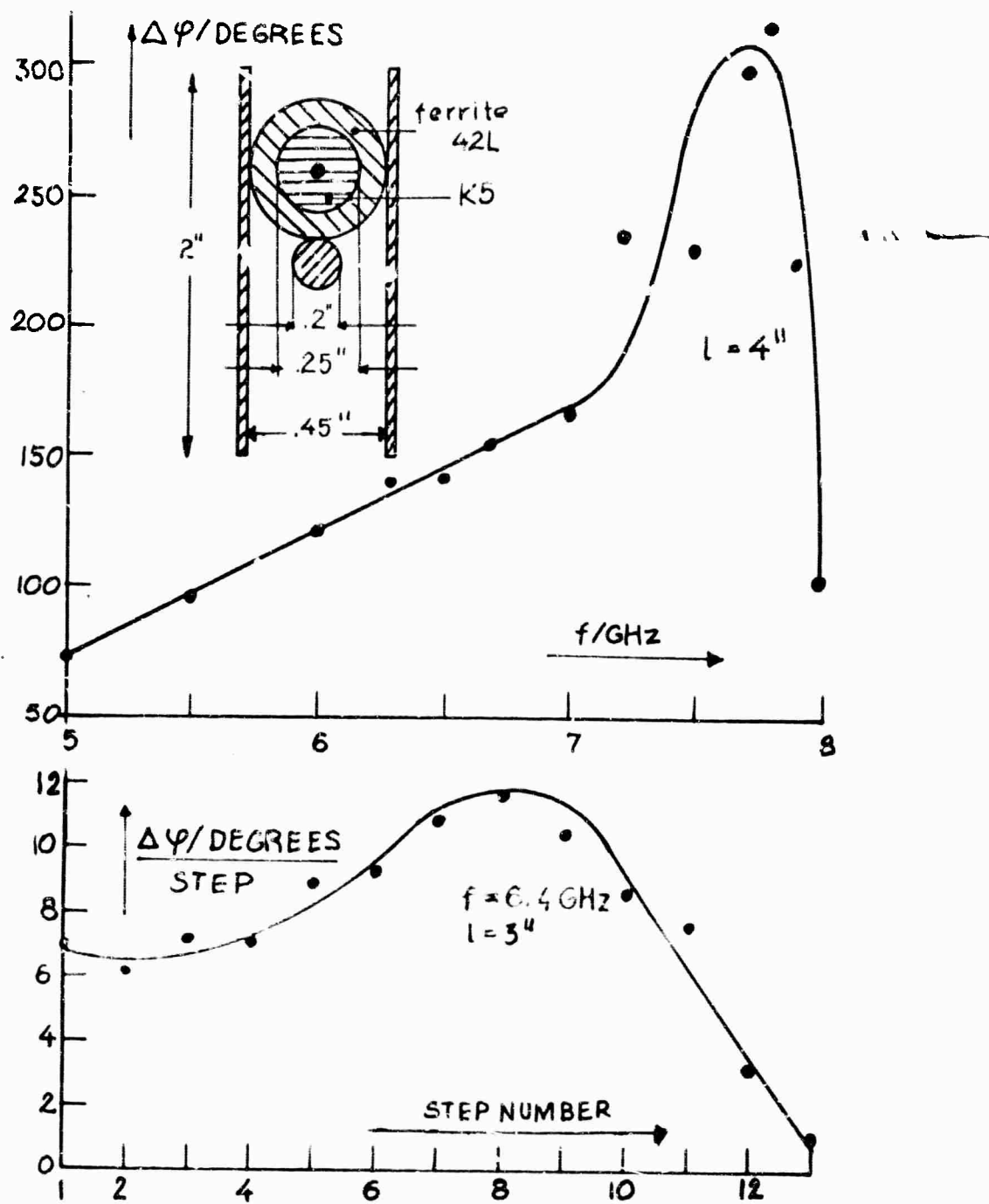
## V. 2 Slab Line Type Phase Shifters

A typical differential phase shift versus frequency characteristic of a slabline type phase shifter is shown in Figure 26. The phase shift increases gradually with frequency below 7 GHz, shows a high peak at 7.5 GHz and then declines suddenly. The phase shift per step of this phase shifter (with a reduced ferrite length) at 6.4 GHz shows a peak at the eighth step.

With some imagination one may use a rectangular waveguide with dielectric slabs parallel to the narrow walls as a model to analyze the fields in the slab line. This is done in Figure 27. The  $TE_{10}$  mode for the centrally loaded guide is the first to appear. With increasing frequency the plane of circular polarization moves from the center of the dielectric towards its outside. In the slabline this is the region of magnetization. A "positive" phase shift increasing with frequency is to be expected. A maximum differential phase shift per step will occur when the plane with circular polarization is magnetized. At high frequencies this is the boundary plane between dielectric and the empty part of the guide, corresponding to the outer ferrite diameter. At these frequencies, however, a  $TE_{01}$  mode corresponding to the side wall loaded rectangular guide is possible and at little higher frequencies the  $TE_{02}$  mode for either the centrally or sidewall loaded guide may also exist. Both these modes produce close to their respective cutoff-frequencies "negative" phase shift. At higher frequencies negatively as well as positively circularly polarized fields exist in the dielectric. The slabline type phase shifter responds in this frequency range with a rather erratic phase shift versus frequency characteristic. The phase shifter should, therefore, work in the frequency range below the cutoff-frequency corresponding to the sidewall loaded guide.

Figures 28 and 29 show that all the slab line type phase shifters with parallel side walls have the frequency characteristic described above. Increasing, in Figure 29, the wall thickness by 80% increased the phase shift by 50% at 6 GHz. The threshold power doubled from 5 to 10 KW.





Differential Phase Shift Characteristics of Slabline  
Type Phase Shifter

FIGURE 26

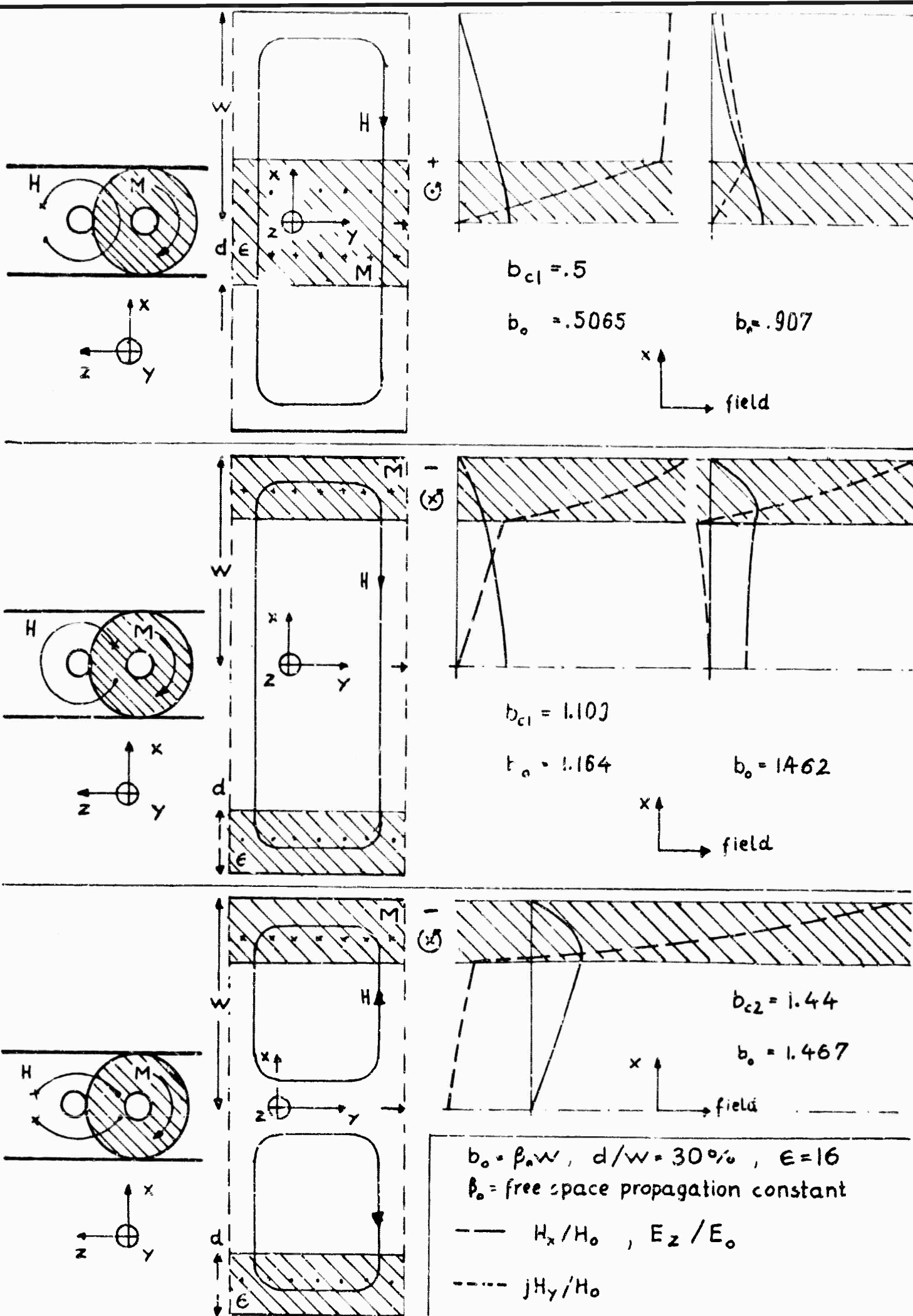
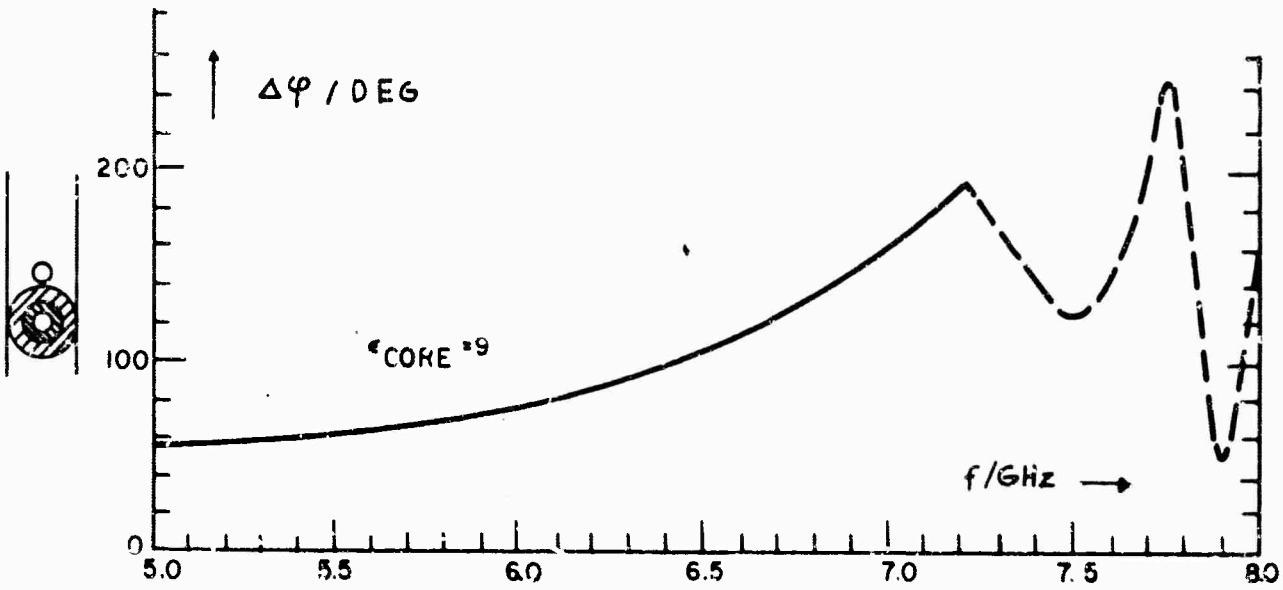
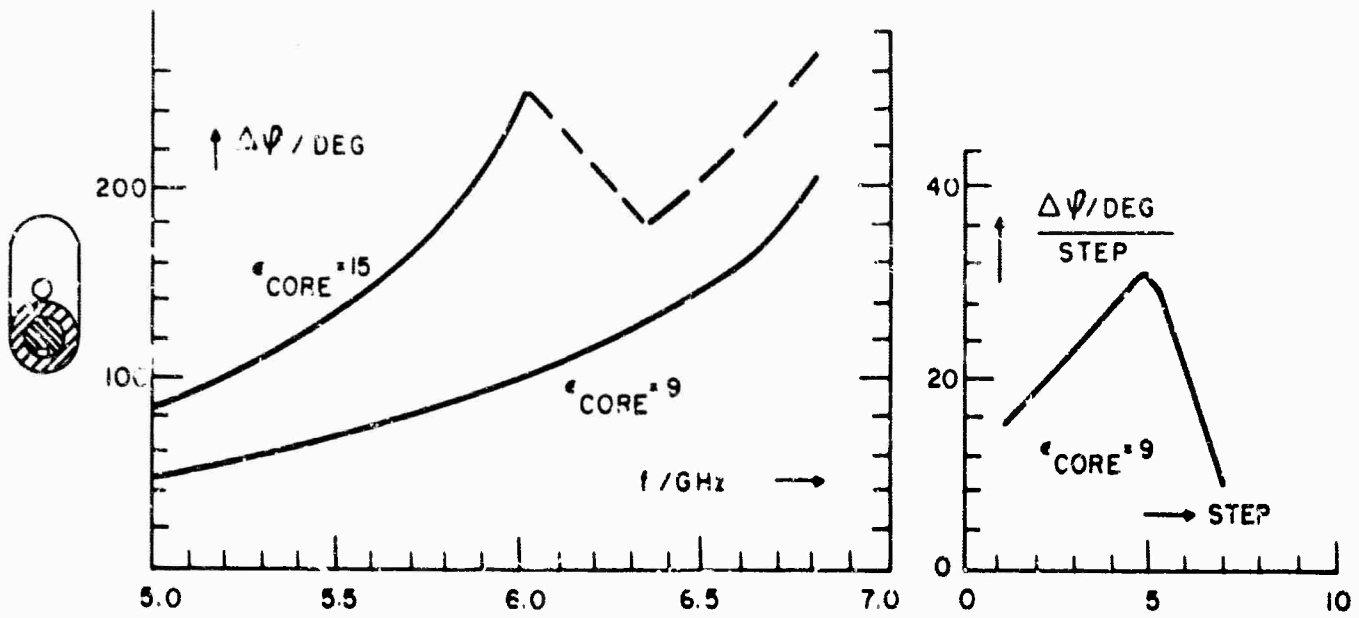
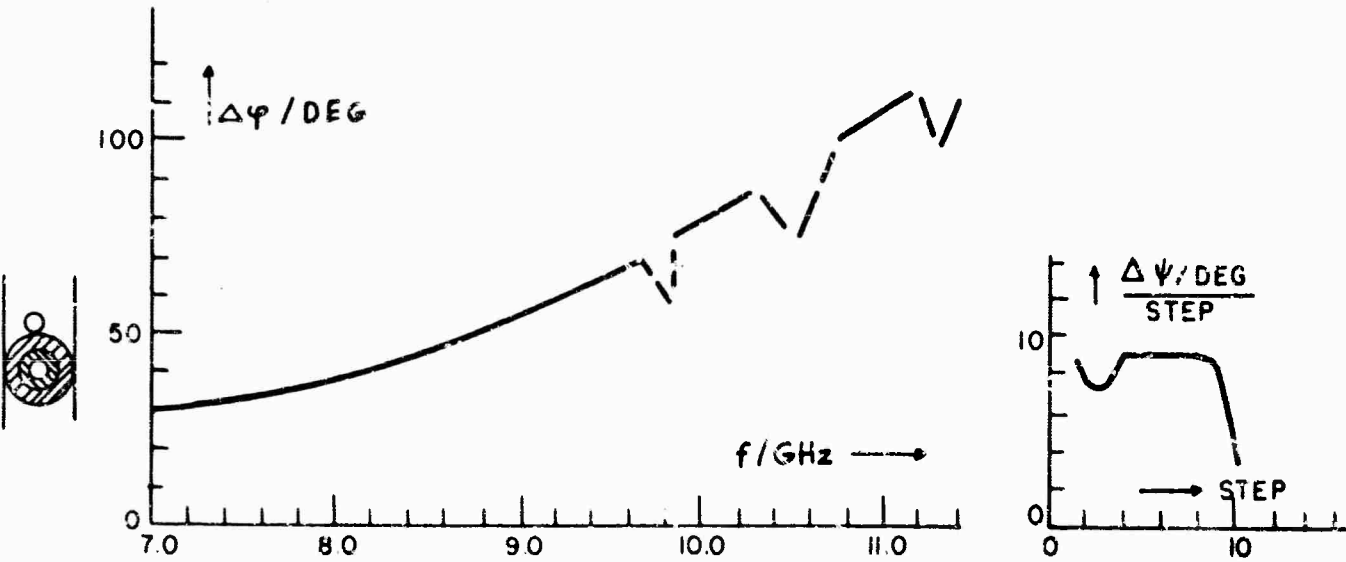


FIGURE 27. Waveguide Model for TE-Modes in Slablines



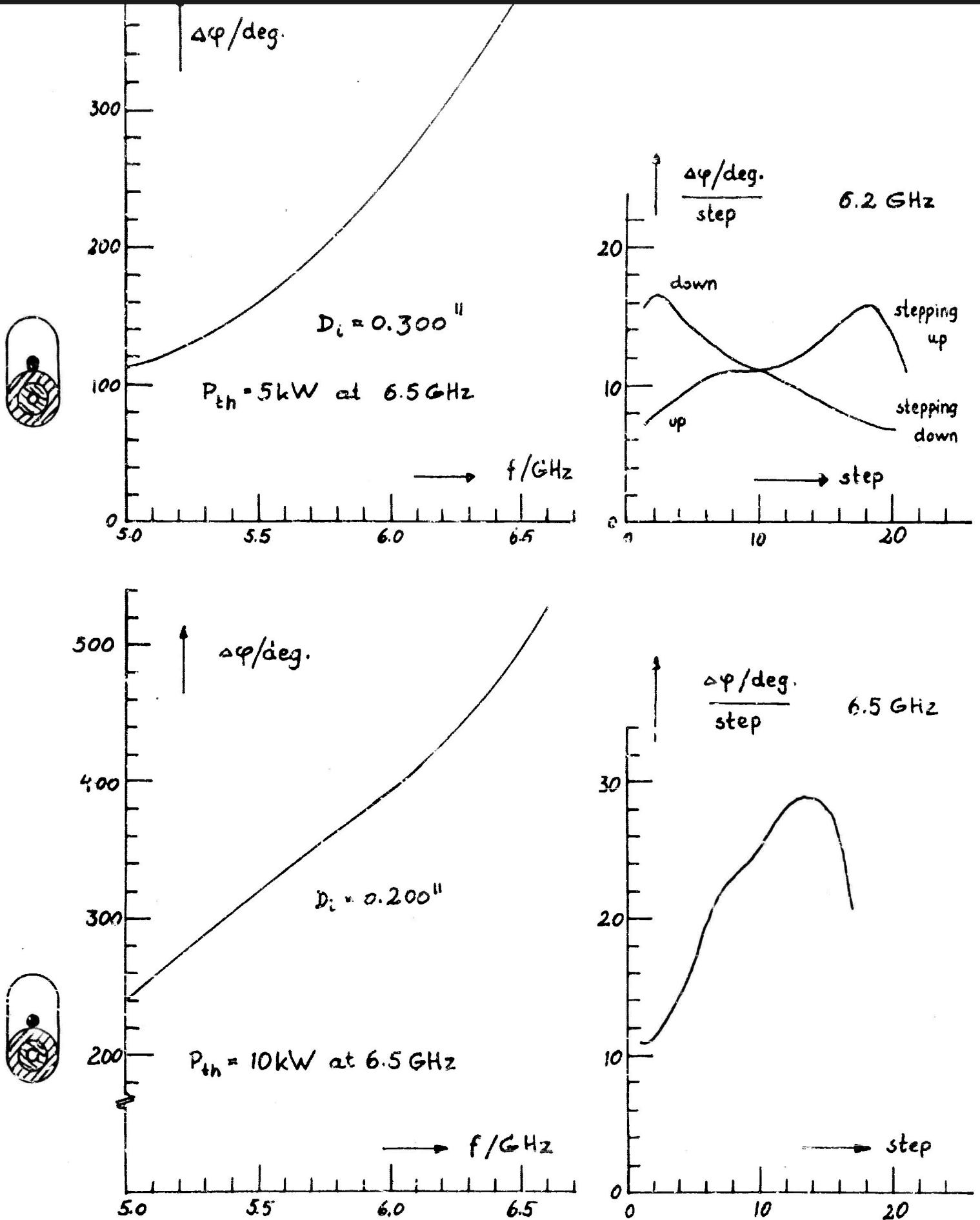
FERRITE TTi-103,  $D_o = 0.437"$ ,  $D_i = 0.29"$ ,  $l = 2.0"$



FERRITE TTi-103,  $D_o = 0.250"$ ,  $D_i = 0.150"$ ,  $l = 3.0"$ ,  $\epsilon_{\text{CORE}} = 9$ .

SLABLINE TYPE PHASE SHIFTER

FIGURE 28



Ferrite: TT1-1400,  $D_o = 0.437''$ ,  $L = 5.0''$ , Core:  $\epsilon = 10$

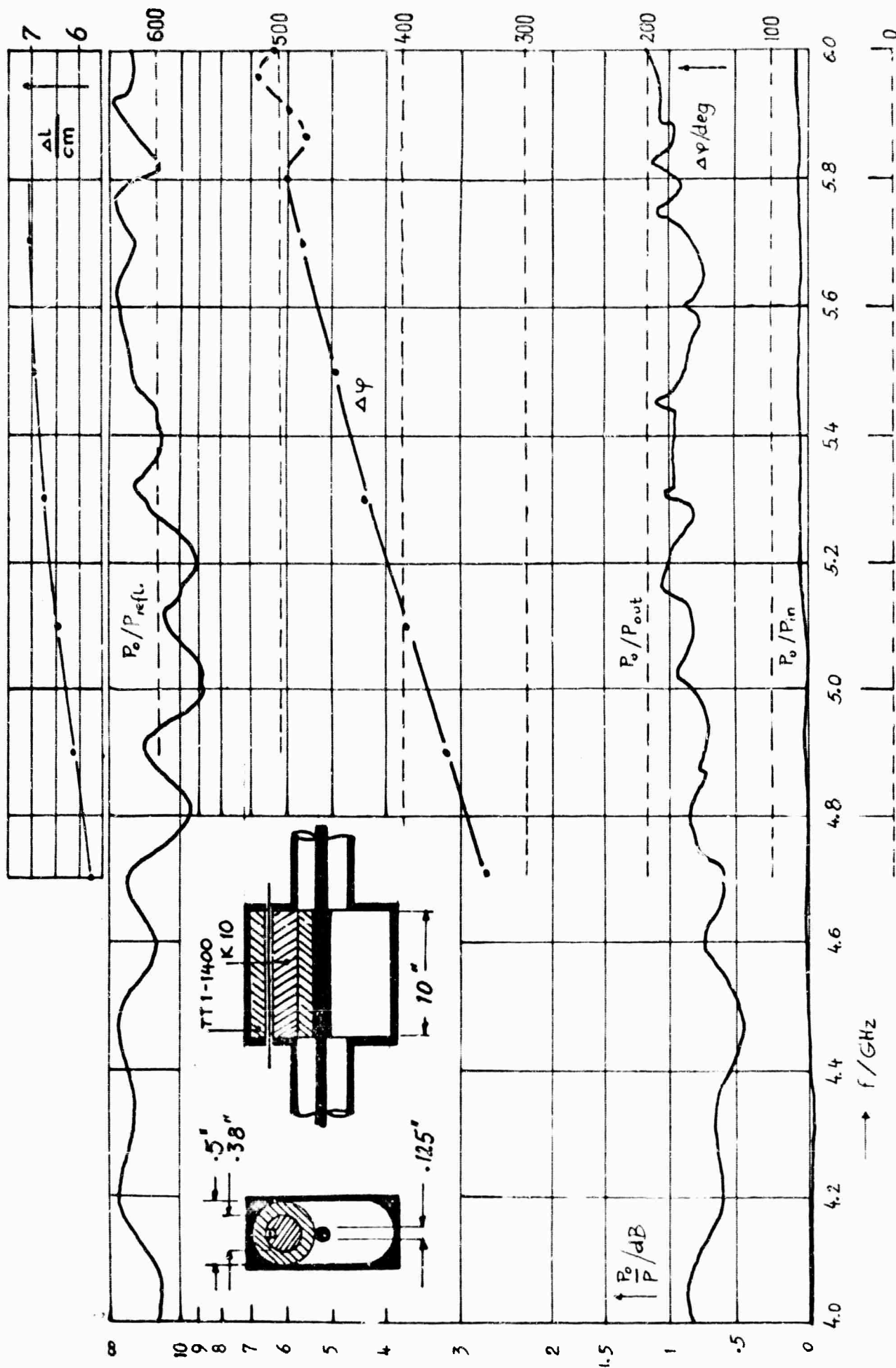
FIGURE 29

The lower value achieved with the thinner wall is somewhat doubtful because moding problems occur at the test frequency. In some other slabline type phase shifters with ferrites with  $D_6 = .5"$ ,  $D_i = .3"$  the threshold powers achieved at 6.5 GHz were

- 6 KW peak with TT1-105 in a shielded slabline
- 9 KW peak with TT1-105 in an open slabline
- 15 KW peak with TT1-1400 in an open slabline

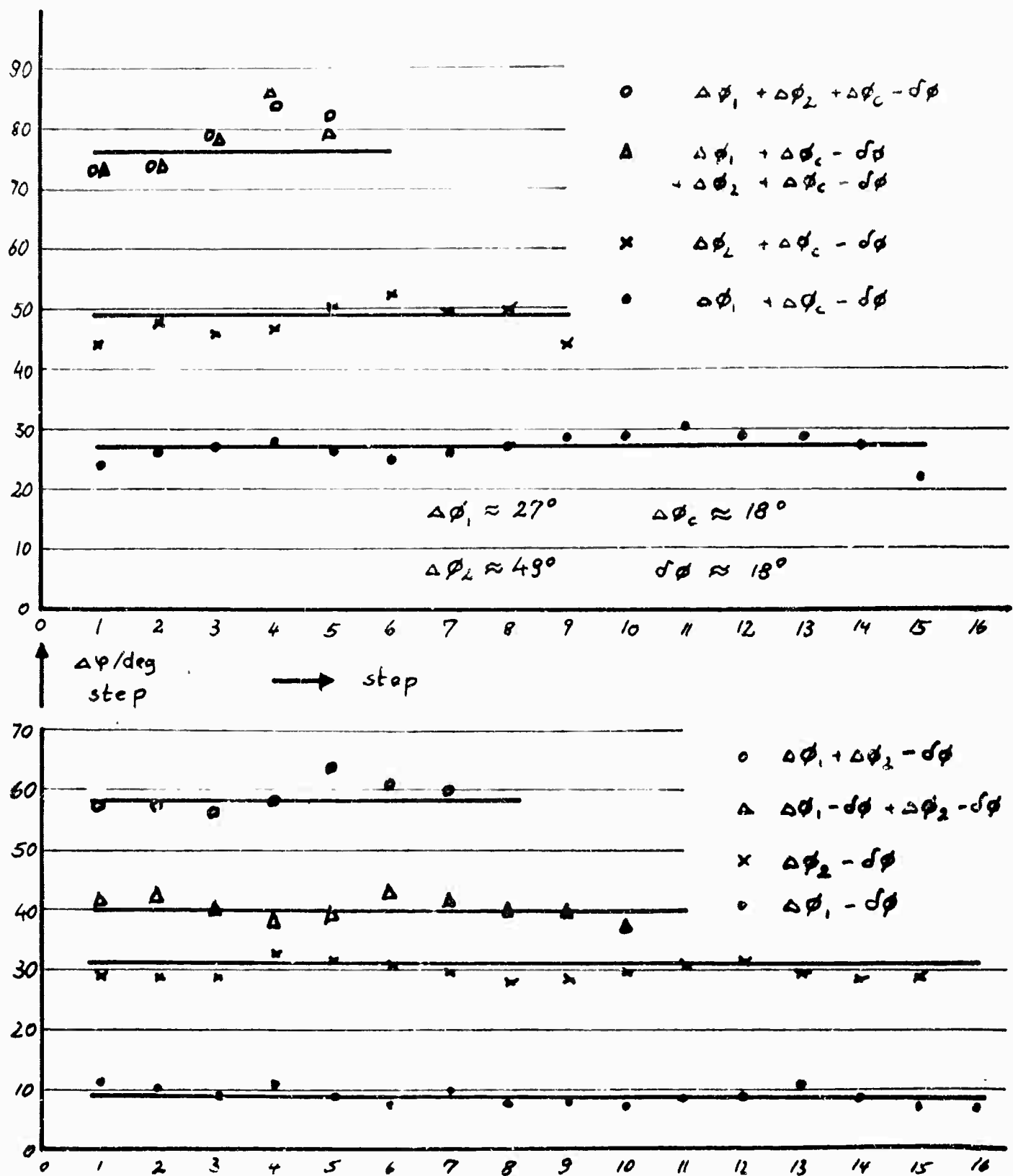
Changing the ferrite and device dimensions led to the phase shifter of Figure 30. The figure of merit is above  $400^\circ/\text{dB}$ . Half the ferrite is in contact with the outer conductor. This facilitates heat transfer problems at high average and peak powers. Figure 31 shows the phase shift per step for this device for two different bit sizes and their combinations (switched one after the other and both together) with and without a correction bit. Though the device performance is not yet ideal, the influence of the correction bit is striking. At 5.7 GHz four equal phase shift bits and a correction bit have been added with a transfer circuit as in Figure 8. The results are:

<u>Bit</u>	<u><math>\Delta\phi/\text{deg}(\text{measured})</math></u>	<u><math>\Delta\phi/\text{deg}(\text{nominal})</math></u>
1	76	77
2	77	
3	76	
4	76	
1 + 2	151	154
1 + 3	151	
1 + 4	151	
2 + 3	152	
2 + 4	152	
3 + 4	151	
1 + 2 + 3	233	231
1 + 2 + 4	233	
1 + 3 + 4	232	
2 + 3 + 4	233	
1 + 2 + 3 + 4	314	308



Overall Performance of Shielded Slabline Type Phase Shifter

FIGURE 30



Adding Phase Shift Bits by Controlled Flux Transfer ( $f = 5.6 \text{ GHz}$ )

FIGURE 31

Figure 32 shows the results achieved with four different bit sizes and the correction bit. The bits are somewhat too small for a  $360^\circ$  phase shifter. The best approximation is given by the  $350^\circ$  line. With respect to this line the error was within  $\pm 6^\circ$  and within  $\pm 3\%$ , whatever is smaller. Better results have been achieved but were not reproducible. For good results the switching wire has to be perfectly parallel to the center conductor and center conductor and switching wire have to be perfectly centered. This phase shifter could not be tested at high power, because the available source works at 6.5 GHz, i.e. beyond the phase shifter range.

To enlarge the bandwidth of the slab line type phase shifter a thin slab was inserted into the slab line as shown in Figure 33 to suppress the mode corresponding to the  $TE_{01}$  mode in the sidewall loaded rectangular guide. The frequency range has well been extended to higher frequencies. The small remaining erratic phase shift response is ascribed to imperfect adjustment of the switching wire and to the mode corresponding to the  $TE_{02}$  mode in the rectangular dielectrically loaded guide. Non-linear losses occurred at about 2.5 KW and voltage breakdown at about 35 KW. The figure of merit is  $> 300$  deg/dB.



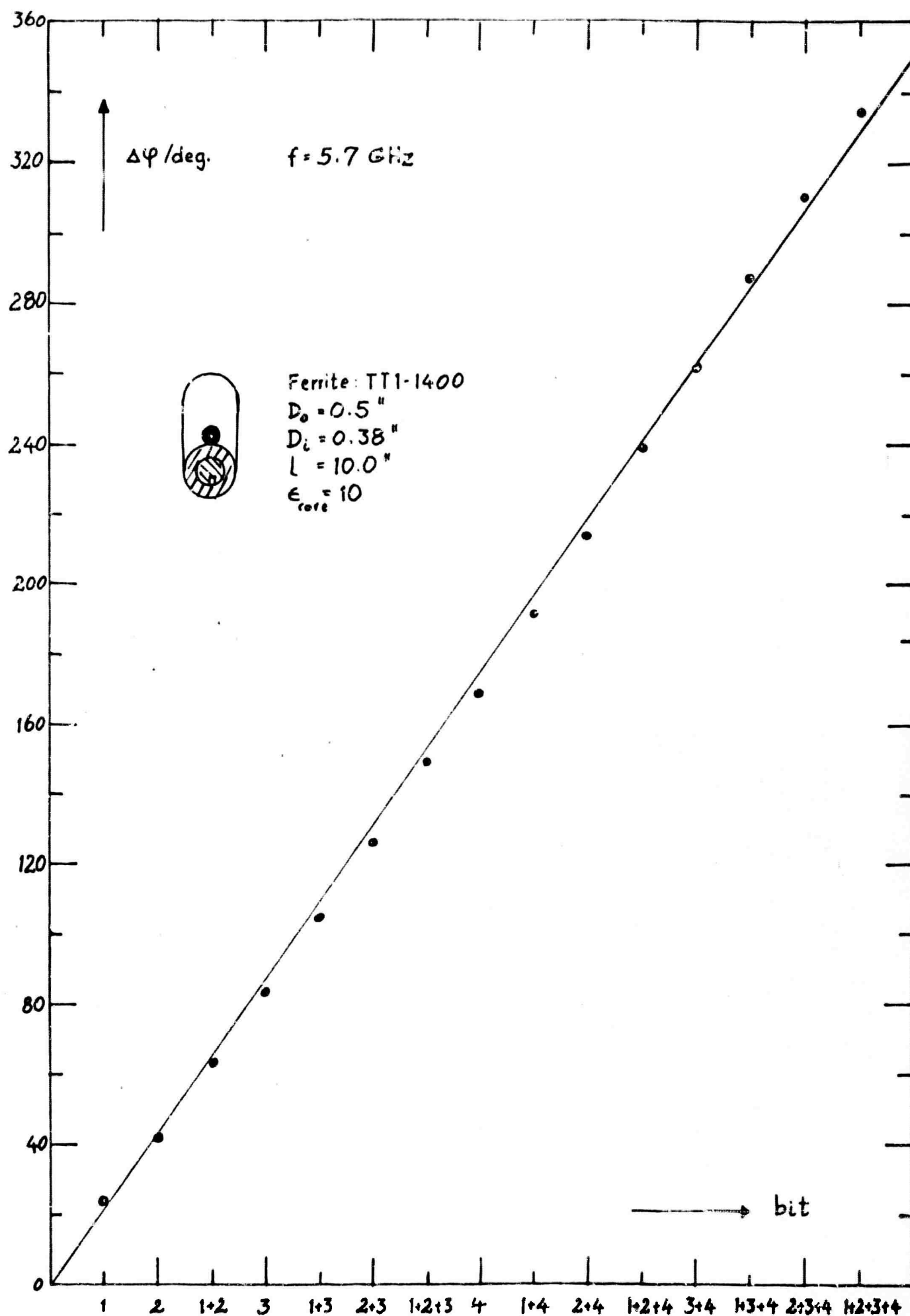


FIGURE 32. Slabline Type Phase Shifter with Flux Controlled Phase Bits

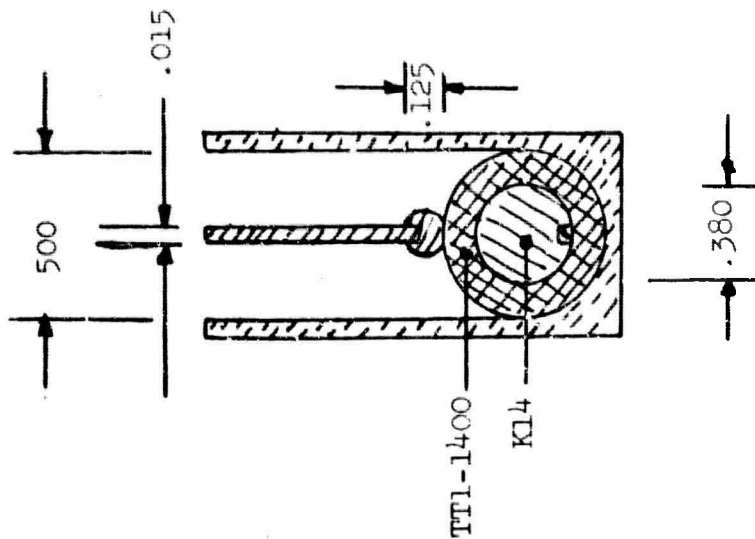
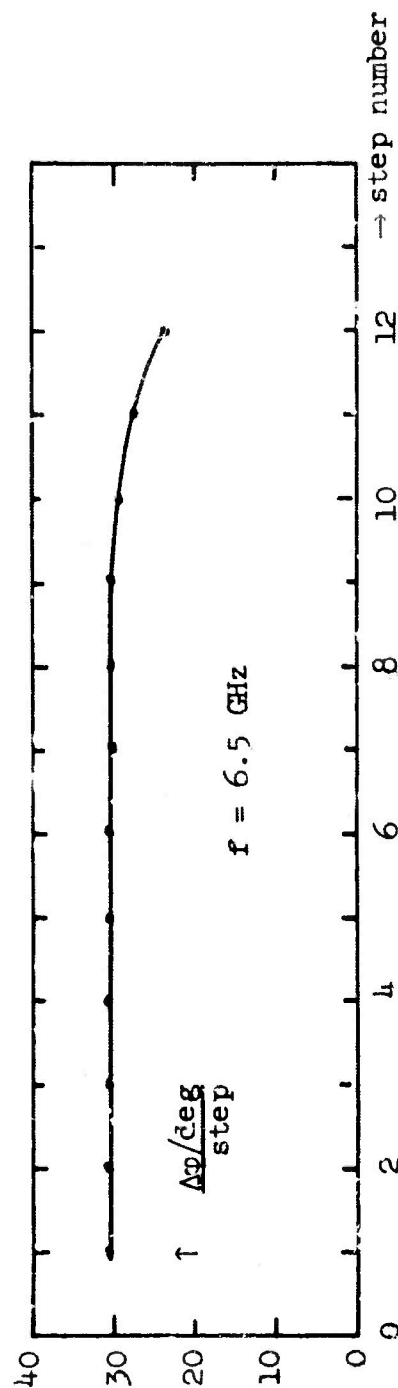
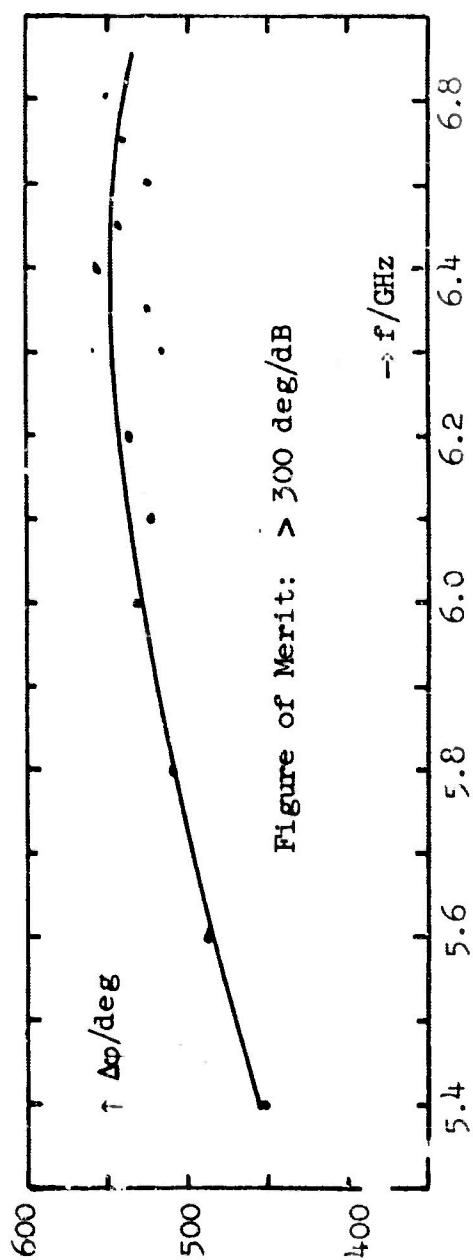


FIGURE 33. Stabline Type Phase Shifter with Mode Suppressor

The differential line lengths is

$$\Delta L = (\Delta\phi/\text{rad}) \cdot (\lambda/2\pi)$$

$$\Delta L/\text{cm} = (1/12) \cdot (\Delta\phi/\text{deg})/(f/\text{GHz})$$

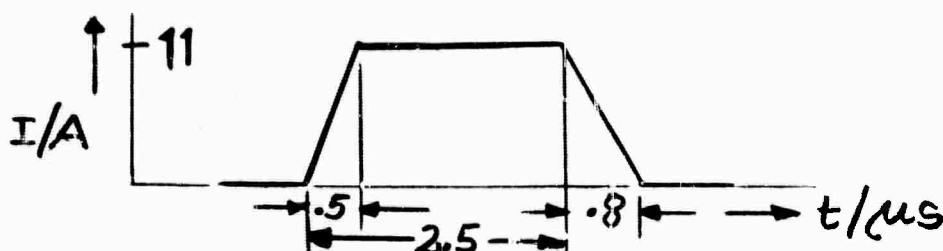
Phase shifters with positive slope  $[\partial(\Delta\phi)/\partial f]$  may thus be useful as delay lines with frequency independent differential line length. All slabline type phase shifters investigated have this positive slope. While the differential phase shift of the phase shifter in Figure 30 increases by 48% between 4.7 and 5.8 GHz, the differential line length increases only by 15%. The waveguide type phase shifter exhibits no slope with a filling factor of about 1/3, positive slope requires a reduced filling factor. A very broad band with constant phase shift has been achieved with this phase shifter. For low power applications the small size of the slab line type phase shifter is an advantage, for high power application voltage breakdown may occur before the ferrite goes unstable because of the small dimensions.

How well a ferrite is used in a phase shifter at a certain frequency may be described in terms of phase shift per flux. Typical for the devices investigated is  $10^\circ/\mu\text{Vs}$ . This value may be constant over a broad band. In helical phase shifters one may obtain  $30^\circ/\mu\text{Vs}$  in the narrow band with maximum phase shift.

Other workers,<sup>(7)</sup> using rectangular ferrite cylinders in waveguide type phase shifters, found that rounding the edges of the ferrite cylinders increased the differential phase shift and the figure of merit. Under the present contract the "rounding effect" was carried to the extreme.

## VI. DRIVING CIRCUITS

Figure 34 shows the basic circuit used in the drivers. It produces an eleven ampere current pulse 2.5  $\mu\text{sec.}$  long with 0.5  $\mu\text{sec.}$  rise time and 0.8  $\mu\text{sec.}$  fall time as sketched below. The power supply voltage was made



as large as possible under the constraints of the 2N1908 transistor.

Initially the reset driver was arranged as in Figure 1. This arrangement did not work well because one of the two parallel inductivities (microwave and driver core) was saturated first and then shorted the other one, so that the other core was not completely reset. Therefore, an arrangement, as indicated in Figure 8, was chosen for the reset circuit. A block diagram for the complete driver circuit is given in Figure 35.

The phase shifter of Figure 30 has a ferrite volume  $V \approx 14 \text{ cm}^3$ , a coercive field strength  $H_c \approx 1 \text{ A/cm}$ , and a saturation magnetization  $M_s \approx 10 \text{ } \mu\text{Vs/cm}^2$ . The required energy per switching cycle is  $E_{\text{requ}} \approx 4H_c M_s V \approx 560 \text{ } \mu\text{joules}$ . The drivers have an efficiency of at least 65%. Thus the total switching energy per cycle is 840  $\mu\text{joules}$  or less. It is believed that for all the phase shifters investigated one needs 0.8 to 1.5  $\mu\text{joules/degree}$ .

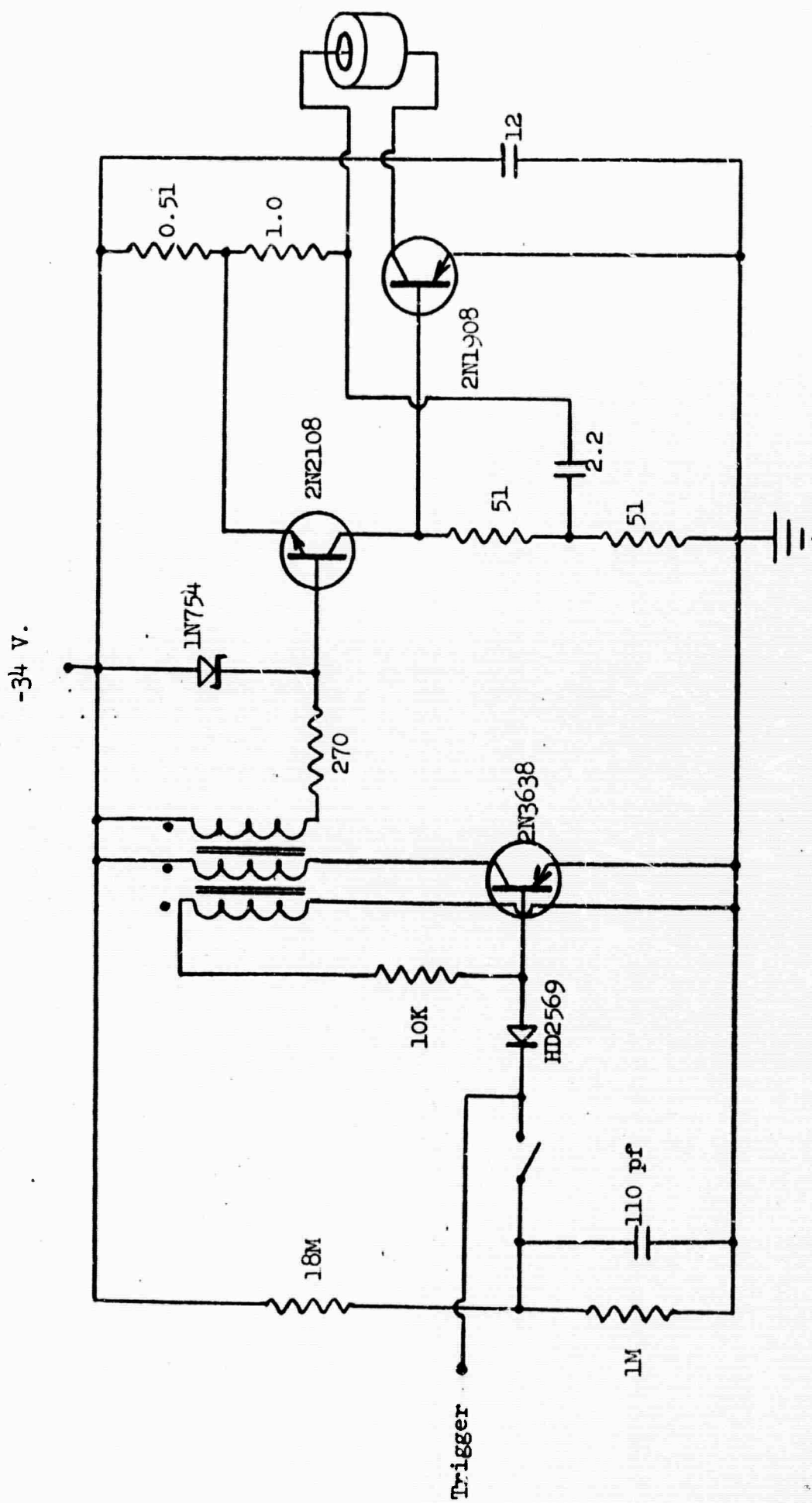
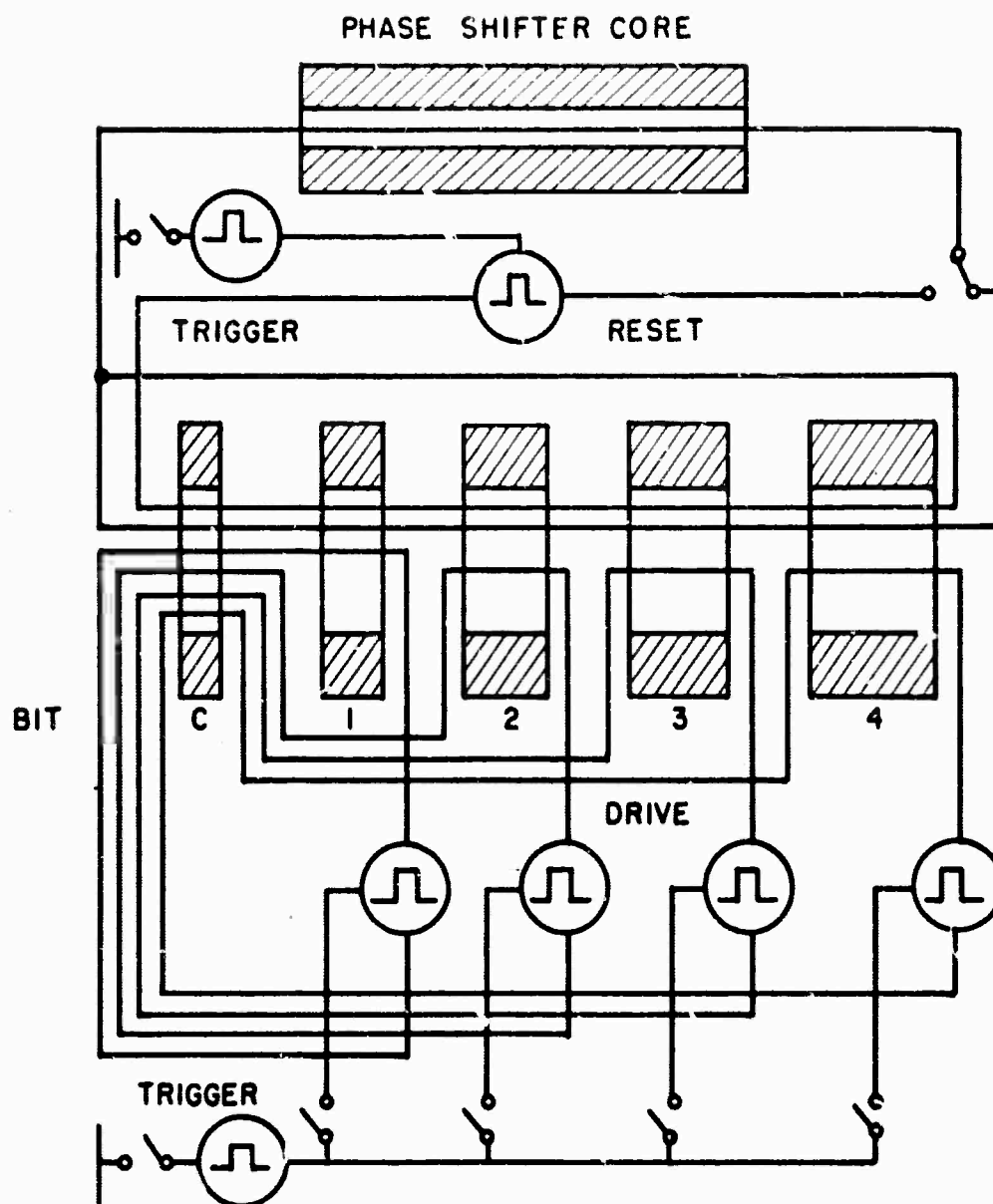


Figure 34. Pulse Driver Circuit



EXPERIMENTAL FOUR BIT DRIVER CIRCUIT

FIGURE 35

## ACKNOWLEDGEMENT

The assistance of Mr. G. T. Roome, who designed and tested the first slabline type phase shifter, and of Mr. E. J. McKinney, who wrote Part II of this report, is highly appreciated. Mr. G. E. Claflin took most of the flux transfer data and designed, together with Mr. T. L. Sly, the driver circuits. Mssrs. E. E. Des Jardins and D. L. Kortz took the bulk of the experimental data. As always, the secretaries, Mrs. I. Maki and Miss L. Jessmore did an excellent job in preparing this report. It is a pleasure to thank all of them.

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## PART II. NONLINEAR LOSSES IN FERRITES

### ABSTRACT

First and second order solutions for the critical field at which the magnetization in a domain becomes unstable are derived. The domain is assumed to be immersed in a larger sample of ferrite which has a spacial average magnetization equal to the remnant value. It is assumed that the magnetization in the domain of interest is uniform and in the direction of the applied field since this case yields the smallest critical field. The first order solution corresponds to the zero order spin mode for the domain (all spins precessing in unison). The second order solution includes spin waves which exist as traveling waves within the domain. A third order solution would include interactions between spin wave modes, particularly the zero order mode and higher order modes. The spin wave which is excited to unstability by the driving field is shown to have half the frequency of the driving field. It is shown that the critical field increases as the saturation magnetization of the material decreases.

<u>Section</u>	<u>Page</u>
I. INTRODUCTION .....	1
II. DEFINITION OF TERMS AND EQUATIONS OF MOTION .....	2
III. FIRST ORDER SOLUTION .....	8
IV. SECOND ORDER SOLUTION .....	12
V. DISCUSSION OF RESULTS AND HIGHER ORDER TERMS .....	19

## NON-LINEAR LOSSES IN FERRITES

### I. INTRODUCTION

The determination of the rf field strength at which nonlinear effects commence in a ferrite has been the object of several investigations. Suhl<sup>(1)</sup> developed the general theoretical approach and applied it to the configuration in which an rf driving field is applied perpendicular to a dc magnetic biasing field. Schlömann and his associates in various publications<sup>(2)</sup> have presented the results for a rf driving field which is parallel to and small compared to the dc magnetic biasing field. This investigation considered the special case of a ferrite toroid biased to saturation by means of a pulse of current in a wire coincident with the axis of the toroid. During the time interval of the current pulse, the magnetization throughout the sample is circumferential and has the saturation magnitude,  $M_s$ . When the current is removed, it is assumed that domains are formed. In each domain the magnetization is unidirectional, but not necessarily circumferential, and has the saturation magnitude. The magnetization from domain to domain is not unidirectional however, and the average circumferential magnetization throughout the toroid is reduced to the remnant magnitude,  $M_R$ .

After the biasing current is removed, a rf driving field  $h_1 \cos \omega_1 t$  is applied in the circumferential direction. The purpose of this analysis is to determine the minimum value of  $h_1$  which will cause instability in the toroid. The critical field will be a minimum in a domain in which the dc part of the magnetization is in the same direction as the applied field, so that only this case will be considered. It is assumed that there exists at least one domain within the toroid for which this applies.

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(1) H. Suhl, Proc. IRE, 44, 1270 (1956).

(2) E. Schlömann and R. I. Joseph, Jour. Appl. Phys., 32, 1006 (1961) for example.

## II. DEFINITION OF TERMS AND EQUATIONS OF MOTION

Consider a domain in which the dc part of the magnetization is in the same direction as the applied rf field. Assume that the domain is small enough to neglect the curvature of the fields in the circumferential direction of the toroid and let the direction of the applied field be defined as the Z axis of a cartesian coordinate system.

The total magnetization within the domain of interest is assumed to be made up of several components defined as follows:

- $\vec{m}$  is the total magnetization vector at a point
- $\vec{M}$  is the dc component of the magnetization vector
- $\vec{m}$  is the spacial average of the ac part of the magnetization averaged throughout the sample
- $\delta\vec{m}$  is the local deviation of the magnetization vector from the spacial average.

Similar definitions will apply to the magnetic field vector terms. The magnetization and field vectors are then expressed in rationalized MKS units by

$$\vec{m} = \vec{M} + \vec{m} + \delta\vec{m} \quad (1)$$

$$\vec{H} = \vec{H} + \vec{h} + \delta\vec{h} = \frac{B}{\mu_0} - \vec{m} \quad (2)$$

The magnetization is assumed to have a dc component only in the z direction. After the initial biasing magnetic field is removed, it is assumed that no applied dc component of magnetic field exists. Let  $\vec{x}$ ,  $\vec{y}$ , and  $\vec{z}$  be unit vectors in the cartesian coordinate directions and Equation (1) can be written in component form as

$$\vec{m} = (m_x + \delta m_x) \vec{x} + (m_y + \delta m_y) \vec{y} + (M_z + m_z + \delta m_z) \vec{z} \quad (3)$$

The magnetic field inside the domain is made up of four terms, the applied rf field in the z direction, the demagnetizing field resulting from the average ac magnetization throughout the domain, the demagnetizing

field due to the magnetization in the toroid sample outside the domain of interest, and the demagnetizing field due to local deviations of the magnetization. The demagnetizing field due to the average ac magnetization in the domain can be expressed in terms of the demagnetizing factors  $N_x$ ,  $N_y$  and  $N_z$ . The local deviations of the magnetization can be expressed in terms of a spin wave spectrum which gives rise to a magnetic field which must satisfy Maxwell's equations. The resultant demagnetizing field can be expressed as<sup>(3)</sup>

$$\underline{\delta h} = - \sum_{\underline{k}} \frac{\underline{k}(\underline{\delta m} \cdot \underline{k})}{k^2} \quad (4)$$

where  $\underline{k}$  is the spin wave vector.

It will be assumed that the only important local deviations of the magnetization are those which are driven in some way by the applied rf field and that these are restricted, in any particular example, to a narrow region of wave vector  $\underline{k}$ . Therefore, Equation (4) can be replaced by the single term in  $\underline{k}$  which dominates. That is

$$\underline{\delta h} \approx - \frac{\underline{k}(\underline{\delta m} \cdot \underline{k})}{k^2} \quad (5)$$

The demagnetizing field due to the toroid magnetization outside the sample can be expressed by<sup>(4)</sup>

$$\underline{h}_d(\underline{r}) = \nabla \int \frac{(\nabla' \cdot \underline{M})}{|\underline{r} - \underline{r}'|} d^3 \underline{r}' \quad (6)$$

where  $\underline{h}_d(\underline{r})$  represents the demagnetizing field at the point  $\underline{r}$  inside the sample,  $\nabla$  represents the gradient operator applied at the point  $\underline{r}$  and  $\nabla' \cdot \underline{M}$  represents the divergence of  $\underline{M}$  at the point  $\underline{r}'$  outside the domain of

(3) Ronald F. Sopho, Theory and Application of Ferrites; Prentice-Hall, Inc., 1960, p. 232.

(4) Ernst Schlömann; Advances in Quantum Electronics; Columbia Univ. Press, 1961, p. 450.

interest. Since the magnetization is assumed to be uniform in any domain, Equation (6) would become a surface integral over all the domain boundaries. Since

$$\nabla \left( \frac{1}{|r-r'|} \right) = \frac{\underline{r}-\underline{r}'}{|\underline{r}-\underline{r}'|^3} \quad (7)$$

the integrand in Equation (6) will decrease as one over the square of the distance from the point at which the field is being calculated to the point at which the source of the field is located. Thus, only the magnetization in domains bordering the domain in which instability occurs are important in determining  $\underline{h}_d(\underline{r})$ . Obviously  $\underline{h}_d(\underline{r})$  will be a function of position within the domain but also of the location of the domain within the toroidal sample. Since the domain structure throughout the sample cannot be known accurately, it will only be possible to approximate  $\underline{h}_d(\underline{r})$ . The simplest approximation is to let it be zero. This is not a bad approximation for points well within the domain of interest (the domain in which instability occurs) because of the one over distance squared dependence mentioned above. A better approximation might be to assume that the magnetization in neighboring domains is nearly z directed and assume a uniform value of  $\underline{h}_d(\underline{r})$  in the z direction of the form  $-NM_R$ . For most domains this would be a good approximation since the presence of residual magnetization implies an average non zero z component of magnetization throughout the toroidal sample. There will be domains, however, in which the magnetization is not nearly z directed, and for these the approximation would be poor. In this analysis, it will be assumed that  $\underline{h}_d(\underline{r})$  is zero.

The magnetic field in the domain can then be expressed as

$$\begin{aligned} \underline{H} = & \left( -N_x m_x - \frac{k_x}{k^2} \underline{\delta m} \cdot \underline{k} \right) \underline{x} + \left( -N_y m_y - \frac{k_y}{k^2} \underline{\delta m} \cdot \underline{k} \right) \underline{y} \\ & + \left( h_1 \cos \omega_1 t - N_z M_z - N_z m_z - \frac{k_z}{k^2} \underline{\delta m} \cdot \underline{k} \right) \underline{z} \end{aligned} \quad (8)$$

The equation of motion for the total magnetization vector contains three terms of interest. These are: the Zeeman term which describes the coupling between the magnetic field and magnetization, the loss term and the exchange term which describes the mutual coupling between spins. Crystalline anisotropy terms are omitted because we are dealing with a polycrystalline material. The equation of motion is given as

$$\dot{\vec{m}} = -\gamma(\vec{h} \times \vec{m}) - \frac{\gamma}{M_s} a^2 H_E \vec{m} \times \nabla^2 \vec{m} + \text{loss term} \quad (9)$$

where  $\gamma$  is the gyromagnetic ratio,  $a^2$  the lattice constant and  $H_E$  is the effective exchange field. The loss term will be of the Landau-Lifshitz form given by Suhl.<sup>(5)</sup> The loss term is not stated explicitly in Equation (9) because the loss parameter  $\alpha$  will not necessarily be the same for  $m$  and for  $\delta m$ . Let  $\alpha_0$  represent the loss term associated with the uniform precession terms  $m$  and  $\alpha_k$  represent the spin wave loss term. Some care is necessary in the choice of sign applied to  $\alpha_0$  and  $\alpha_k$  to insure that the homogeneous solutions ( $h_1 = 0$ ) result in decaying waves. The following assumption is made regarding the order of magnitude of the magnetization

$M_z$  is a zero order quantity

$m_x$  and  $m_y$  are first order quantities

$m_z$ ,  $\delta m_x$  and  $\delta m_y$  are second order quantities

$\delta m_z$  is a third order quantity

Since it has been assumed that the domain is saturated, the magnetization terms can be related by

$$(m_x + \delta m_x)^2 + (m_y + \delta m_y)^2 + (M_z + m_z + \delta m_z)^2 = M_s^2 \quad (10)$$

Equating terms of equal order of magnitude in Equation (10) gives the approximations

$$\begin{aligned} M_z &\approx M_s \\ m_z &\approx -\frac{1}{2} \frac{m_x^2 + m_y^2}{M_s} \\ \delta m_z &\approx -\frac{1}{M_s} (m_x \delta m_x + m_y \delta m_y) \end{aligned} \quad (11)$$

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(5) H. Suhl, Proc IRE, 44, (1956) p. 1271.

Define  $\omega_s$  and  $\omega_m$  by

$$\omega_s \equiv \gamma h_1 \quad (12)$$

$$\omega_M \equiv \gamma M_s \approx \gamma (M_z + m_z) \approx \gamma' M_z \quad (13)$$

Let  $Z_x$  and  $Z_y$  be the x and y components of the Zeeman term in Equation (9). Then from Equations (3) and (8) and the order of magnitudes assumed,  $Z_x$  and  $Z_y$  can be written to third order as

$$\begin{aligned} Z_x = & - \left[ \omega_s \cos \omega_1 t - \omega_M (N_z - N_y) \right] m_y && \text{First Order} \\ & - \left[ \omega_M \frac{k_x k_y}{k^2} \right] \delta m_x - \left[ \omega_s \cos \omega_1 t - \omega_M N_z + \omega_M \frac{k_y^2}{k^2} \right] \delta m_y && \text{Second Order (14)} \\ & - \left[ \omega_M \frac{k_y k_z}{k^2} \right] \delta m_z + \left[ \gamma \frac{k_x k_z}{k^2} \right] \delta m_x m_y + \left[ \gamma \frac{k_y k_z}{k^2} \right] \delta m_y m_y && \text{Third Order} \end{aligned}$$

$$\begin{aligned} Z_y = & \left[ \omega_s \cos \omega_1 t - \omega_M (N_z - N_x) \right] m_x && \text{First Order} \\ & + \left[ \omega_M \frac{k_x k_y}{k^2} \right] m_y + \left[ \omega_s \cos \omega_1 t - \omega_M N_z + \omega_M \frac{k_y^2}{k^2} \right] \delta m_x && \text{Second Order (15)} \\ & + \left[ \omega_M \frac{k_x k_z}{k^2} \right] \delta m_z - \left[ \gamma \frac{k_x k_z}{k^2} \right] \delta m_x m_x - \left[ \gamma \frac{k_y k_z}{k^2} \right] \delta m_y m_y && \text{Third Order} \end{aligned}$$

Similarly the components  $L_x$  and  $L_y$  of the loss term and  $E_x$  and  $E_y$  of the exchange term can be written

$$L_x = \alpha_o \dot{m}_y - \alpha_k \dot{\delta m}_y + \frac{\alpha_o}{M_s} \dot{m}_z m_y \quad (16)$$

$$L_y \approx -\alpha_o \dot{m}_x + \alpha_k \dot{\delta m}_x - \frac{\alpha_o}{M_s} \dot{m}_z m_x \quad (17)$$

where the first term on the right is the first order term, the second the second order term and the third the third order term and

$$E_x = a^2 \omega_E \nabla^2 \delta m_y = - a^2 \omega_E k^2 \delta m_y \quad (18)$$

$$E_y = - a^2 \omega_E \nabla^2 \delta m_x = a^2 \omega_E k^2 \delta m_x \quad (19)$$

where

$$\omega_E \equiv \gamma H_E \quad (20)$$

The exchange terms  $E_x$  and  $E_y$  are both of second order. The exchange term contains no first or third order terms.



### III. FIRST ORDER SOLUTION

Collecting only the first order terms from Equations (14) through (19) gives

$$\begin{aligned}(1 + i\alpha_0)\dot{m}^+ + iAm^+ + iBm^- &= i\omega_s (\cos\omega_1 t)m^+ \\ (1 - i\alpha_0)\dot{m}^- - iAm^- - iBm^+ &= -i\omega_s (\cos\omega_1 t)m^-\end{aligned}\quad (21)$$

where the terms are defined as follows

$$\begin{aligned}m^\pm &= m_x \pm im_y \\ A &= \omega_M(N_z - \frac{1}{2}N_x - \frac{1}{2}N_y) \\ B &= \frac{1}{2}\omega_M(N_y - N_x)\end{aligned}\quad (22)$$

The coupled equations of motion of Equation (21) can be decoupled in the homogeneous case by making the change of variables

$$\begin{aligned}m^+ &= an^+ - bn^- \\ m^- &= -b^*n^+ + an^-\end{aligned}\quad (23)$$

where

$$\begin{aligned}\cosh x &= A/\omega_0 \\ \sinh x &= B/\omega_0 \\ a &= \cosh x/2 \\ b &= (1 - i\alpha_0) \sinh x/2 \\ b^* &= (1 + i\alpha_0) \sinh x/2 \\ \omega_0^2 &= A^2 - B^2 = \omega_M^2 (N_z - N_x)(N_z - N_y)\end{aligned}\quad (24)$$

Neglecting all terms containing  $\alpha_0^2$ , the decoupled equations of motion become

$$\begin{aligned}\dot{n}^+ + (i\omega_0 + \alpha_0 A)n^+ &= (i\frac{\omega_s}{\omega_0} A \cos\omega_1 t)n^+ - (i\frac{\omega_s}{\omega_0} B \cos\omega_1 t)n^- \\ \dot{n}^- - (i\omega_0 - \alpha_0 A)n^- &= - (i\frac{\omega_s}{\omega_0} A \cos\omega_1 t)n^- + (i\frac{\omega_s}{\omega_0} B \cos\omega_1 t)n^+\end{aligned}\quad (25)$$

The homogeneous solution is obtained by equating the right sides of Equations (26) to zero and is given by

$$\begin{aligned} n_o^+ &= n_o^+ e^{-(i\omega_o + \alpha_o A)t} \\ n_o^- &= n_o^- e^{-(i\omega_o + \alpha_o A)t} \end{aligned} \quad (26)$$

The particular solution is obtained by differentiating Equations (26) with respect to time, with  $n_o^+$  and  $n_o^-$  considered time variable, and substituting into Equations (25). The driving terms on the right sides of the resulting equations will contain time variations of the form  $\omega_1 \pm \omega_o$  and  $-\omega_1 \pm \omega_o$ . Only those terms which have time variations nearly equal to the natural frequencies,  $\pm \omega_o$ , will be effective. Such terms exist only for

$$\omega_1 \approx 2\omega_o \quad (27)$$

Assuming Equation (27) and neglecting off resonant driving terms gives

$$\dot{n}_o^+ \approx -i\left(\frac{\omega_s}{\omega_o}\right)\left(\frac{B}{2}\right) n_o^- e^{-i(\omega_1 - 2\omega_o)t} \quad (28)$$

$$\dot{n}_o^- \approx i\left(\frac{\omega_s}{\omega_o}\right)\left(\frac{B}{2}\right) n_o^+ e^{i(\omega_1 - 2\omega_o)t}$$

Differentiating Equations (28) again with respect to time and rearranging terms gives

$$n_o^+ \pm i(\omega_1 - 2\omega_o)N_o^+ - \left(\frac{B}{2}\right)^2 \left(\frac{\omega_s}{\omega_o}\right)n_o^+ = 0 \quad (29)$$

Assuming a solution of the form

$$n_o^+ = N_o^+ e^{-(i\omega \pm \beta)t} \quad (30)$$

where  $\omega$  is an oscillatory part and  $\beta$  a damping part of the time variation and substituting into Equation (29), gives, by equating real and imaginary terms,

$$\omega = -\left(\frac{\omega_1}{2} - \omega_o\right) \quad (31)$$

$$\beta^2 = \left(\frac{B}{2}\right)^2 \left(\frac{\omega_s}{\omega_o}\right)^2 - \left(\frac{\omega_1}{2} - \omega_o\right)^2$$

The two solutions for  $n_o^+$  and for  $n_o^-$  then consist of a low frequency oscillation with an exponentially increasing amplitude for one and an exponentially decreasing amplitude for the other.

From Equation (25) it is seen that instability will occur if  $n_o^+$  grows more rapidly than the exponential decay terms in  $n_o^+$ . The threshold of instability is given by

$$\beta^2 = (\alpha_o A)^2 \quad (32)$$

Substituting the value of  $\beta^2$  from Equation (31) and the value of B from Equation (21) into Equation (32) and solving for  $\omega_s$  gives

$$\omega_{s \text{ crit}}^2 = \frac{16 \left( \frac{\omega_o}{\omega_m} \right)^2}{(N_y - N_x)^2} \left[ \left( \frac{\omega_1}{2} - \omega_o \right)^2 + (\alpha_o \omega_o)^2 - \alpha_o \left( \frac{\omega_M}{2} \right)^2 (N_y - N_x)^2 \right] \quad (33)$$

The minimum instability threshold occurs when  $\omega_1 = 2\omega_o$ . For this case we have

$$\left( \frac{h_1}{\Delta H} \right)_{\text{crit}} = \frac{\left[ 1 - \left( \frac{\omega_M}{\omega_1} \right)^2 (N_y - N_x)^2 \right]^{\frac{1}{2}}}{\left( \frac{\omega_M}{\omega_1} \right) (N_y - N_x)} \quad (34)$$

where

$$\Delta H \equiv \frac{\alpha_o \omega_o}{\gamma} \quad (35)$$

is the low power resonance linewidth.

Let  $N_x = 0$  and recall that

$$N_x + N_y + N_z = 1 \quad (36)$$

Then solving for  $(N_y - N_x)$  in terms of  $\left( \frac{\omega_M}{\omega_1} \right)$ , using the definition of  $\omega_o$  from Equation (24) and the assumption that  $\omega_1 = 2\omega_o$ , gives the normalized minimum critical field as

$$\left(\frac{h_1}{\Delta H}\right)_{\text{crit}} = \frac{4 \sqrt{1 - \frac{1}{8} \left(\frac{\omega_1}{\omega_M}\right)^2} \left[ 5 - 3 \sqrt{1 + 2 \left(\frac{\omega_1}{\omega_M}\right)^2} + \left(\frac{\omega_1}{\omega_M}\right)^2 \right]}{\left(\frac{\omega_M}{\omega_1}\right) \left[ 3 - \sqrt{1 + 2 \left(\frac{\omega_1}{\omega_M}\right)^2} \right]} \quad (37)$$

The critical field obtained by this approximation is the field which will cause the entire domain to go unstable simultaneously. In this approximation, the spin system throughout the domain is in synchronism. For this model the critical field approaches infinity as the ratio  $\left(\frac{\omega_M}{\omega_1}\right)$  approaches 0.5 and is imaginary for smaller values of  $\left(\omega_M/\omega_1\right)$ . Since finite critical fields have been measured below  $\left(\omega_M/\omega_1\right) = 0.5$ , this model is obviously too crude to be useful in the range of small  $\left(\omega_M/\omega_1\right)$ . As the ratio  $\left(\omega_M/\omega_1\right)$  becomes larger the critical field given by Equation (37) grows smaller and the approximation is somewhat better. A more accurate approximation of the critical field for all ranges of  $\left(\omega_M/\omega_1\right)$  is given by the second order differential equations which are taken up in the next section.

#### IV. SECOND ORDER SOLUTION

Retaining the definitions of Equation (22) and including first and second order terms from Equations (14) through (19) gives the coupled equations of motion as

$$\begin{aligned} (1 \pm i\alpha_0)m^{\pm} + (1 \mp i\alpha_k)\dot{\delta m}^{\pm} = \pm (i\omega_s \cos\omega_1 t)m^{\pm} \mp iA_m^{\pm} \mp iB_m^{\pm} \\ \pm (i\omega_s \cos\omega_1 t)\dot{\delta m}^{\pm} \pm iA_k^{\pm}\delta m^{\pm} \pm iB_k^{\pm}\delta m^{\pm} \end{aligned} \quad (38)$$

where

$$\begin{aligned} A_k^{\pm} = -\omega_{Mz} + a^2\omega_x k^2 + \frac{\omega_M}{2} \frac{k^+ k^-}{k^2} = -\omega_{Mz} + a^2\omega_E k^2 + \frac{\omega_M}{2} \sin^2\theta \\ B_k^{\pm} = \frac{\omega_M}{2k^2} (k^{\pm})^2 = \frac{\omega_M}{2} e^{\pm i2\varphi} \sin^2\theta \end{aligned} \quad (39)$$

The angle  $\theta$  is the angle between the  $z$  axis of the coordinate system and the direction of the spin wave vector  $\underline{k}$ . The angle  $\varphi$  is the angle between the projection of  $\underline{k}$  on the  $xy$  plane and the  $x$  axis. Equation (38) can be solved approximately by applying the technique used in solving the differential equations in Part II. First the homogeneous equations for  $m^{\pm}$  and  $\delta m^{\pm}$  are solved separately. These solutions are then inserted into the particular equation with the constants from the homogeneous solution allowed to be time variable. Next all driving terms which do not have time variations nearly equal to the resonant frequencies from the homogeneous solution are discarded. The resultant equations can then be decoupled by a transformation similar to the one in Part II and the resulting equations solved directly.

The homogeneous solution for  $m^{\pm}$  was obtained in Part II and is given by Equation (26). The homogeneous equations for  $\delta m^{\pm}$  are given by

$$(1 \mp i\alpha_k)\dot{\delta m}_k^{\pm} = \pm iA_k^{\pm}\delta m^{\pm} \pm iB_k^{\pm}\delta m^{\mp} \quad (40)$$

Making the transformation of variables

$$\delta m^+ = a_k \delta n^+ - b_k \delta n^- \quad (41)$$

$$\delta m^- = -b_{-k} \delta n^+ + a_k \delta n^-$$

gives

$$\delta \dot{n}^+ + (i\omega_k + \alpha_k A_k) \delta n^+ = 0 \quad (42)$$

where

$$\cosh x_k = A_k / \omega_k$$

$$\sinh x_k = B_k / \omega_k \quad (43)$$

$$a_k = \cosh x_k / 2$$

$$b_{\pm k} = (1 \pm i\alpha) e^{\pm i2\varphi} \sinh x_k / 2$$

$$\omega_k^2 = A_k^2 - B_k^2 = (a^2 \omega_x k^2 - N_z \omega_M) (\omega_M \sin^2 \theta + a^2 \omega_x k^2 - N_z \omega_M)$$

Equation (42) has solutions of the form

$$\delta n^+ = \delta n_0^+ e^{+(i\omega_k + \alpha_k A_k)t} \quad (44)$$

From Equations (23) and (26) it can be seen that  $m^+$  and  $m^-$  have frequency components at  $+\omega_0$  and at  $-\omega_0$ . Similarly, from Equations (41) and (44),  $\delta m^+$  have frequency components at  $+\omega_k$  and at  $-\omega_k$ . Assuming that these terms dominate, consider the frequency components of the terms in Equation (38). These are summarized in Table I. The equations of motion are now best solved in approximate form by separately considering various cases in the frequency relationships.

TABLE I. APPROXIMATE FREQUENCIES OF TERMS IN  
SECOND ORDER MAGNETIZATION EQUATION

$(1 + i\alpha)_m^+$	$(1 + i\alpha_k)\delta m^+$	$i\omega_s \cos \omega_1 t m^+$	$iA m^+$	$iB m^+$	$i\omega_s \cos \omega_1 t \delta m^+$	$A_k \delta m^+$	$B_k \delta m^+$
$\omega_0$	$\omega_k$	$\omega_1 + \omega_0$	$\omega_0$	$-\omega_0$	$\omega_1 + \omega_k$	$\omega_k$	$-\omega_k$
$-\omega_0$	$-\omega_k$	$\omega_1 - \omega_0$	$-\omega_0$	$+\omega_0$	$\omega_1 - \omega_k$	$-\omega_k$	$\omega_k$
		$-\omega_1 + \omega_0$			$-\omega_1 + \omega_k$		
		$-\omega_1 - \omega_0$			$-\omega_1 - \omega_k$		

A. Case 1:  $\omega_1 \neq 2\omega_0$ ;  $\frac{1}{2}\omega_1 \approx \omega_k$ ;  $\omega_0 \neq \omega_k$

From Equation (24) it is seen that  $\omega_0$  is a function of the ferrite material through  $\omega_M$  and a function of the domain shape through the demagnetizing factors  $N_x$ ,  $N_y$  and  $N_z$ . For a given domain then,  $\omega_0$  is a constant. The driving frequency,  $\omega_1$ , is controlled so that the first inequality of this case can always be established. Spin waves in the sample can exist over a very broad frequency spectrum. The approximate equality  $\omega_k \approx \frac{1}{2}\omega_1$  will always be satisfied for a portion of the spin wave spectrum regardless of the value of  $\omega_1$  and will be the dominant term. Other spin waves with frequencies  $\omega_k \neq \frac{1}{2}\omega_1$  will not be driven by the applied rf field and can be neglected.

Considering only those terms of Equation (38) which vary at  $\omega_k$  as determined from Table I gives

$$(1 + i\alpha_k)\delta m^+ = \pm i\omega_s \cos \omega_1 t \delta m^+ \pm iA \delta m^+ B_k \delta m^+ \quad (45)$$

Again making the change of variables indicated in Equation (41) and assuming a solution of the form given in Equation (44) with  $\delta n_o^+$  time varying, gives

$$\begin{aligned}\dot{\delta n_o^+} &\approx -\frac{i}{2} |B_k| \frac{\omega_s}{\omega_k} e^{i2\varphi} e^{i(\omega_1 - 2\omega_k)t} \delta n_o^- \\ \dot{\delta n_o^-} &\approx \frac{i}{2} |B_k| \frac{\omega_s}{\omega_k} e^{-i2\varphi} e^{-(\omega_1 - 2\omega_k)t} \delta n_o^+\end{aligned}\quad (46)$$

Taking second time derivatives and combining equations gives

$$\left[ \frac{\partial^2}{\partial t^2} + i(\omega_1 - 2\omega_k) - \left( \frac{|B_k|}{2} \frac{\omega_s}{\omega_k} \right)^2 \right] \delta n_o^+ = 0 \quad (47)$$

Assuming a solution of the form

$$\delta n_o^+ = \delta N_o^+ e^{+i(\omega \pm \kappa)t} \quad (48)$$

where  $\omega$  is the oscillatory and  $\kappa$  the loss parts of the time variation and solving for  $\omega$  and  $\kappa$  gives

$$\begin{aligned}\omega &= \left( \frac{\omega_1}{2} - \omega_k \right) \\ \kappa^2 &= \left( \frac{|B_k|}{2} \frac{\omega_s}{\omega_k} \right)^2 - \left( \frac{\omega_1}{2} - \omega_k \right)^2\end{aligned}\quad (49)$$

The instability threshold for this case occurs when

$$\kappa^2 = \alpha_k^2 A_k^2 \quad (50)$$

Solving Equation (50) for  $\omega_s^2$  gives

$$\omega_s^2_{\text{crit}} = (\gamma h_1)_{\text{crit}}^2 = \frac{4\omega_k^2}{|B_k|^2} \left[ \left( \frac{\omega_1}{2} - \omega_k \right)^2 + \alpha_k^2 A_k^2 \right] \quad (51)$$



The minimum critical field occurs when  $\omega_1 = 2\omega_k$  which gives

$$\left(\frac{h_1}{\Delta H_k \text{ crit}}\right) = \frac{\left(\frac{\omega_1}{\omega_M}\right)}{\sin^2 \theta} \frac{\left(1 - \frac{1}{2} \frac{\omega_M \sin^2 \theta}{a^2 \omega_k^2 - \omega_M N_z}\right)}{\sqrt{1 - \frac{\omega_M \sin^2 \theta}{a^2 \omega_k^2 - \omega_M N_z}}} \quad (52)$$

where

$$\Delta H_k \equiv \frac{2\alpha_k \omega_k}{\gamma} \quad (53)$$

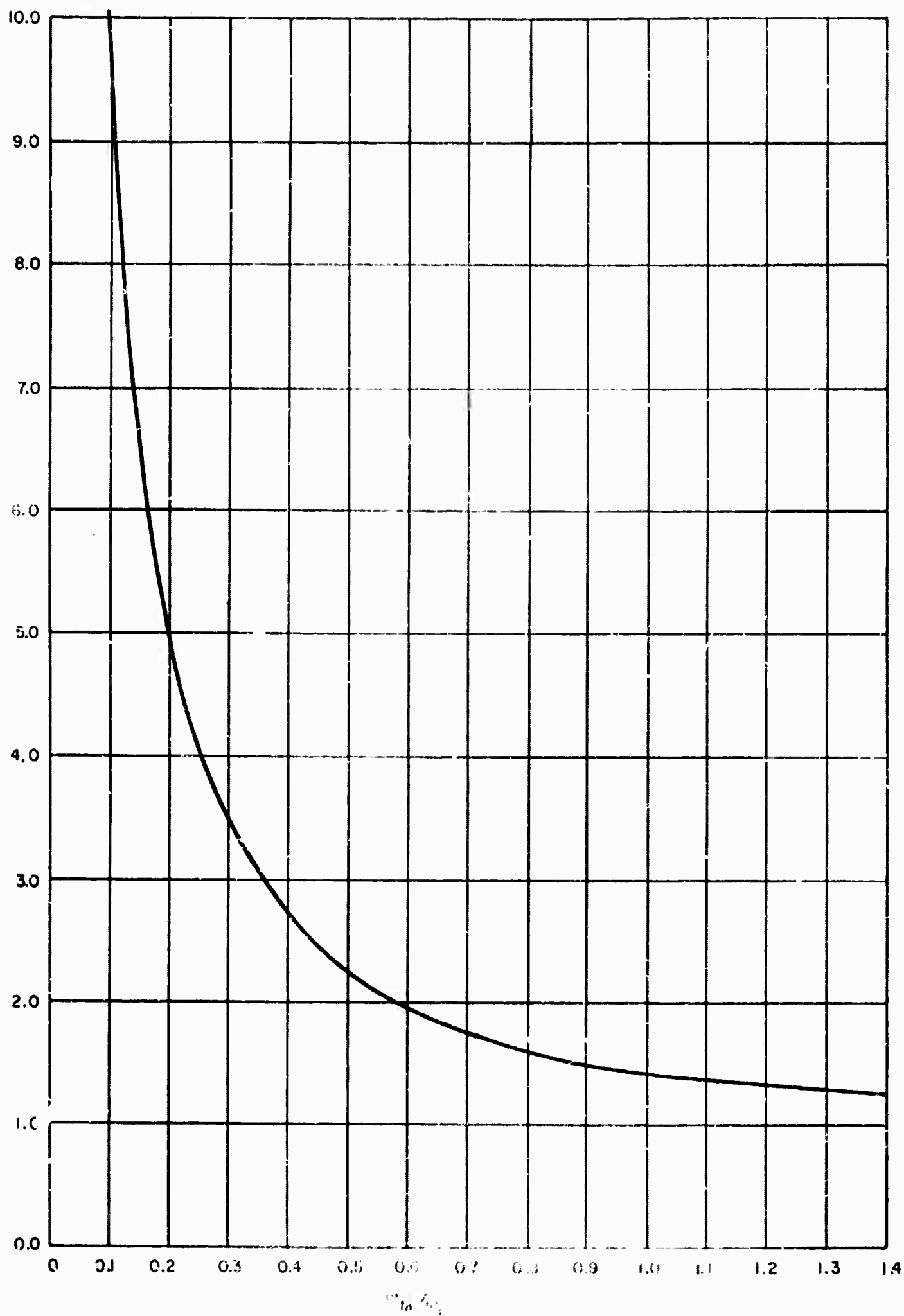
is the spin wave linewidth. Using the definition of  $\omega_k$  from Equation (43) and the assumption that  $\omega_1 = 2\omega_k$ ,  $(a^2 \omega_k^2 - \omega_M N_z)$  can be evaluated and substituting into Equation (42) to obtain the result

$$\left(\frac{h_1}{\Delta H_k \text{ crit}}\right) = \frac{\sqrt{1 + \left(\frac{\omega_M}{\omega_1}\right)^2 \sin^4 \theta}}{\left(\frac{\omega_M}{\omega_1}\right) \sin^2 \theta} \quad (54)$$

Equation (54) is further minimized when  $\sin^2 \theta = 1$  so that

$$\left(\frac{h_1}{\Delta H_k \text{ crit}}\right) = \frac{\sqrt{1 + \left(\frac{\omega_M}{\omega_1}\right)^2}}{\left(\frac{\omega_M}{\omega_1}\right)} \quad (55)$$

This expression is plotted in Figure 1.



Case 2:  $\omega_1 - \omega_0 \approx \omega_k \neq \omega_0$

The equations of motion for this case are given by

$$(1 + \alpha_k) \dot{m}^+ + iA_k m^+ + iB_k \dot{m}^+ = \pm i\omega_s \cos \omega_1 t m^+ \quad (56)$$

which correspond to a coupled pair of lossy harmonic oscillators on the left with a harmonic driving term on the right. The equations can be decoupled by the coordinate transformation of Equation (41), however, one need not do so to determine that the solutions of Equations (56) are stable harmonic oscillators as long as the right sides of the equations have constant amplitude, i.e.  $m^+$  constant in amplitude. The critical field for this case would then be given by Equation (37) which is the field required to make  $m^+$  grow exponentially.

Case 3:  $\omega_k = \omega_0 = \frac{1}{2} \omega_1$

For this case all of the terms from Equation (38) must be retained in part and a direct solution is too complex algebraically to be solved. A reasonable approximation to this case might be given by Equation (55) for low values of  $\omega_M/\omega_1$  and by the lower of the values from Equations (37) and (55) for large values of  $\omega_M/\omega_1$ .

## V. DISCUSSION OF RESULTS AND HIGHER ORDER TERMS

The first and second order solutions both indicated that the normalized critical field increases monotonically as the ratio  $\omega_M/\omega_1$  decreases. The first order solution does not have a finite critical field for  $\omega_M/\omega_1$  less than 0.5 and is thus inadequate for low values of  $\omega_M/\omega_1$ . The second order solution has finite values for the normalized critical field for all non-zero values of  $\omega_M/\omega_1$ . For this reason, and because it considers more of the interacting energy terms, the second order solution is to be preferred.

There are, however, several possibilities for error in the second order solution. One of the most serious problems is the normalizing term  $\Delta H_k$  which is the spin wave linewidth. This term is related to the spin wave loss or relaxation, and is expected to be a function of frequency and material and possibly even the shape of the sample and domain rather than a constant. Until this term is known more explicitly, Equation (55) and Figure 1 can only be considered as rough approximations.

It was assumed that the demagnetizing field due to spacial average magnetization in the domain could be expressed in terms of the demagnetizing factors  $N_x$ ,  $N_y$  and  $N_z$ . These factors apply precisely only for ellipsoidal samples which are small compared to the wavelength of the externally applied fields. For other shapes constant N factors can be assumed throughout the bulk of the domain which will closely approximate the physical situation. However, near the domain boundaries and particularly near sharp edges, the demagnetizing field can be a rapidly varying function of position. This can be accounted for by allowing the demagnetizing factors to be functions of position near the domain boundaries and constant throughout most of the domain. Unfortunately, computation of the functional dependence of the demagnetizing field (or factors) is a prohibitively difficult task for almost all shapes of interest.

The effect of the variation in  $N_z$  can be seen from the spin wave dispersion diagram shown in Figure 2. The abscissa scale in Figure 2 is normalized by the factor  $\sqrt{\omega_M/a^2\omega_E}$  which depends on the material constants  $\gamma$ ,  $M_s$ ,  $a^2$ , and  $H_E$  and which makes Figure 2 a general curve applicable to all ferrites.

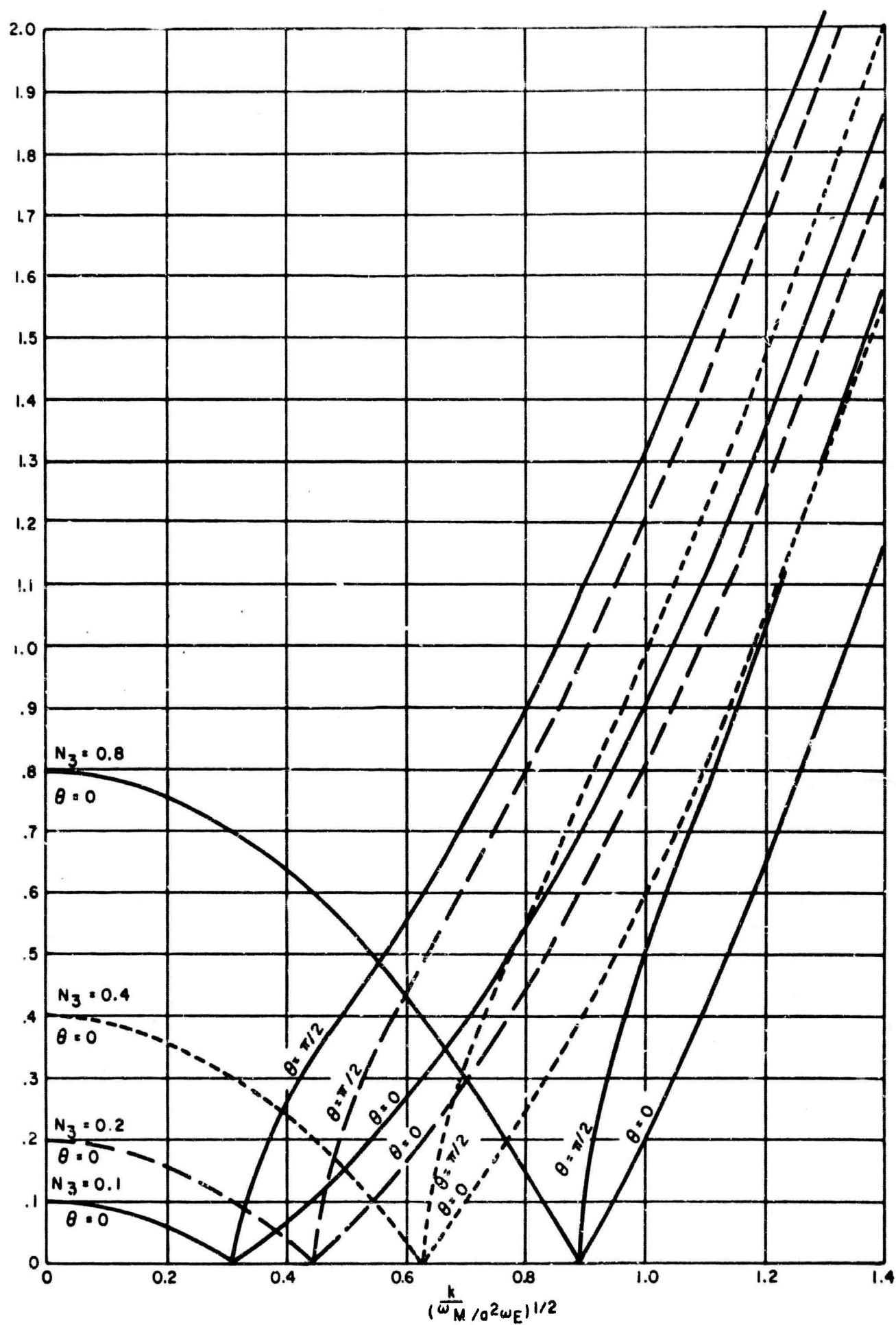


FIGURE 2. Spin Wave Dispersion Diagram with Remnant Magnetization

For a fixed value of  $\omega_k/\omega_M$  (recalling that the applied  $\omega_1 = 2\omega_k$ ) it is seen that for various values of  $N_z$  spin waves of various wave numbers or wave lengths can be generated in different regions of the sample. If  $N_z$  varies rapidly enough to give interaction between the local demagnetizing fields,  $\delta h$ , Equation (5) may no longer be valid. In this case Equation (4) would not be valid either since it represents the case in which a spectrum of spin waves is generated in a given local region. For the case in which spin waves of different wave numbers  $k$  are generated in neighboring regions within the domain Equation (4) should be replaced by a weighted sum in which the weighting terms would be at least a function of the distance between regions. It can also be noted from Figure 2 that for small values of  $\omega_k/\omega_M$  and for some angles  $\theta$  two spin waves with the same frequency can be generated with different values of wave number  $k$ . This is possible for values of  $\theta$  which satisfy

$$\sin^2 \theta < N_z \quad (57)$$

and values of  $\frac{\omega_k}{\omega_M}$  which satisfy

$$\frac{\omega_k}{\omega_M} < N_z \quad (58)$$

This situation again makes Equation (5) invalid unless a second term is added to account for the degeneracy. This situation is not too serious however, since it is expected  $\omega_k/\omega_M > N_z$  for most practical cases of interest here.

So far only the first and second order solutions to the instability threshold problem have been considered. Some qualitative statements regarding the third order solution can be made. The exchange terms contain no third order quantities and so remain unchanged. The third order loss

terms are off frequency and can be neglected. The third order Zeeman terms allow direct coupling from the average magnetization terms to the local deviation terms. Neglecting off frequency terms, the third order equation for the magnetization can be written as

$$(1 - i\alpha_k) \delta m^+ - iA_k \delta m^+ - iB_k^+ \delta m^- = i \left[ \omega_s \cos \omega_1 t - \frac{k}{k^2} (k^- m^+ + k^+ m^-) \right] \delta m^+ - i\gamma \frac{k_- k^+}{k^2} m^+ \delta m^- \quad (59)$$

If the driving term containing  $\omega_s$  is much larger than the remaining driving terms, Equation (59) reduces to the second order solution. If the term  $\omega_s$  negligibly small compared to the other driving terms, Equation (59) will reduce to Equation (21) of Suhl's analysis. In general, the local spin system can gain energy from the driving field through the  $\omega_s$  term or gain energy from or lose energy to the spacial average spin system through the cross product terms. A complete solution should indicate that each driving term will result in one exponentially increasing local spin wave and one exponentially decreasing one above some threshold value. If the wave which grows from the  $\omega_s$  pumping term corresponds to the decaying wave from the cross product terms, the threshold would be expected to be higher than that computed in the second order solution resulting from the transfer of pump energy from the local spin deviations to the spacial average. This would tend to distribute the total spin energy in the sample more uniformly and thus seem to be a correct evaluation from a thermodynamic approach. Except in the case where  $\omega_1 = 2\omega_0$  as well as  $\omega_1 = 2\omega_k$ , an increase in the spacial average magnetization due to direct pumping is not expected. In this case a noticable increase in the average magnetization due to indirect coupling from locally deviating spin system which is being pumped would not be expected because of the large difference in the number of spins involved in the local variation as compared to the domain as a whole. Therefore, except in the special case of pumping at twice the bulk resonance frequency of the domain  $\omega_0$ , the normalized critical field can be expected to be somewhat larger than that computed in the second order solution due to the third order Zeeman terms and the instability threshold will be caused by the direct coupling from the applied rf magnetic field to the local spin system.

### PART III. PROPAGATION OF TE MODES IN DIELECTRICALLY LOADED WAVEGUIDES

#### ABSTRACT

The propagation of TE modes in rectangular waveguides which contain two dielectric slabs parallel to the narrow wall and extending over the full height of the guide is investigated. Waveguide and dielectric are assumed to be lossless and infinitely long. Apart from these restrictions the dielectric slabs may have arbitrary thickness, position and dielectric constant. The analysis is restricted to TE modes with the E-field parallel to the narrow guidewall. The guide containing only one dielectric slab is covered by this analysis. The even modes  $n = 2, 4, 6, \dots$  of the guide with two slabs correspond to the odd modes  $n' = n/2 = 1, 2, 3, \dots$  of the guide with one slab and half the width of the guide with two slabs.

For the relative dielectric constants  $\epsilon = 2.25, 4, 9, 12.25, 16, 25$  the normalized cutoff frequencies for TE 10, 20, 30, 40, 60 modes have been computed for fifteen slab thicknesses (including 0 and 100% filling factor) and a maximum of eleven slab positions, normalized propagation constants have been computed for five slab thicknesses and a maximum of eight slab positions for TE 10 and TE 20 modes between their respective cutoff frequencies and a frequency slightly above the second and fourth order mode cutoff frequency of the empty guide, respectively. The results are presented graphically and numerically.

These results are discussed. The parametric dependence of field distributions, the ratio of magnetic field components (ellipticity), the ratio of cutoff frequencies (fractional bandwidth), and normalized wave impedances are illustrated.

<u>Contents</u>	<u>Page</u>
Introduction .....	1
Quantities and Symbols .....	1
Theory .....	3
Computer Results .....	7
Discussion of Results .....	9
Acknowledgements .....	12
Bibliography .....	16
Normalized Cutoff-Frequencies (graphical) .....	17
Normalized Propagation Characteristics (graphical) ..	23
Normalized Cutoff-Frequencies (numerical) .....	53
Normalized Propagation Constants (numerical) .....	59



## PROPAGATION OF TE MODES IN DIELECTRICALLY LOADED WAVEGUIDES

### INTRODUCTION

Previous analyses of propagation of  $TE_{no}$  modes in rectangular waveguides which contain dielectric slabs have in general been restricted to two cases, where the dielectric slab is placed a) against a waveguide wall (1, 2, 3, 4), or b) in the center plane of the waveguide (2,3,4,5,6). A few more special cases have been considered by workers dealing with ferrite applications in the microwave region (4, 8). For a rather general position of the dielectric slab certain phase shift characteristics of the loaded guide have been calculated (7).

The present analysis deals with a rectangular waveguide which contains two dielectric slabs parallel to the narrow walls and extending over the full height of the guide. The slabs are placed symmetrically with respect to the center E-plane of the guide. Apart from this restriction the slabs have arbitrary position, thickness and dielectric constant. Only  $TE_{no}$  -modes with E-fields parallel to the narrow guide wall are considered. The guide and the dielectric are assumed to be lossless and infinitely long. Waveguides containing only one dielectric slab are covered by this analysis.

### QUANTITIES AND SYMBOLS

Figure 1 shows the cross-section of a rectangular, dielectrically loaded waveguide in a rectangular coordinate system. The broad dimension of the guide extends along the x-axis, y is the direction of propagation of fields in the guide. a, c, d, w are waveguide dimensions along the x-axis.  $\beta = \omega \sqrt{\mu_0 \epsilon_0}$  is the free-space propagation constant, k is the propagation constant in the guide in the direction of the guide, p in the empty guide region and q in the dielectric are propagation constants transverse to the direction of the guide and the electric field. Instead of these symbols dimensionless quantities will be used throughout the analysis. These are obtained by either multiplying or dividing the above given quantities by w. Quantities, their dimensions (MKSA-sys) and symbols used here are:

<u>Symbol</u>	<u>Unit</u>	<u>Quantity</u>
$\alpha = a/w$	-	waveguide dimensions as shown in Figure 1
$\delta = d/w$	-	
$\gamma = c/w$	-	$\alpha + \delta + \gamma = 1$
$x, y, z$	m	right hand coordinate system as shown in Fig. 1
$\varphi = x/w$	-	normalized x-coordinate
$i_x, i_y, i_z$	-	unit vectors
$t$	s	time
$\omega$	$s^{-1}$	angular frequency
$B = \beta w$	-	free-space propagation number(frequency parameter)
$B_c = \beta_c w$	-	normalized cutoff frequency
$K = kw$	-	longitudinal propagation number in the guide
$P = pw$	-	transverse propagation number in empty part of guide
$Q = qw$	-	transverse propagation number in the dielectric
$\rho_1, \dots, \rho_6$	radian	electrical widths of waveguide sections
$\theta$	radian	phase angle
$C, D$	-	relative amplitudes
$\mu_o$	$VsA^{-1}m^{-1}$	free-space permeability
$\epsilon_o$	$AsV^{-1}m^{-1}$	free-space dielectric constant
$\epsilon$ or DK	-	relative dielectric constant of dielectric (DK appears in the computer printed tables)
$\vec{E} = i_z E$	$Vm^{-1}$	electric field
$E_o$	$Vm^{-1}$	normalizing electric field
$\vec{H} = i_x H_x + i_y H_y$	$Am^{-1}$	magnetic field
$H_o$	$Am^{-1}$	normalizing magnetic field
$Y_W$	$AV^{-1}$	wave admittance of loaded guide
$Y_o$	$AV^{-1}$	wave admittance of free space
$n$	-	fractional bandwidth
ELL	-	ellipticity

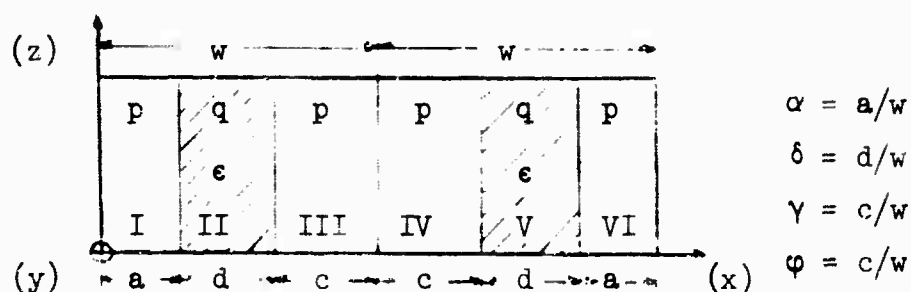


Figure 1. Waveguide with Dielectric Slabs

## THEORY

It completely suffices, because of the symmetry of the loaded guide, to consider the regions I, II, III of Figure 1, so that  $0 \leq x \leq w$  or  $0 \leq \varphi \leq 1$ . All fields vary as  $\exp(\omega t - ky)$ , so that  $\partial/\partial t = j\omega$  and  $\partial/\partial y = -jk$ . This  $t$  and  $y$  dependence is omitted in all equations. The relative permeability of the dielectric is assumed to be 1. Maxwell's equations for the problems considered here reduce to

$$(1) \quad \nabla \times \vec{E} = i_x \frac{\partial E}{\partial y} - i_y \frac{\partial E}{\partial x} = -j\omega\mu_0 (i_x H_x + i_y H_y)$$

$$(2) \quad \nabla \times \vec{H} = i_z \left( \frac{\partial H_y}{\partial x} - \frac{\partial H_x}{\partial y} \right) = i_z j\omega\epsilon_0 \epsilon E$$

$$(3) \quad \nabla \cdot \vec{E} = 0$$

$$(4) \quad \nabla \cdot \vec{H} = 0$$

and from equation (1, 2, 3)

$$(5) \quad \nabla \times \nabla \times \vec{E} = - \left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} \right) \vec{E} = \omega^2 \mu_0 \epsilon_0 \epsilon E$$

The boundary conditions are

$$(6) \quad E_{\text{tang}} = E \text{ is continuous, } H_{\text{tang}} = H_y \text{ is continuous}$$

The fields in the various regions of the guide can be described as in Table 1 with dimensionless quantities:

Region	Field Modes	$H_x/H_0 = H_x \omega\mu_0/KE_0, E/E_0$			$jH_y/H_0 = jH_y \omega\mu_0/KE_0$		
		$0 \leq K/B < 1$	$K/B = 1$	$1 < K/B \leq \sqrt{\epsilon}$	$0 \leq K/B < 1$	$K/B = 1$	$1 < K/B \leq \sqrt{\epsilon}$
I $0 \leq \varphi \leq \alpha$	all	$\sin R\varphi$	$\varphi$	$\text{sh }  P \varphi$	$\frac{P}{K} \cos R\varphi$	$\frac{1}{K}$	$\frac{ P }{K} \text{ch }  P \varphi$
II $\alpha \leq \varphi \leq 1-\gamma$	all	$D \sin(Q\varphi + \theta)$			$D \frac{Q}{K} \cos(Q\varphi + \theta)$		
III $1-\gamma \leq \varphi \leq 1$	odd	$C \cos P(1-\varphi)$	$C$	$C \text{ch }  P (1-\varphi)$	$C \frac{P}{K} \sin P(1-\varphi)$	$0$	$-C \frac{ P }{K} \text{sh }  P (1-\varphi)$
	even	$C \sin P(1-\varphi)$	$C(1-\varphi)$	$C \text{sh }  P (1-\varphi)$	$-C \frac{P}{K} \cos P(1-\varphi)$	$-C \frac{1}{K}$	$-C \frac{ P }{K} \text{ch }  P (1-\varphi)$

Table 1. Normalized Field Distribution in Waveguide which Contains Dielectric Slabs

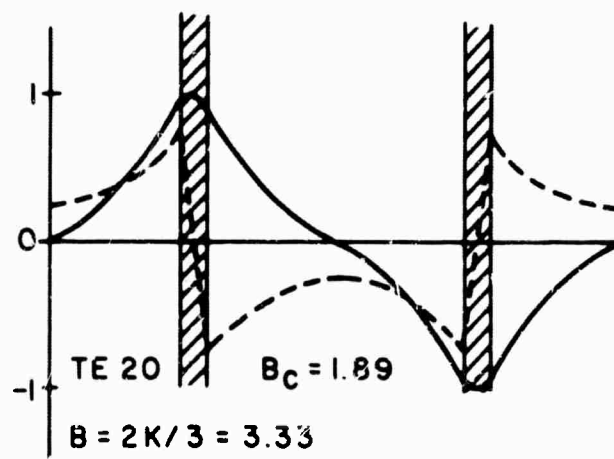
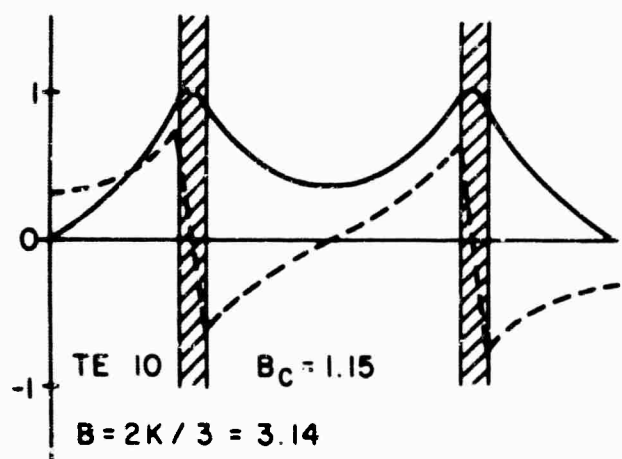
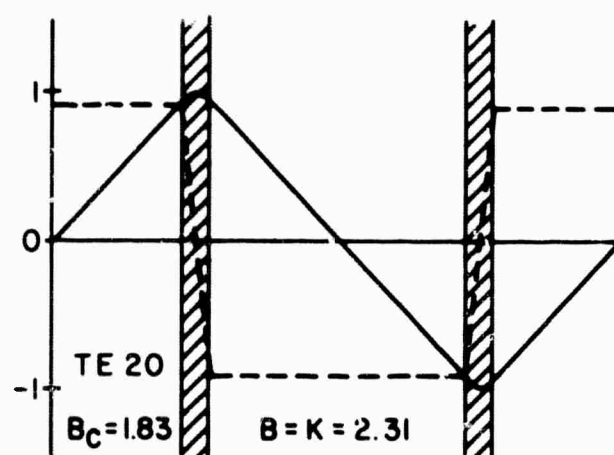
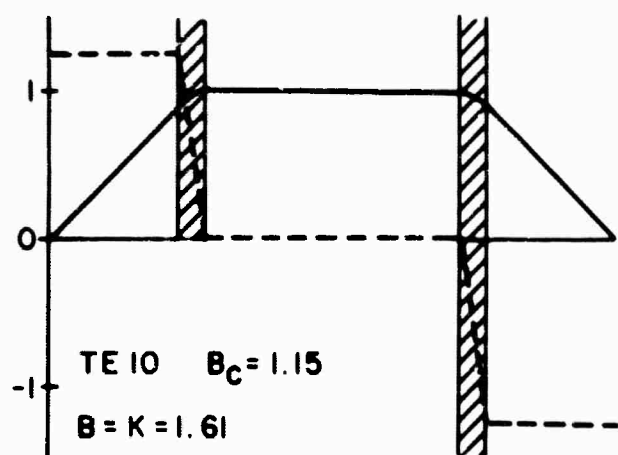
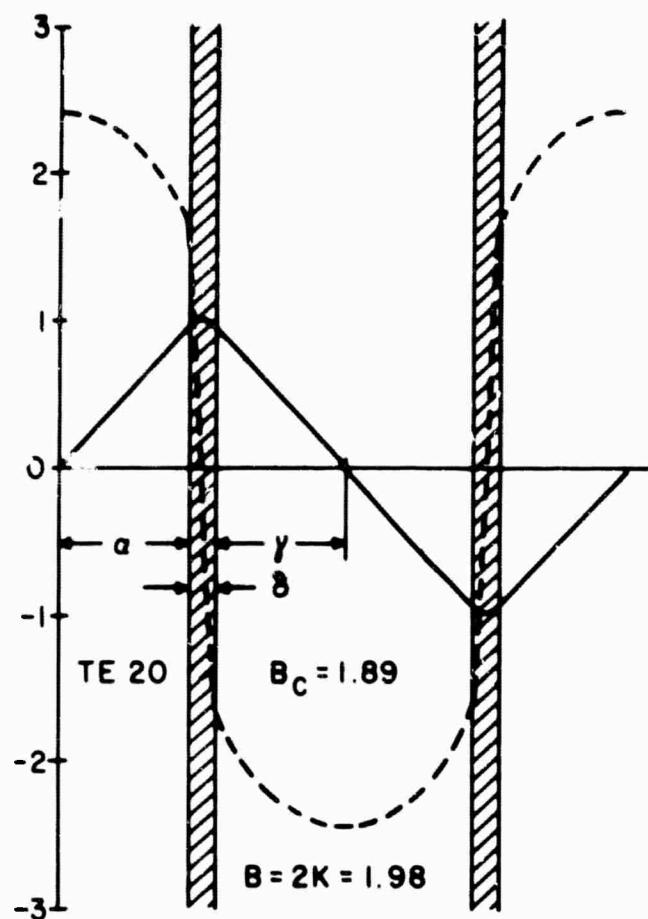
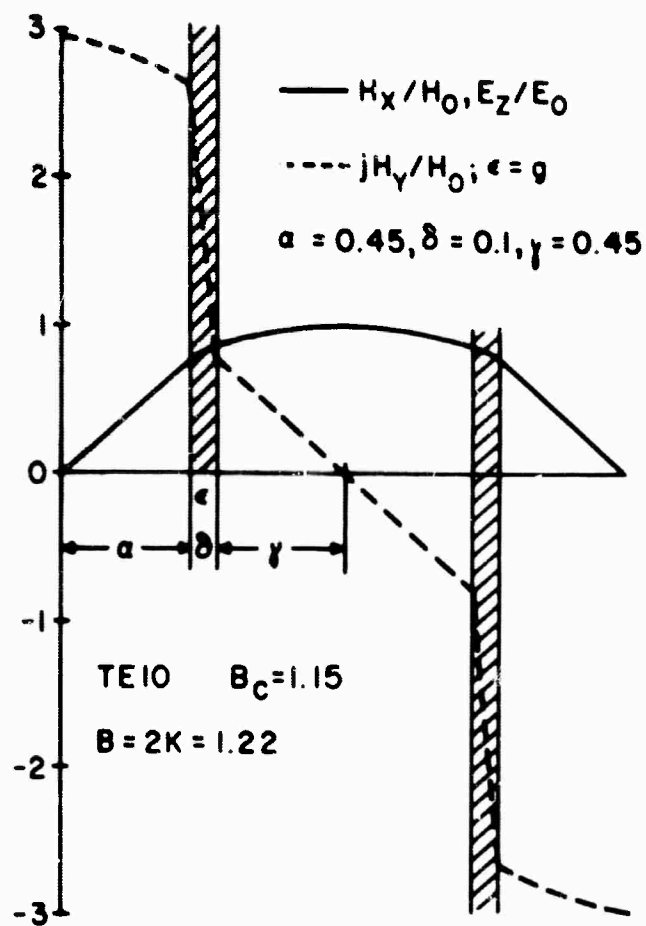


Figure 2. Normalized Field Distribution in Waveguide with Dielectric Slabs

These field distributions are illustrated by Figure 2 for the first odd ( $n = 1$ ) and first even ( $n = 2$ ) mode. The even modes ( $n = 2, 4, 6, \dots$ ) of the waveguide under consideration with two symmetrically placed dielectric slabs, where  $2w$  is the waveguide width, correspond to the odd and even modes ( $n' = n/2 = 1, 2, 3, \dots$ ) of a waveguide with only one dielectric slab, where  $w$  is the waveguide width; that is the left half of Figure 1 only.

The fields of Table 1, together with equation (5) give for the propagation constants the relations

$$(7) \quad P^2 = B^2 - K^2$$

$$(8) \quad Q^2 = \epsilon B^2 - K^2$$

$$(9) \quad P^2 = Q^2 - (\epsilon - 1) B^2$$

For any  $Q^2 > 0$  there exists a  $P^2 \geq 0$ , the sign depending on  $(\epsilon - 1)B^2$ .  $Q^2 > 0$  describes a sinusoidal field distribution in the dielectric slab,  $P^2 > 0$  describes a sinusoidal field distribution in the empty waveguide region. For  $P^2 < 0$  the field distributions in the empty waveguide regions are hyperbolic functions which describe a quasi-exponential decay.  $P^2 = 0$  gives the intermediate function where  $E$  and  $H_x$  have a constant slope in the empty part of the guide and where  $H_y$  is constant. In this particular case the fields in the empty region are described by functions as  $\varphi \exp(R\varphi) = \varphi$ , since  $P = 0$ . When  $Q^2 < 0$ , also  $P^2 < 0$ . Then the boundary conditions (6) are violated.  $H_y$  is no longer continuous. Therefore,  $Q^2 \geq 0$  for all  $TE_{no}$  modes.  $Q^2 = 0$  is reached at infinitely high frequencies. The assumptions of Table 1 cover only and all allowed field distributions for  $TE_{no}$  modes and are thus justified. While the frequency increases from cutoff to infinity,  $K/B$  increases from 0 to  $\sqrt{\epsilon}$ . For  $K = B$  the field distribution between the two dielectric slabs represents a pure TEM field for all odd  $TE_{no}$  modes in the waveguide considered.

The determinantal equations for  $B(K)$  - e.g. cutoff frequencies for  $k = 0$  - or  $K(B)$  - propagation constants for given frequencies - are found by expressing the widths of the empty sections of the guide in equivalent widths of a guide completely filled with the dielectric. The field distribution in the slabs is not changed by this replacement. The total electrical width ( $2w_e$ ) of this equivalent guide is  $n\pi$  for a  $TE_{no}$  mode.

At the boundary between regions I and II of Figure 1, Table 1 and equation (6) yield for frequencies  $0 \leq K/B < 1$  and with  $\rho_1 = P\alpha$ ,  $\rho_2 = Q\alpha + \theta$

$$(10) \quad E_0 \sin \rho_1 = E_0 D \sin \rho_2$$

$$(11) \quad jE_0 P \cos \rho_1 / \omega \mu_0 = jE_0 D Q \cos \rho_2 / \omega \mu_0$$

$$(12) \quad Z_x(\alpha) = -E_z(\alpha) / jH_y(\alpha) = j \frac{\omega \mu_0}{P} \operatorname{tg} \rho_1 = j \frac{\omega \mu_0}{Q} \operatorname{tg} \rho_2$$

$\omega \mu_0 / P = Z_P$  and  $\omega \mu_0 / Q = Z_Q$  are the transverse wave-impedances of the empty and loaded waveguide regions, respectively, for waves travelling in  $\pm x$  direction.  $Z_x(\alpha)$  is the impedance experienced by a wave travelling into the shorted waveguide of impedance  $Z_P$  and length  $\rho_1$  or impedance  $Z_Q$  and length  $\rho_2$ . From equation (12) follows the equivalent length

$$(13) \quad \rho_2 = \operatorname{arctg}\left(\frac{Q}{P} \operatorname{tg} \rho_1\right)$$

Similar considerations for higher frequencies, odd and even modes, and both slab boundaries lead for the various waveguide regions to the actual and equivalent electrical widths of Table 2.

Width \ Frequency		$0 \leq K/B < 1$	$K/B = 1$	$1 < K/B \leq \sqrt{\epsilon}$
I	$\rho_1$	$P\alpha$	$\alpha$	$P \alpha $
Equivalent I	$\rho_2$	$\operatorname{arctg}\left(\frac{Q}{P} \operatorname{tg} \rho_1\right)$	$\operatorname{arctg}(Q\rho_1)$	$\operatorname{arctg}\left(\frac{Q}{ P } \operatorname{th} \rho_1\right)$
II	$\rho_3$	$Q\delta$	$Q\delta$	$Q\delta$
III	$\rho_4$	$P\gamma$	0 $\gamma$	$P \gamma $
		+ odd modes even		
Equivalent III	$\rho_5$	$\operatorname{arctg}\left[\left(\frac{Q}{P}\right)^{(-1)^n} \operatorname{tg} \rho_4\right]$	0 $\operatorname{arctg}(Q\rho_4)$	$\operatorname{arctg}\left[(-1)^n \left(\frac{Q}{ P }\right)^{(-1)^n} \operatorname{th} \rho_4\right]$
I + I + III	$\rho_6$	$\rho_2 + \rho_3 + \rho_5$	$\rho_2 + \rho_3 + \rho_5$	$\rho_2 + \rho_3 + \rho_5$

Table 2. Electrical Widths in Waveguides, which Contain Dielectric Slabs

$\rho_5$  is defined for  $0 \leq K/B < 1$  for odd modes by  $1/P \tan \rho_4 = 1/Q \tan \rho_5$  and for even modes by  $\tan \rho_4/P = \tan \rho_5/Q$  and for other frequencies accordingly. Table 2 is the skeleton for the computer program used to determine the cutoff frequencies  $B(K = 0)$  and propagation constants  $K(B)$ . The determinantal equation is in both cases

$$(14) \quad \rho_6 = n\pi/2 \quad \text{for } TE_{no} \text{ -modes}$$

Applying equation (6) to the boundaries of regions I/II and II/III of Figure 1 yields the phase angle  $\theta$  and the relative amplitudes of Table 1. They are given in Table 3 in terms of the electrical widths defined in Table 2.

Frequency Modes		$0 \leq K/B < 1$	$K/B = 1$	$1 < K/B \leq \epsilon$
$\theta$	all	$\rho_2 - Q\alpha$	$\rho_2 - Q\alpha$	$\rho_2 - Q\alpha$
$1/D$	all	$\sin \rho_2 / \sin \rho_1$	$\sin \rho_2 / \rho_1$	$\sin \rho_2 / \sinh \rho_1$
$C/D$	odd	$\sin(\rho_2 + \rho_3) / \cosh \rho_4$	$\sin(\rho_2 + \rho_3)$	$\sin(\rho_2 + \rho_3) / \cosh \rho_4$
$C/D$	even	$\sin(\rho_2 + \rho_3) / \sinh \rho_4$	$\sin(\rho_2 + \rho_3) / \rho_4$	$\sin(\rho_2 + \rho_3) / \sinh \rho_4$

Table 3. Phase Angle and Relative Amplitude of Fields in Waveguides, Which Contain Dielectric Slabs

### COMPUTER RESULTS

Normalized cutoff frequencies have been calculated for  $TE_{no}$  modes of the guide of Figure 1 with  $n = 1, 2, 3, 4, 6$  for six relative dielectric constants (2.25, 4, 9, 12.25, 16, 25), fifteen slab thicknesses (including 0% and 100% filling factor) and a maximum of eleven positions of the slab in the guide. Normalized propagation constants have been calculated for  $TE_{10}$  and  $TE_{20}$  modes between their respective cutoff frequency and a frequency somewhat above the second and fourth order mode cutoff frequency of the empty guide, respectively, again for the relative dielectric constants given above, five slab thicknesses (5, 10, 15, 25, 40% filling factor) and a

maximum of eight slab positions. The position parameter  $\alpha + \delta/2$  gives the distance between the left guidewall and the center plane of the left slab as a fraction of half the guide width. For the odd modes the position is varied between the slabs touching the guide walls ( $\alpha + \delta/2 = \delta/2$ ) and the slabs touching each other in the center of the guide. ( $\alpha + \delta/2 = 1 - \delta/2$ ). For even modes it suffices, because of the symmetry of the field distribution, to vary the position between the slabs touching the wall and moving them half way towards each other ( $\alpha + \delta/2 = 0.5$ ). For even modes the cutoff-frequencies and propagation constants are the same for  $\alpha + \delta/2 = \tau$  and  $\alpha + \delta/2 = 1 - \tau$ . The normalized cutoff-frequencies are given graphically in Figures 8 ... 13 and numerically in Tables 4 ... 9. In these tables the slab thickness ( $\delta$ ) is the horizontal parameter, the position of the slab ( $\alpha + \delta/2$ ) and the order of modes (N) are the vertical parameters.

The normalized propagation constants K are given graphically in Figures 14 ... 43 with one set of curves for every dielectric constant and slab thickness combination and the position as parameter for every curve. They are given numerically in Tables 10 ... 29. The horizontal parameters are position ( $\alpha + \delta/2$ ) and order of mode ( $TE_{no}$ ). The vertical parameter is the normalized frequency (B). The normalized cutoff frequency ( $E_c$ ) for each set of parameters is given on top of each column. Propagation constants -0.0 indicate that the frequency for this K is below cutoff.

Examples:

- a) Guide WR 137, width 1.372 inches, two slabs each 0.069 inches thick, 0.206 inches between left wall and center of left slab, relative dielectric constant  $\epsilon = 9$ ,  $w = \text{width}/2 = 0.686$  inches,  $\delta = 0.1$ ,  $\alpha + \delta/2 = 0.3$ . Wanted:  $TE_{10}$  and  $TE_{20}$  cutoff-frequencies and guide wavelengths at 5.46 GHz. One finds:  $B_c(TE_{10}) = 1.33$ ,  $B_c(TE_{20}) = 2.05$ . With  $\lambda_c = 2\pi w/E_c$  one gets  $\lambda_c(TE_{10}) = 8.27\text{cm}$ ,  $\lambda_c(TE_{20}) = 5.35\text{ cm}$ ,  $f_c(TE_{10}) = 3.63\text{ GHz}$ ,  $f_c(TE_{20}) = 5.6\text{ GHz}$ . At 5.46 GHz one finds  $B = 2$  and  $K(TE_{10}) = 1.3$ , corresponding to  $\lambda_g = 2\pi w/K = 5.76\text{ cm}$ , no propagation for the  $TE_{20}$  mode.
- b) Guide WR 90, width 0.9 inches, one slab with  $\epsilon = 12.25$ , 0.135 inches thick, 0.18" between wall and center of slab. Wanted: cutoff frequencies of the two lowest order modes, guide wavelength at 10 GHz for the lowest order mode.  $w' = \text{width} = 0.9\text{ inch}$ ,  $\delta = 0.15$ ,  $\alpha + \delta/2 = 0.2$ .  $B_c(TE_{10}') = B_c(TE_{20})$



of guide with two slabs and with 1.8 inches and  $w = \text{width}/2$ . One finds  $B_c(\text{TE}'_{10}) = 1.83$ ,  $B_c(\text{TE}'_{20}) = B_c(\text{TE}_{40}) \approx 4.25$ ; with  $\lambda_c = 2\pi w'/B_c$  one gets  $\lambda_c(\text{TE}'_{10}) \approx 7.85$  cm,  $\lambda_c(\text{TE}'_{20}) \approx 3.38$  cm and  $f_c(\text{TE}'_{10}) \approx 3.82$  GHz,  $f_c(\text{TE}'_{20}) \approx 7.83$  GHz. At 10 GHz one finds  $B = 2\pi w'/\lambda_0 \approx 4.8$  and  $K \approx 12.71$ , yielding  $\lambda_g \approx 1.12$  cm.

c) The  $\text{TE}_{20}$  solutions with  $\alpha + \delta/2 = 0.5$  are equivalent to the  $\text{TE}_{10}$  solutions with the two slabs touching each other, where  $\alpha + \delta/2 = 1 - \delta/2$ , i.e.  $2K(\text{TE}_{10}, B, \alpha + \delta/2 = 1 - \delta/2) = K(\text{TE}_{20}, 2B, \alpha + \delta/2 = 0.5)$   
 $2B_c(\text{TE}_{10}, \alpha + \delta/2 = 1 - \delta/2) = B_c(\text{TE}_{20}, \alpha + \delta/2 = 0.5)$

### DISCUSSION OF RESULTS

The influence of the dielectric slab on the field distribution in the waveguide is to concentrate within the slab with increasing frequency an increasing fraction of the total energy flowing through the guide. The phase velocity approaches asymptotically " $\sqrt{\epsilon}$  times the velocity of light in free space". The parametric dependence of the E field distribution is illustrated by Figure 3.

The influence of a thin slab on the cutoff frequencies is the stronger the closer the slab position is to the lines of maximum electric field strength in the empty guide. With increasing slab thickness the influence is less and less related to the empty guide field distribution. For thick slabs it is weakest when the slabs touch the guide walls; for odd modes it is strongest when the slabs are in the center of the guide and touch each other, for even modes it is strongest when each slab is in center of half a guide width.

At a certain thickness the cutoff frequency is nearly independent of the slab position (see  $\text{TE}_{30}, 40, 60$ ). For thinner slabs a similar behavior can be found in the propagation characteristics. Figure 4 shows  $K$  versus  $B$  curves for the  $\text{TE}_{40}$  mode with  $\epsilon = 9$  and  $\delta = 0.2$ . At  $B = 5.42$  the propagation constant is independent of the slab position. Above this crossover frequency the same order of curves is obtained as for cutoff frequencies

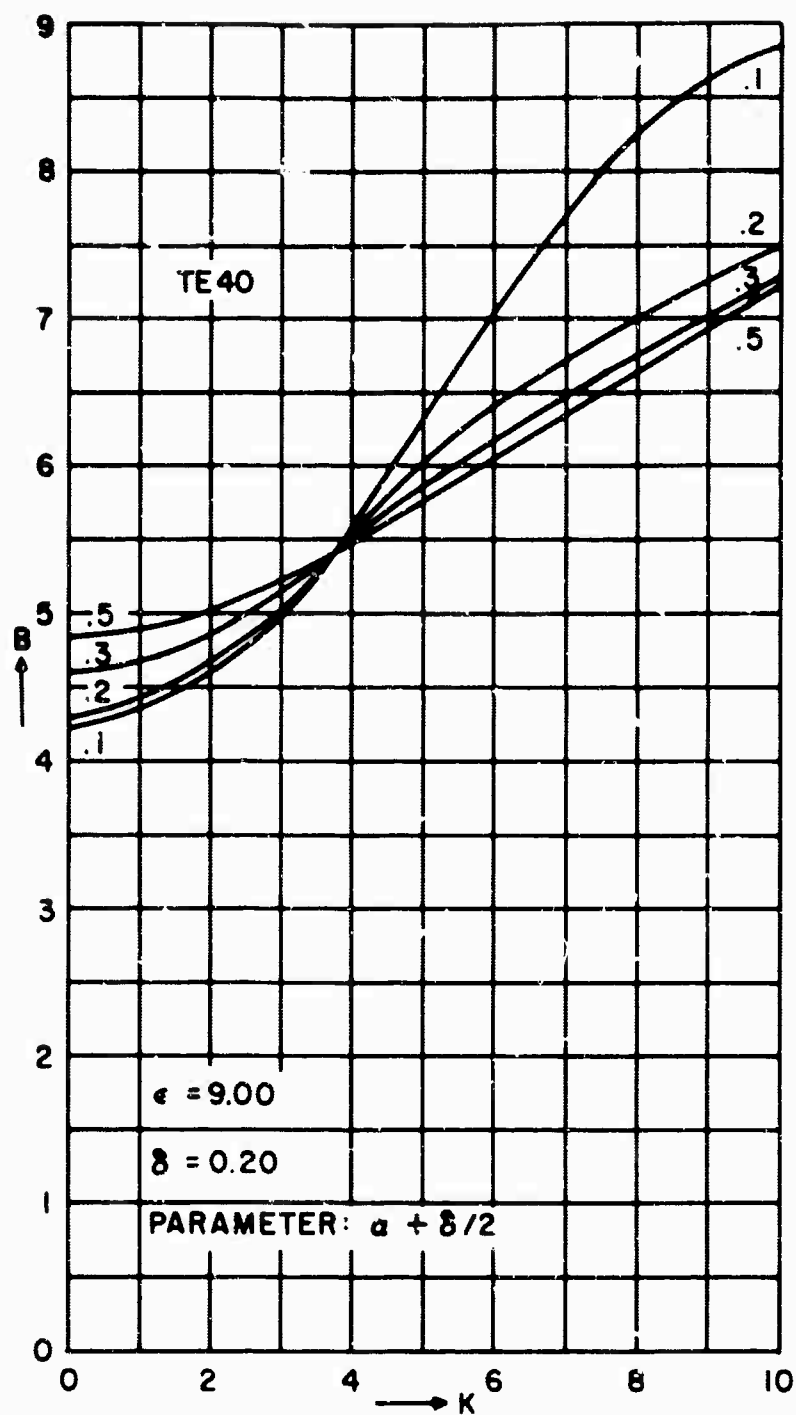


Figure 3. Propagation Characteristics in Waveguide with Dielectric Slabs

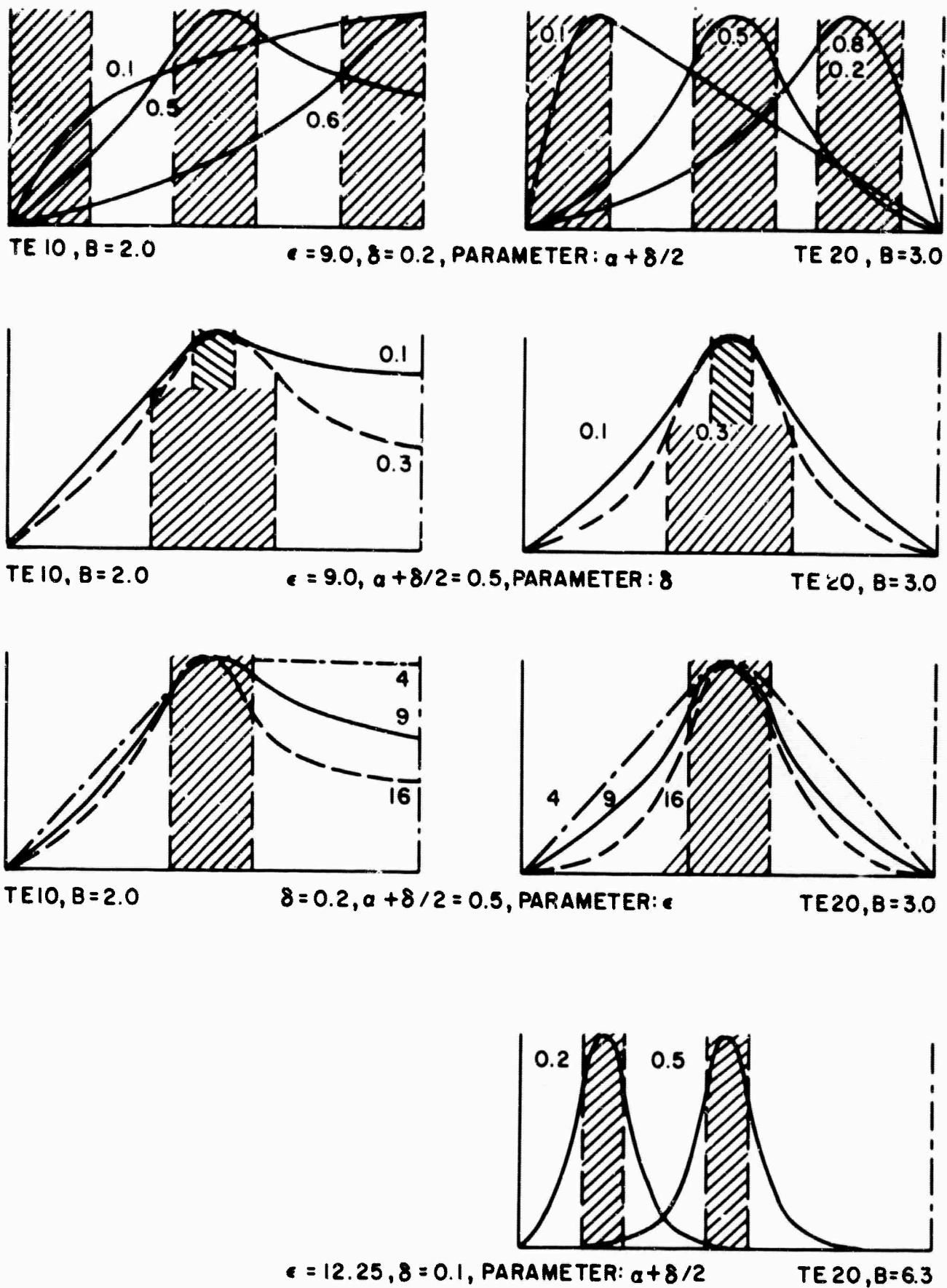


Figure 4. E-Field Distribution in a Waveguide with Dielectric Slabs

with thick slabs. The propagation characteristics for  $TE_{20}$  modes for e.g.  $\epsilon = 12.25$  and  $\delta = 0.1$  show, however, that this order, too, may change at even higher frequencies. If the empty regions in half the waveguide (I and III,  $\alpha$  and  $\gamma$  in Figure 1) are not equal, a wave may treat them at high frequencies as essentially equal with the value of the narrower region. The waveguide shows then a larger, propagation constant and field concentration at a given frequency than a guide with  $\alpha = \gamma$  does. An E-field distribution for such a case is shown in the lowest part of Figure 3.

The ratio of the magnetic field components, the ellipticity,

$$(15) \quad ELL(\varphi) = H_x(\varphi)/jH_y(\varphi)$$

at the slab boundaries I/II and II/III approaches "one" asymptotically, as illustrated in Figure 5 for  $TE_{10}$  and  $TE_{20}$  modes for various parameters. How the ratios of cutoff frequencies, the fractional bandwidths

$$(16) \quad \alpha_{21} = \frac{B_c(TE_{20})}{B_c(TE_{10})} \quad \text{and} \quad \alpha_{42} = \frac{B_c(TE_{40})}{B_c(TE_{20})}$$

vary with different parameters is shown in Figure 6. Finally, the parametric dependance of the normalized wave admittance

$$(17) \quad Y_w/Y_o = (H_x/E_z)/Y_o = K/B = \lambda_o/\lambda_g \quad \text{with} \quad Y_o = \sqrt{\epsilon_o/\mu_o} = (377 \text{ ohms})^{-1}$$

is shown in Figure 7.

#### ACKNOWLEDGEMENTS

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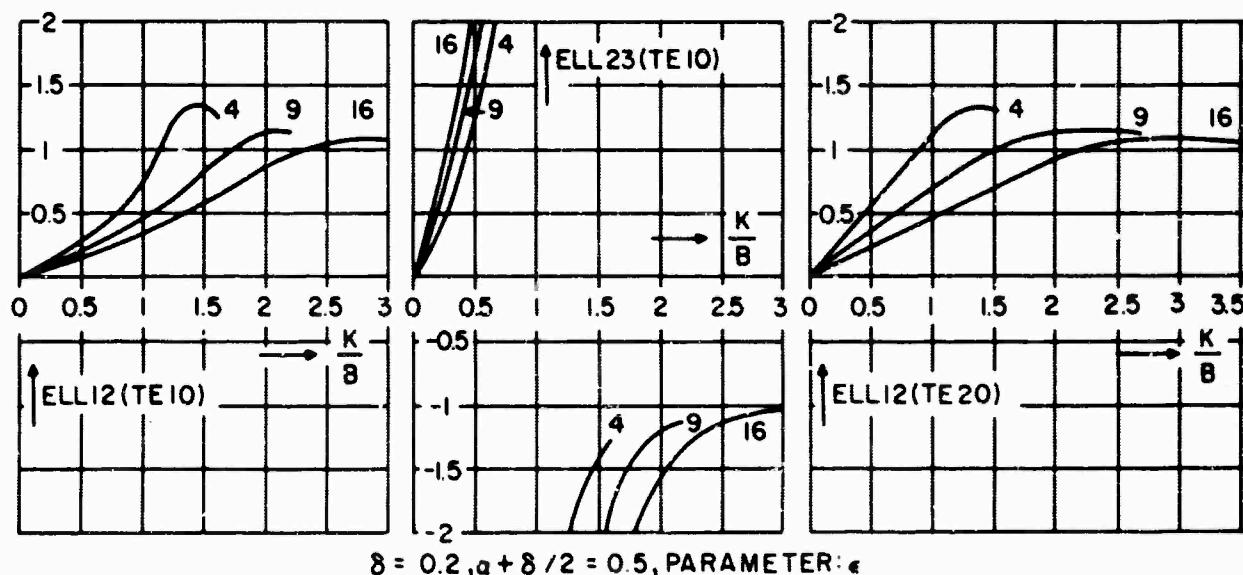
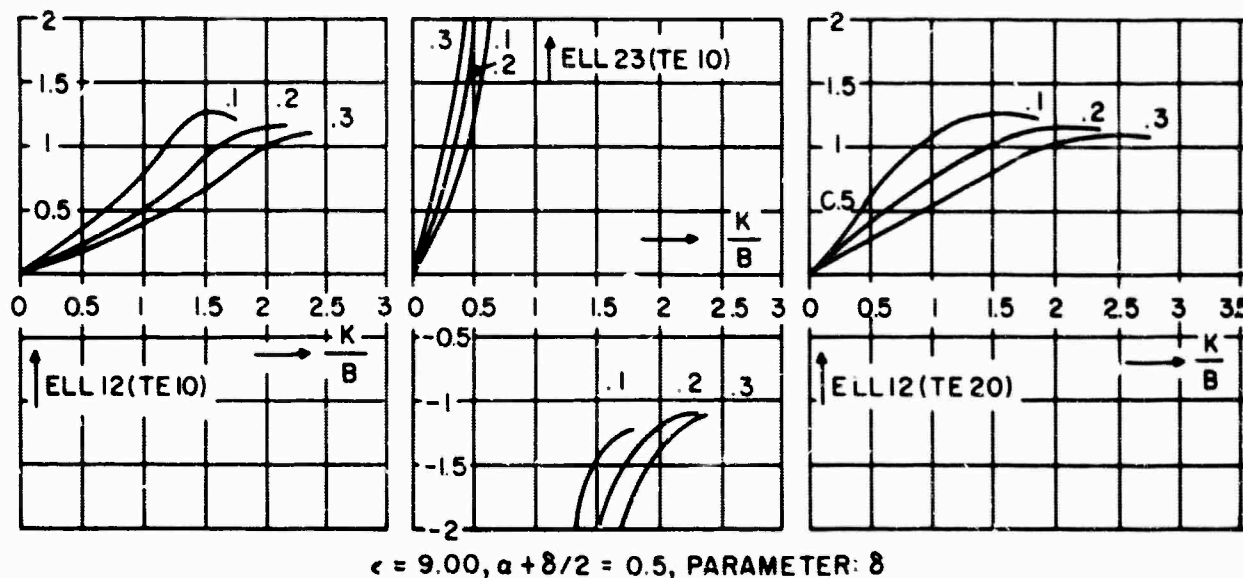
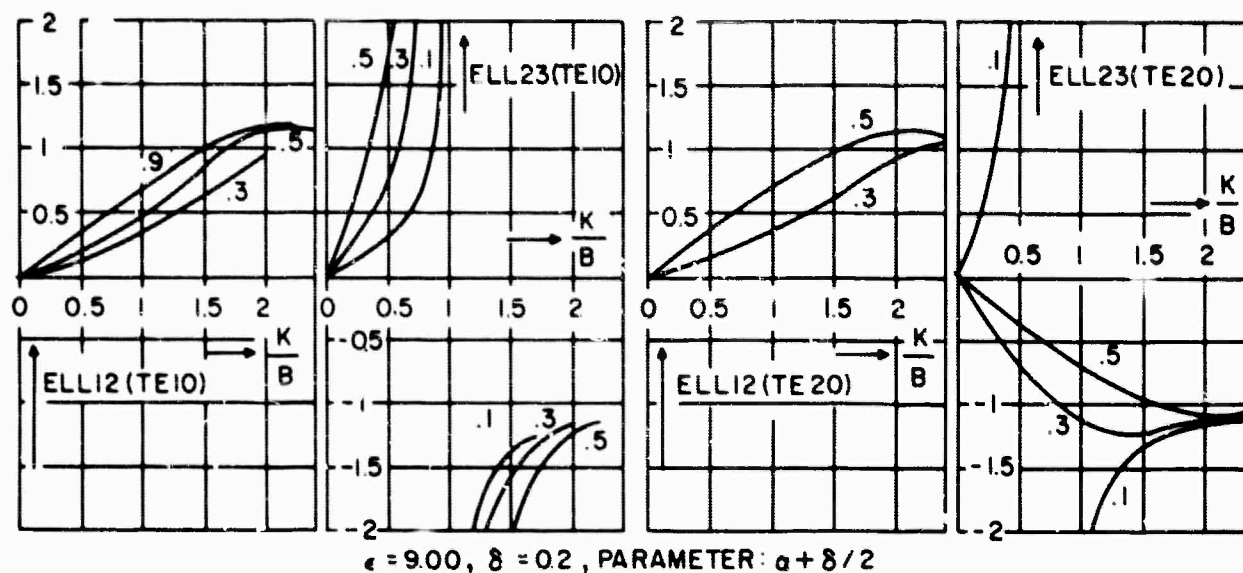


Figure 5. Ellipticity at Slab Boundaries in Waveguide with Dielectric Slabs

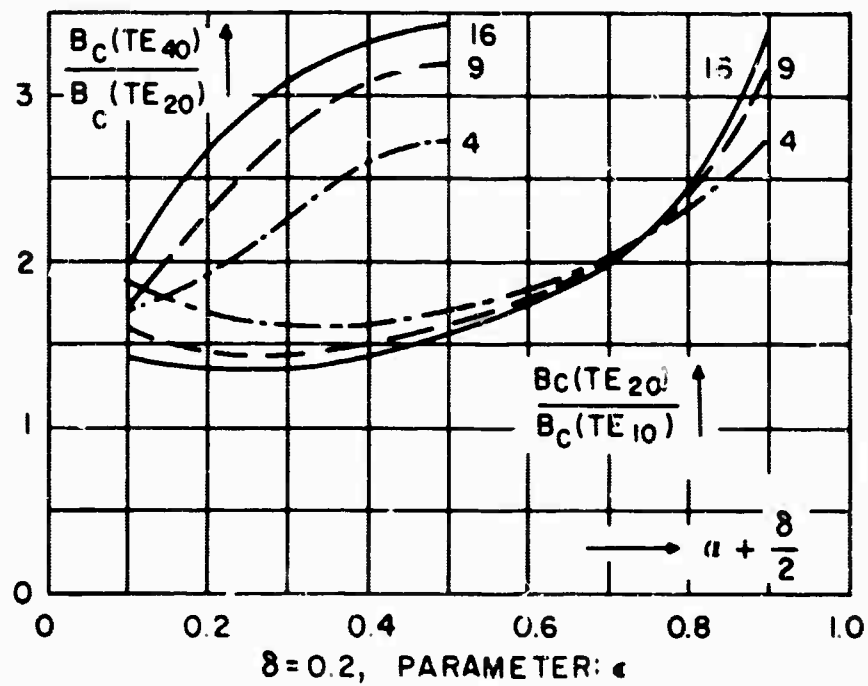
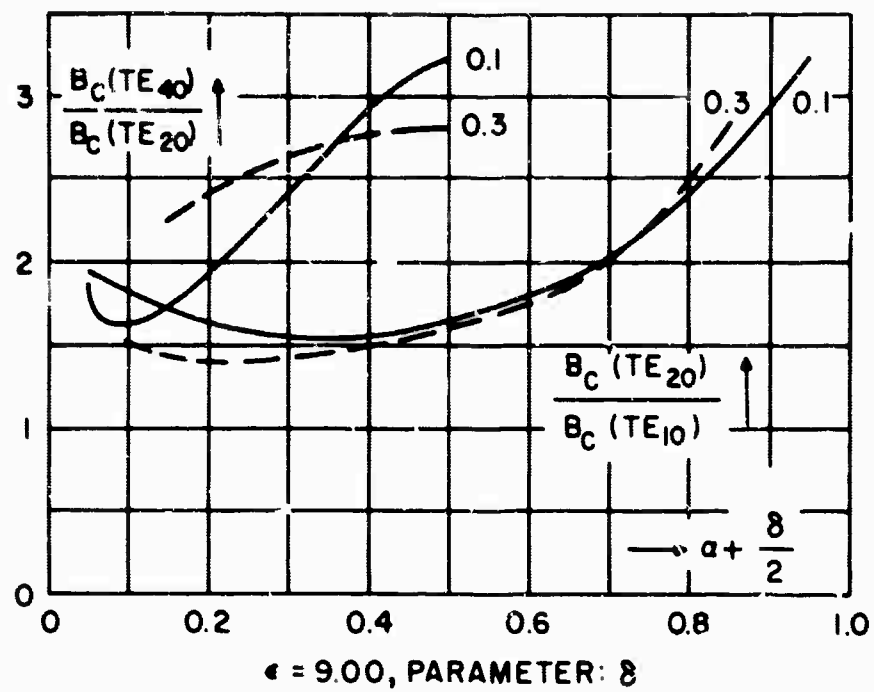


Figure 6. Fractional Bandwidths of Waveguide with Dielectric Slabs

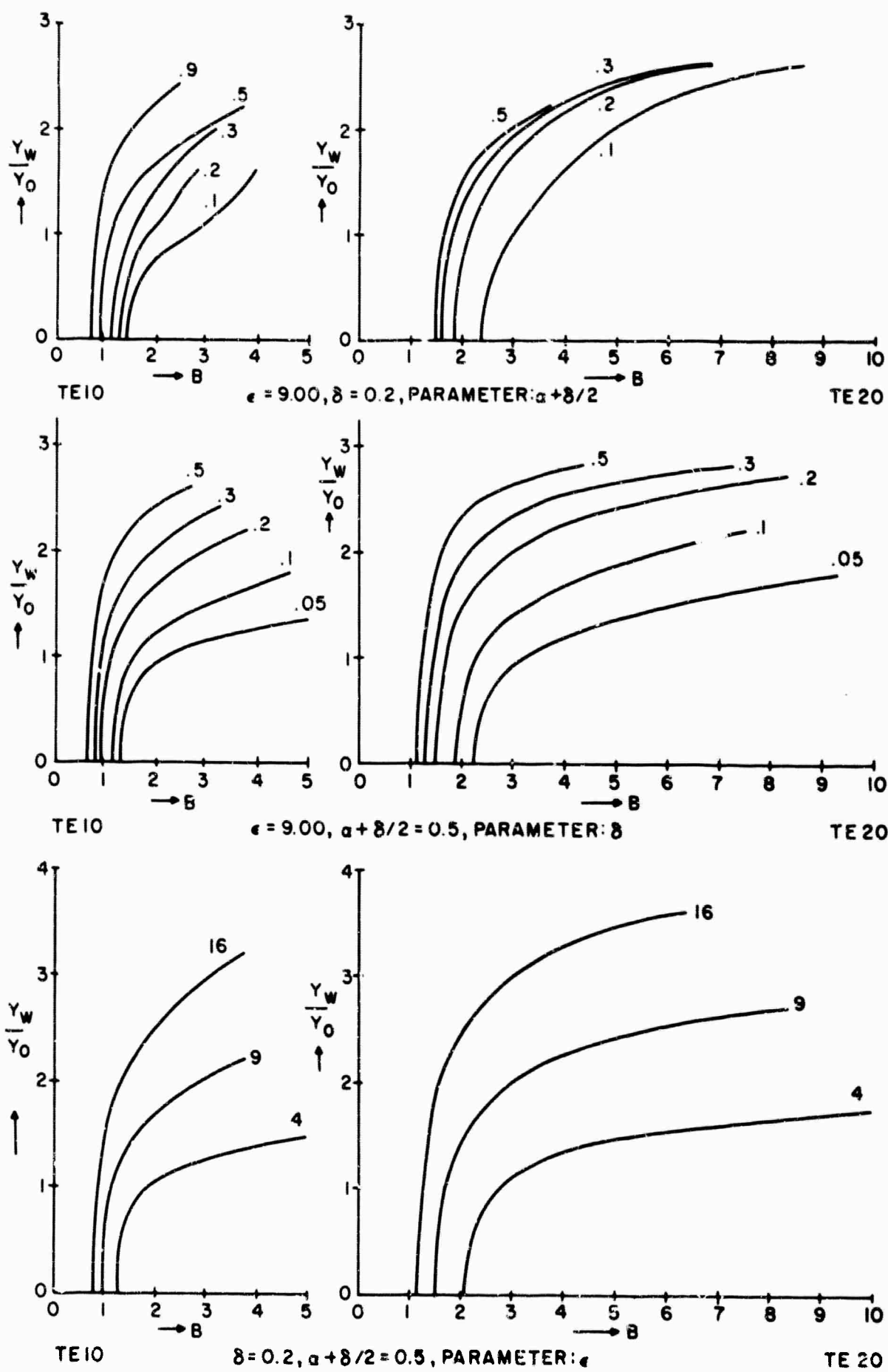


Figure 7. Normalized Wave Admittance of Waveguide with Dielectric Slabs

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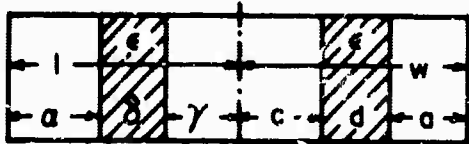


# NORMALIZED CUTOFF-FREQ.

$B_c = \beta_c w$  FOR  $TE_{n0}$ -MODES

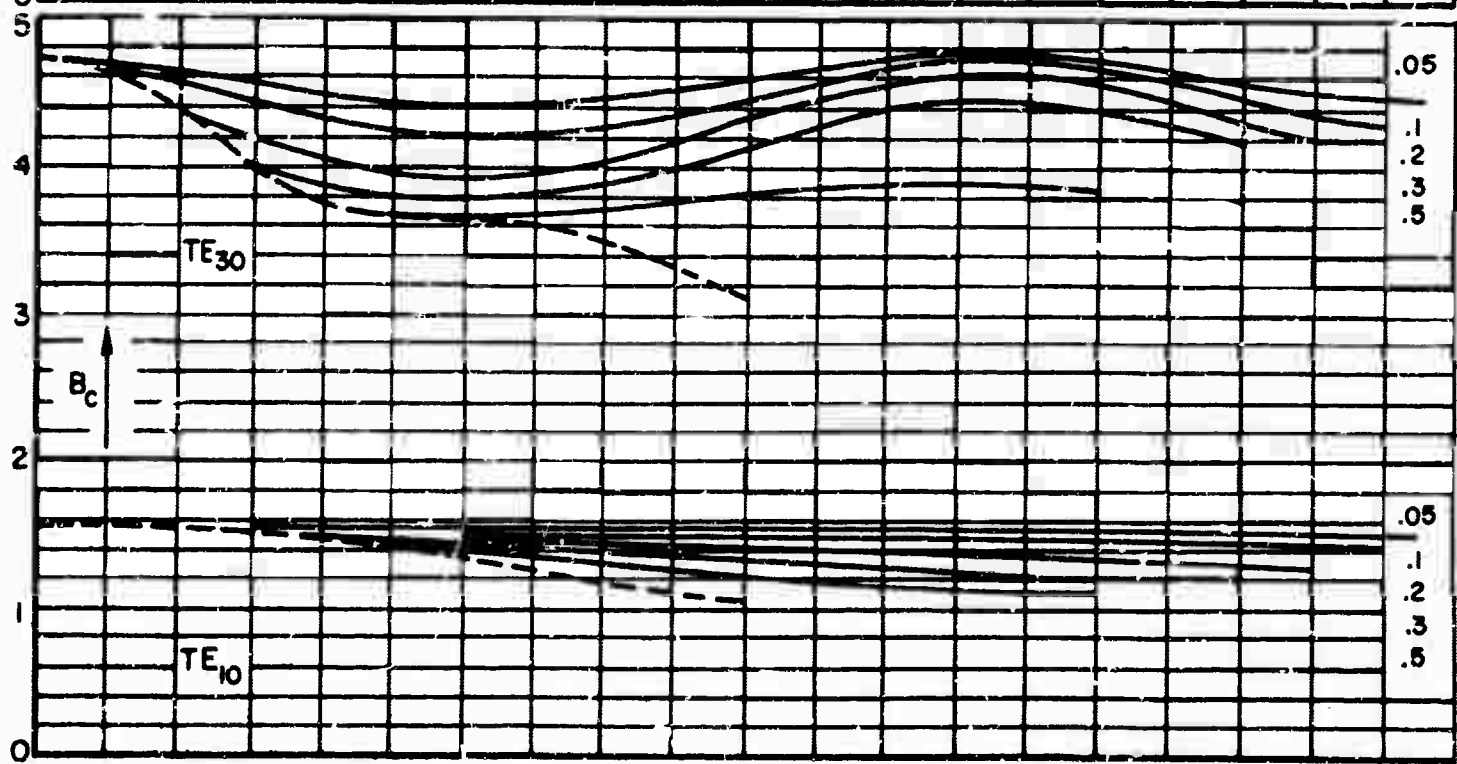
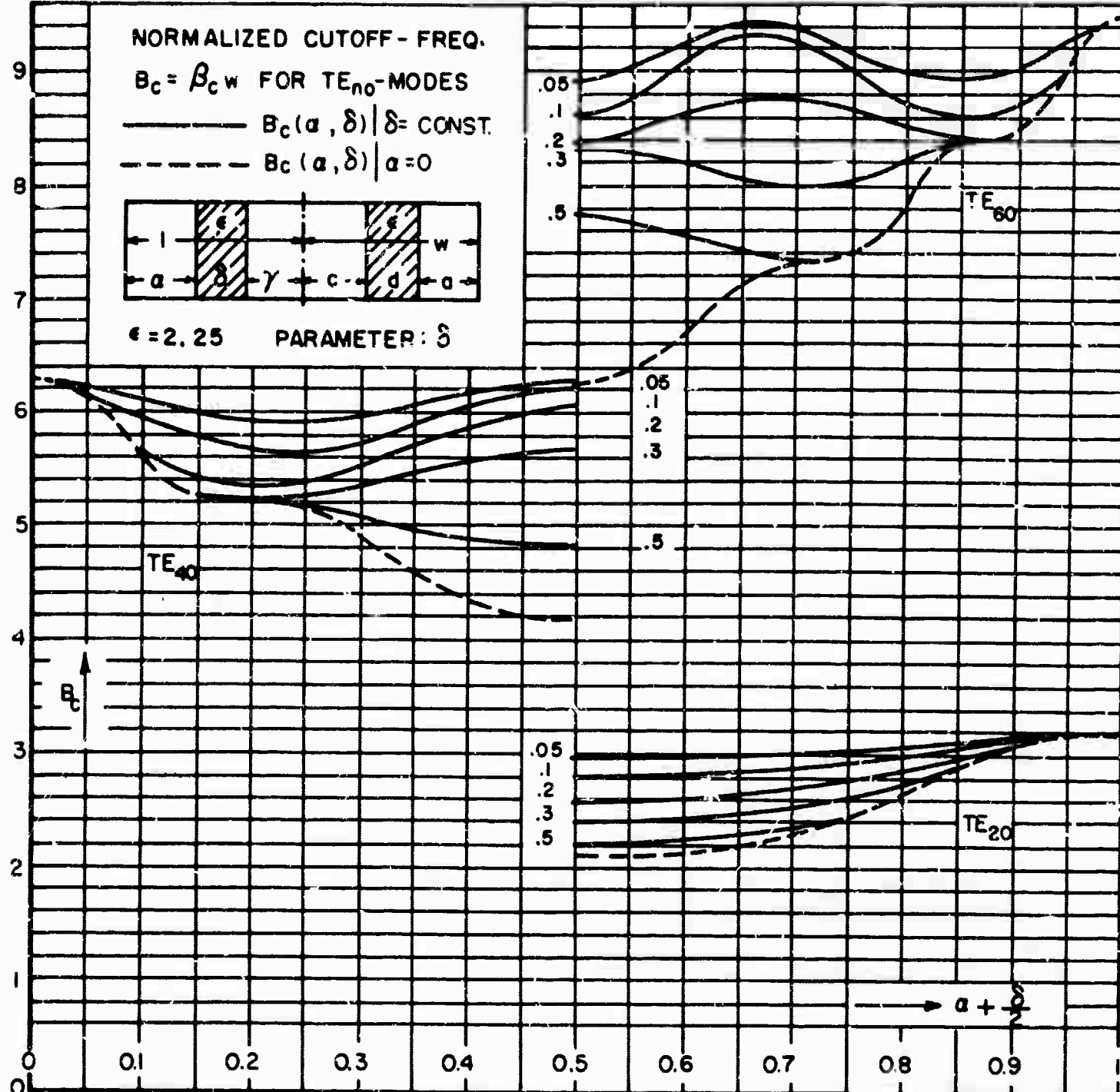
————  $B_c(a, \delta) | \delta = \text{CONST.}$

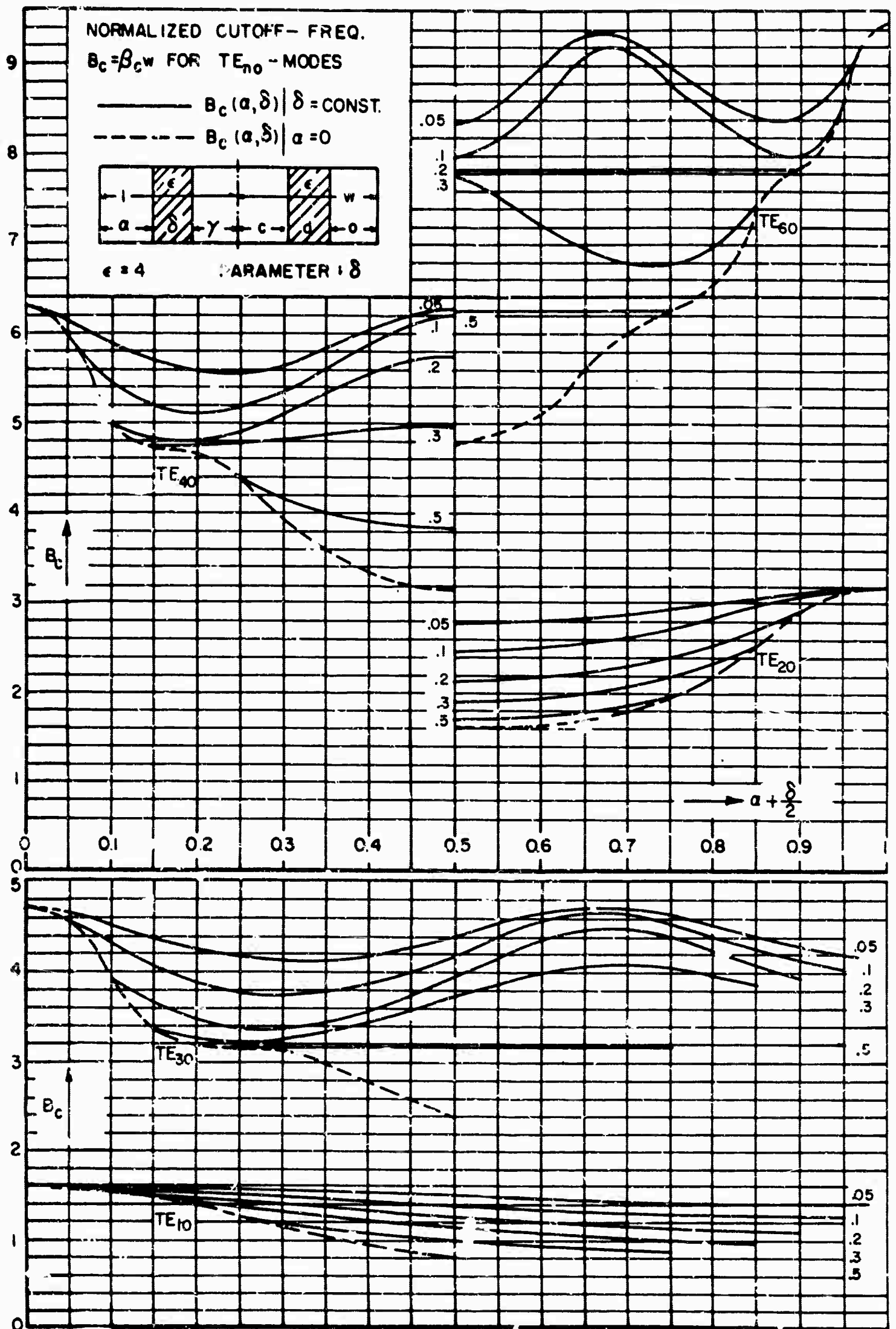
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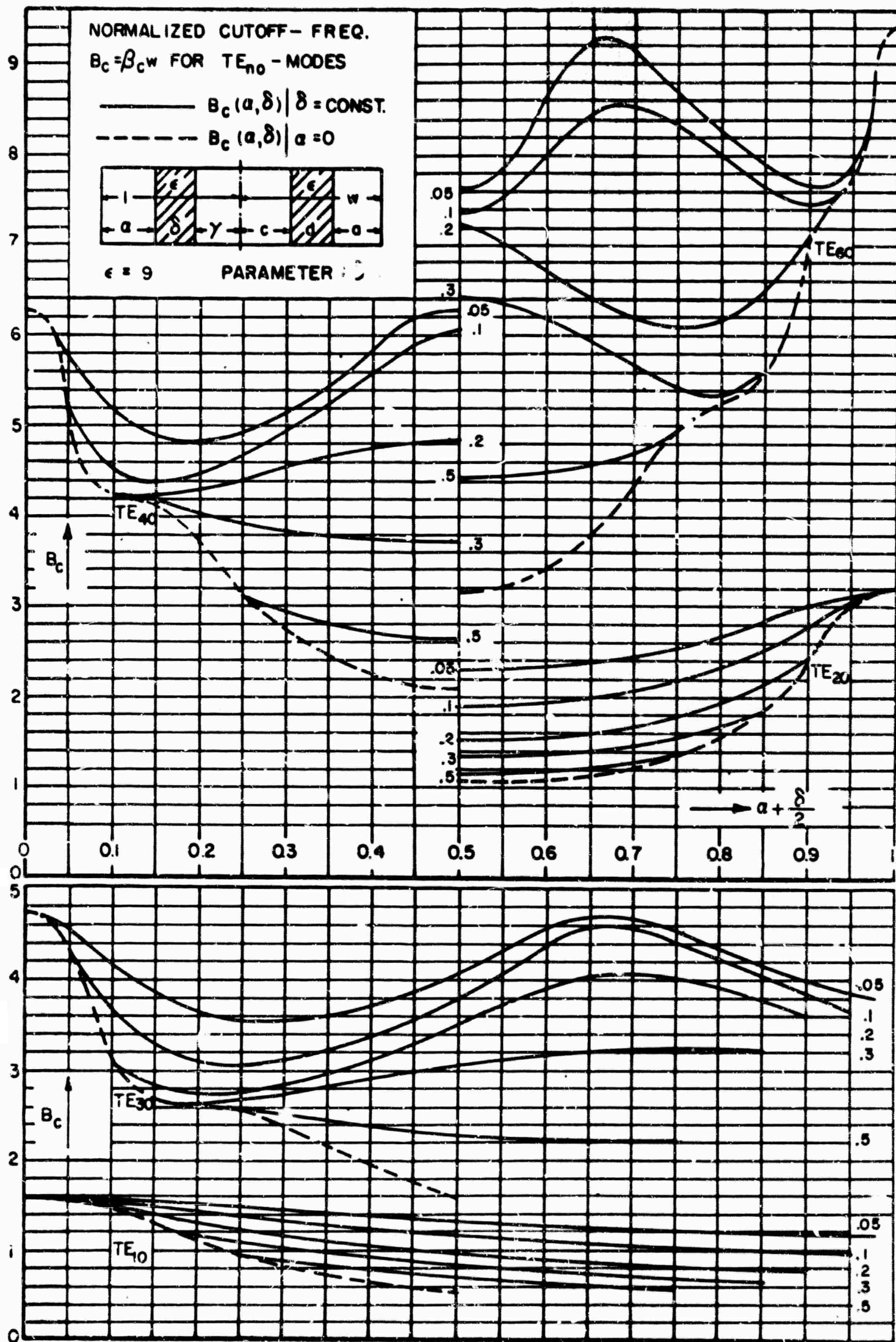


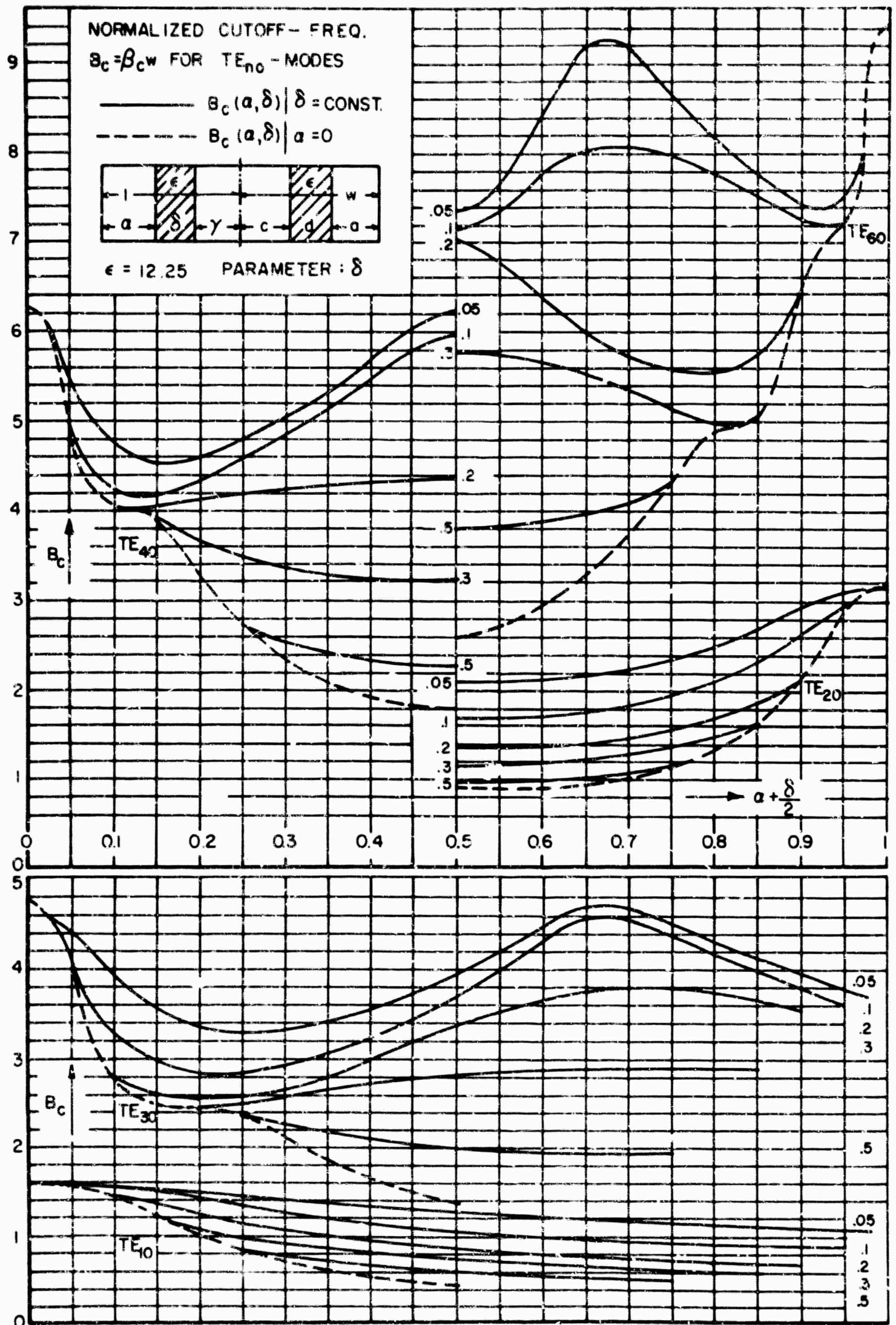
$\epsilon = 2.25$

PARAMETER:  $\delta$

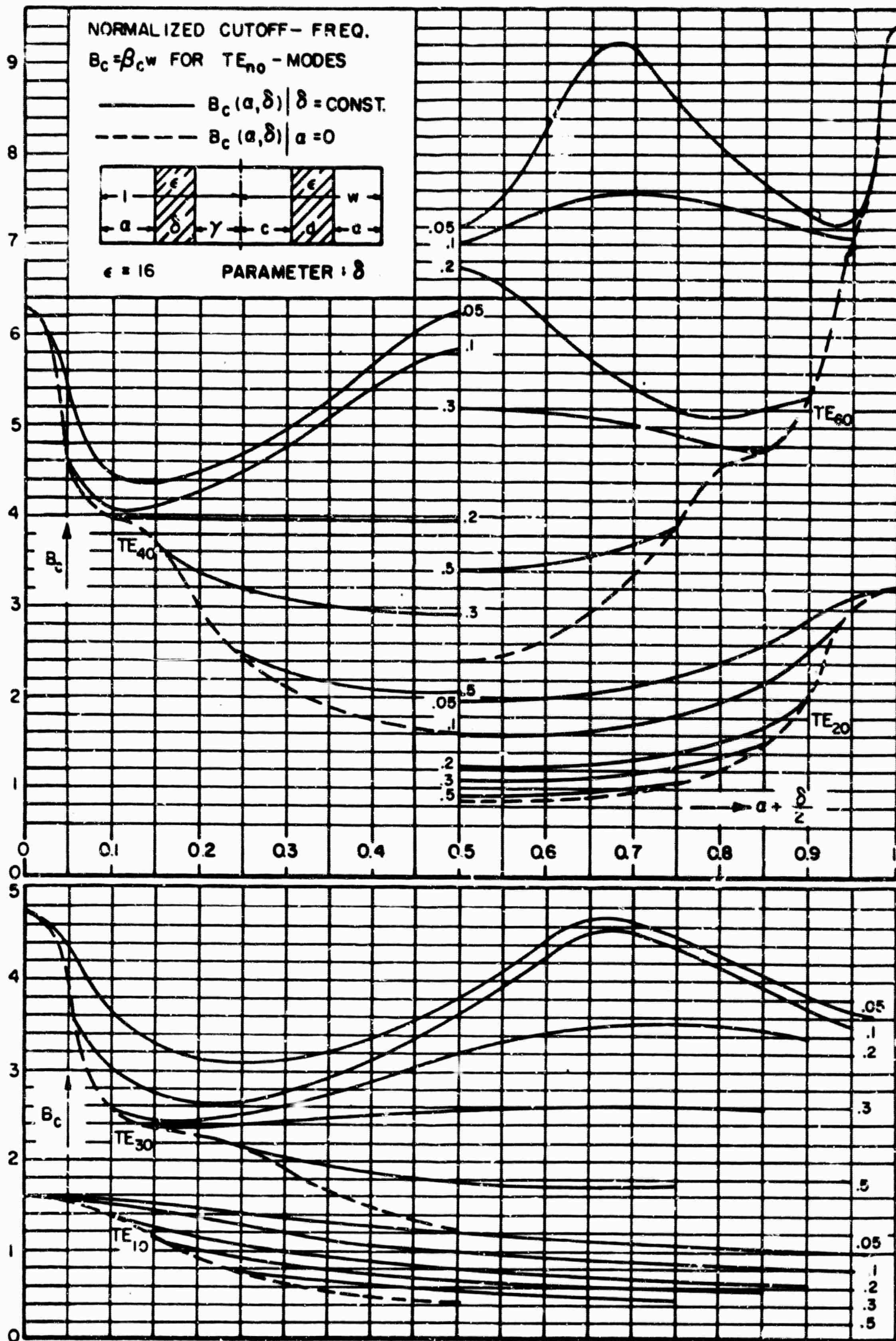


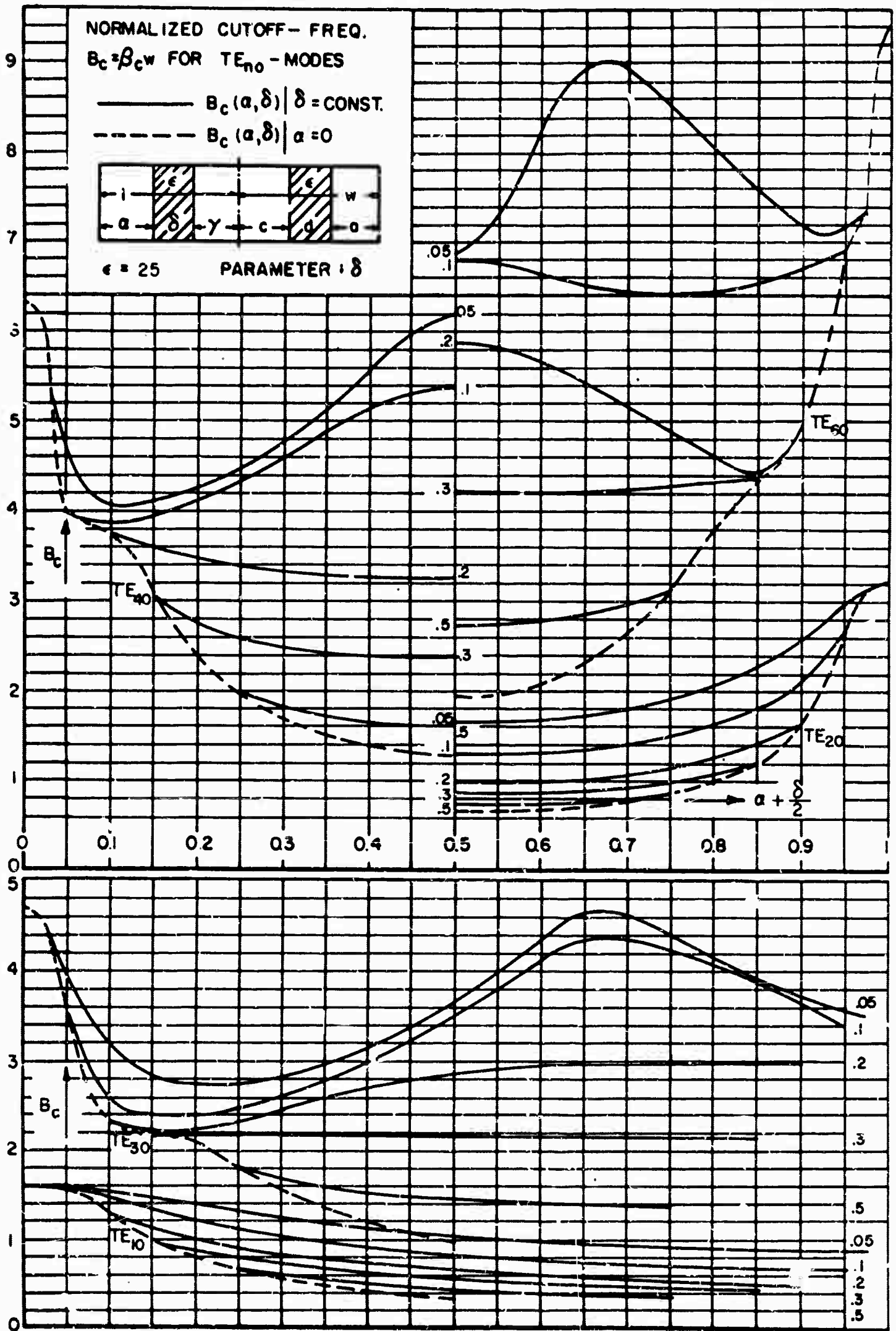


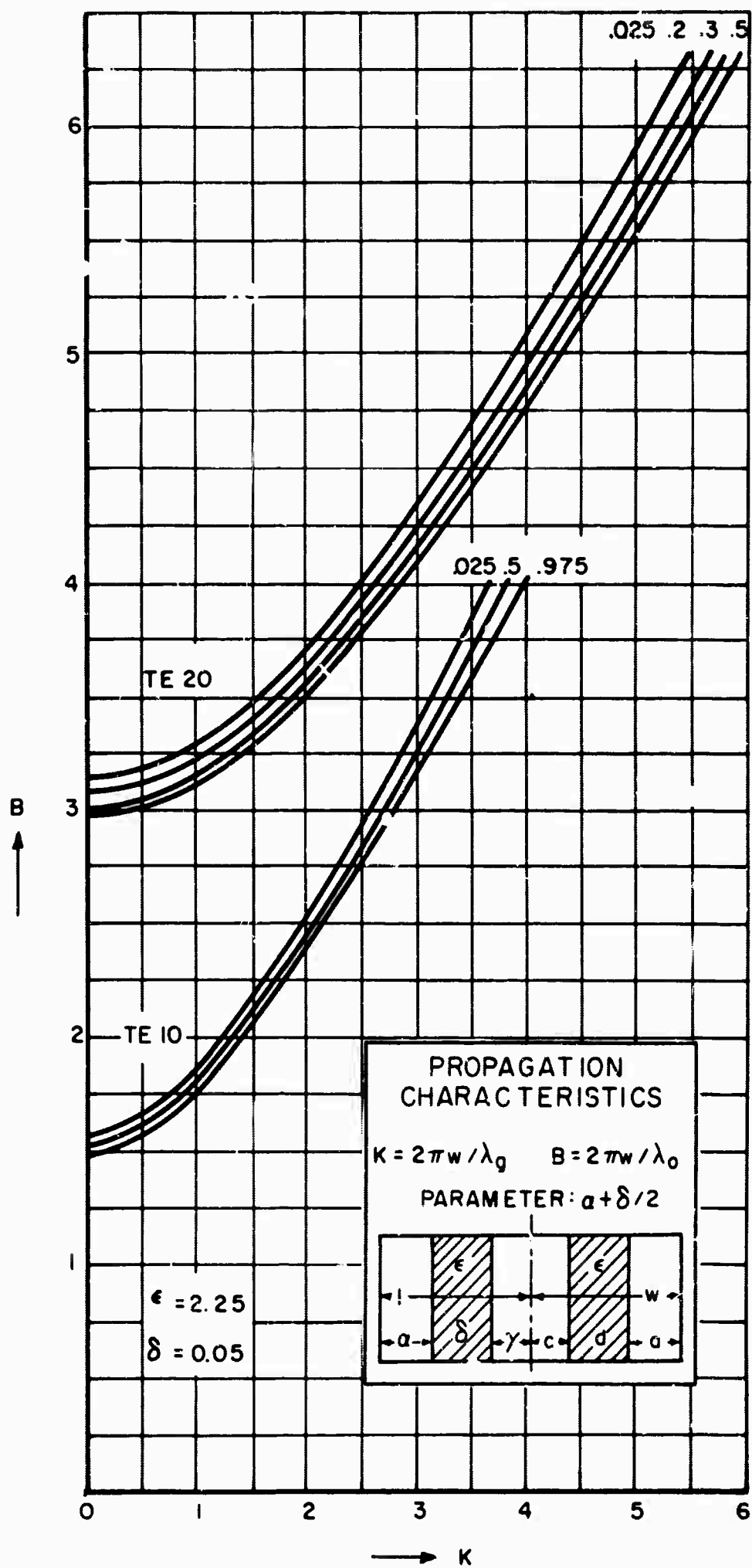


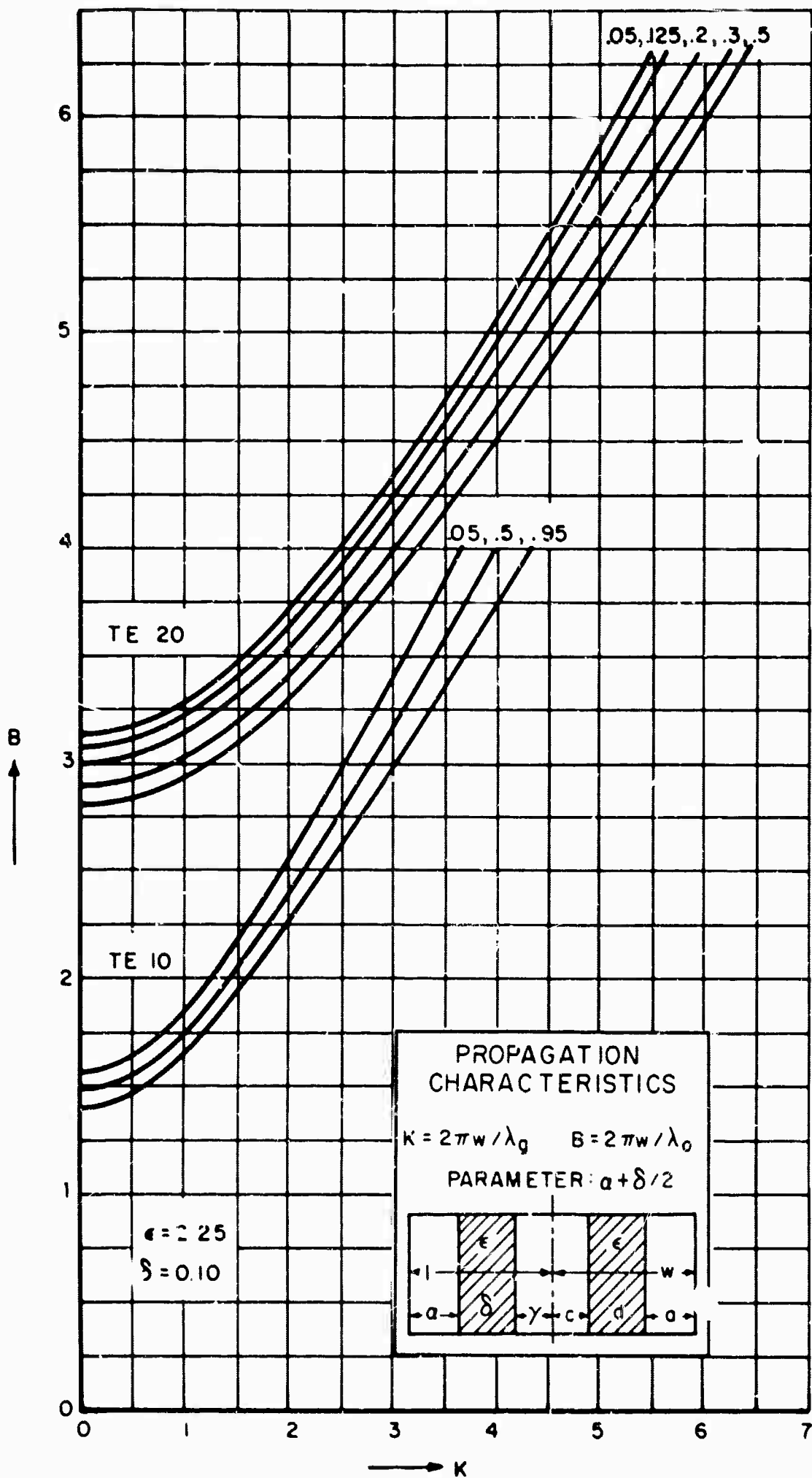




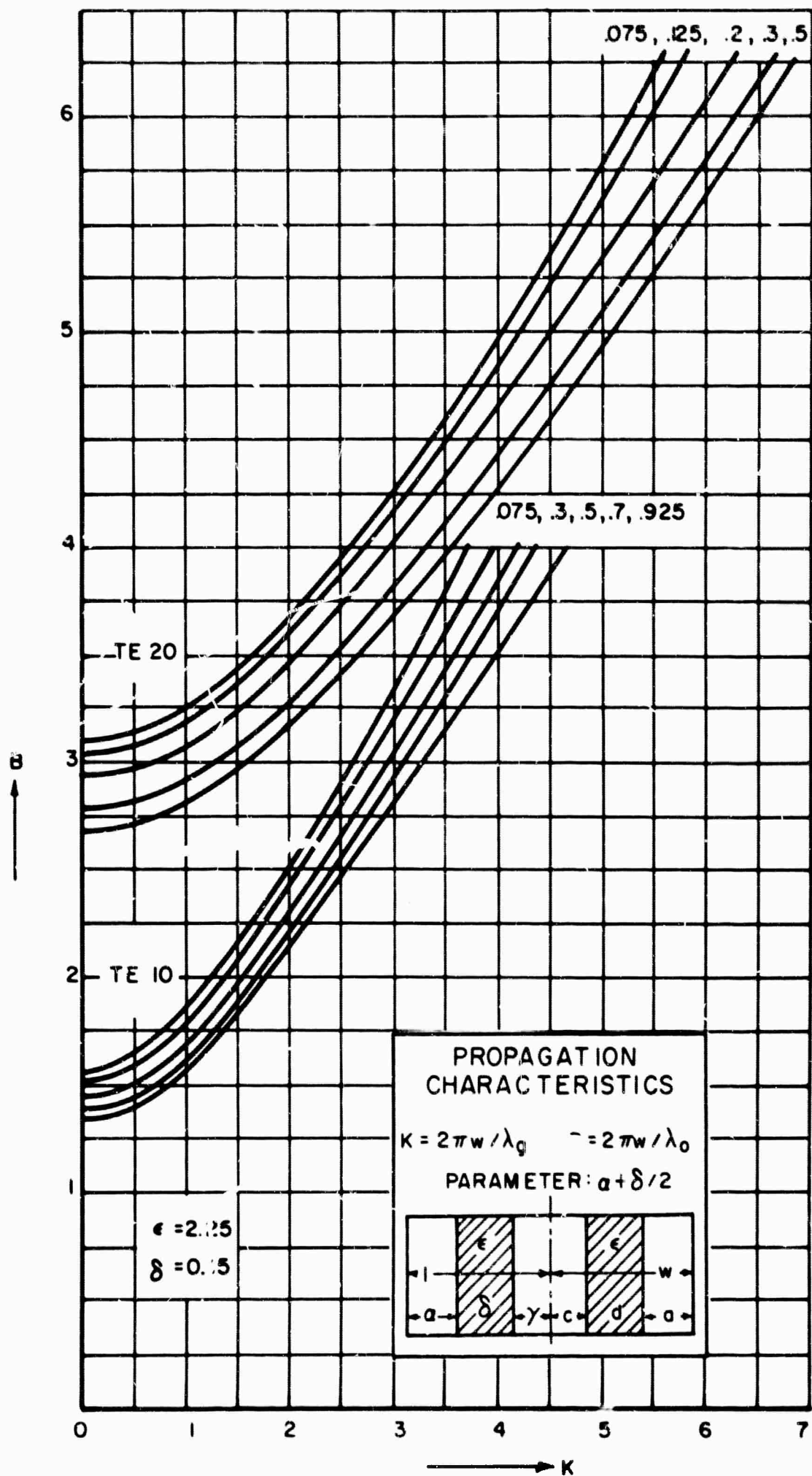


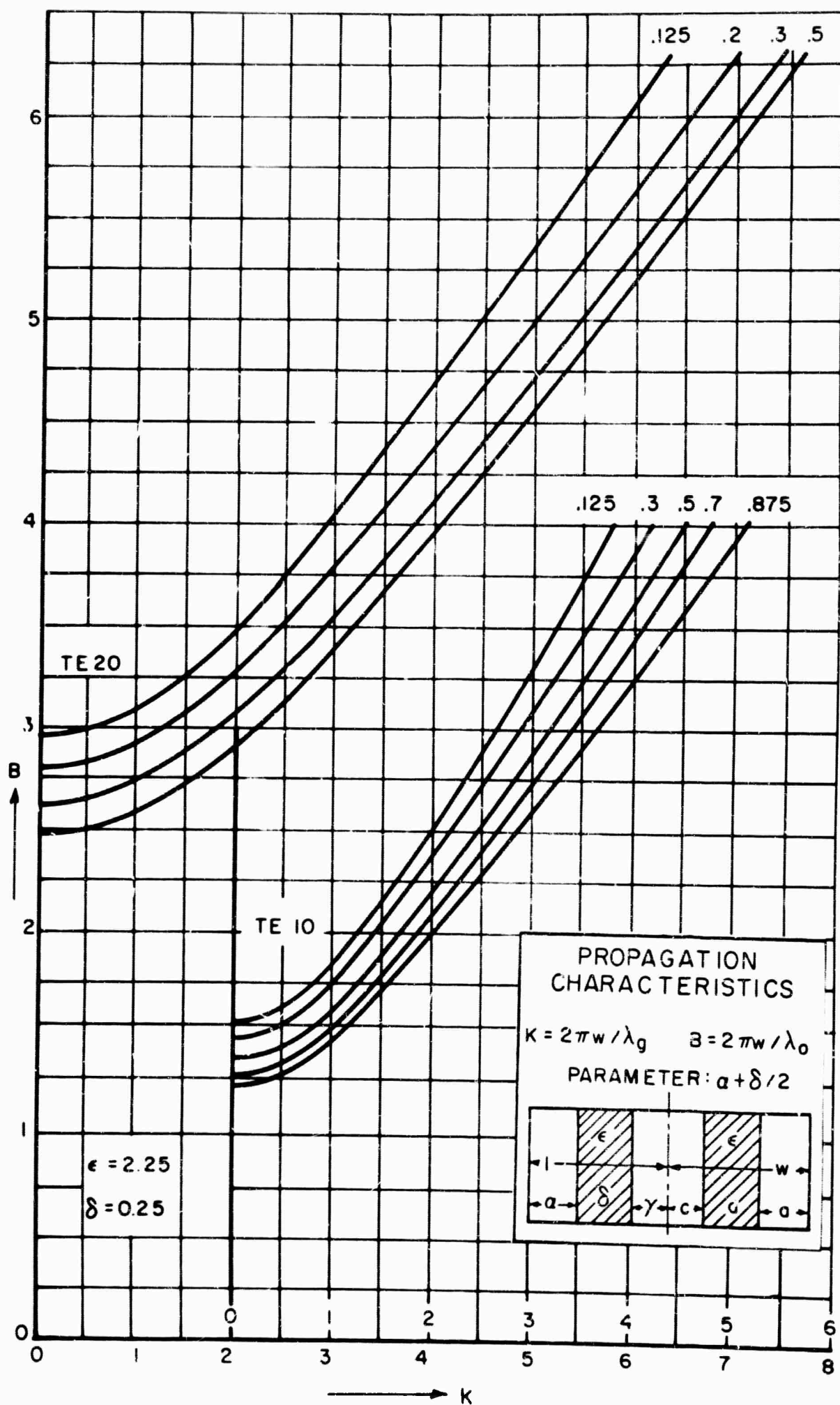


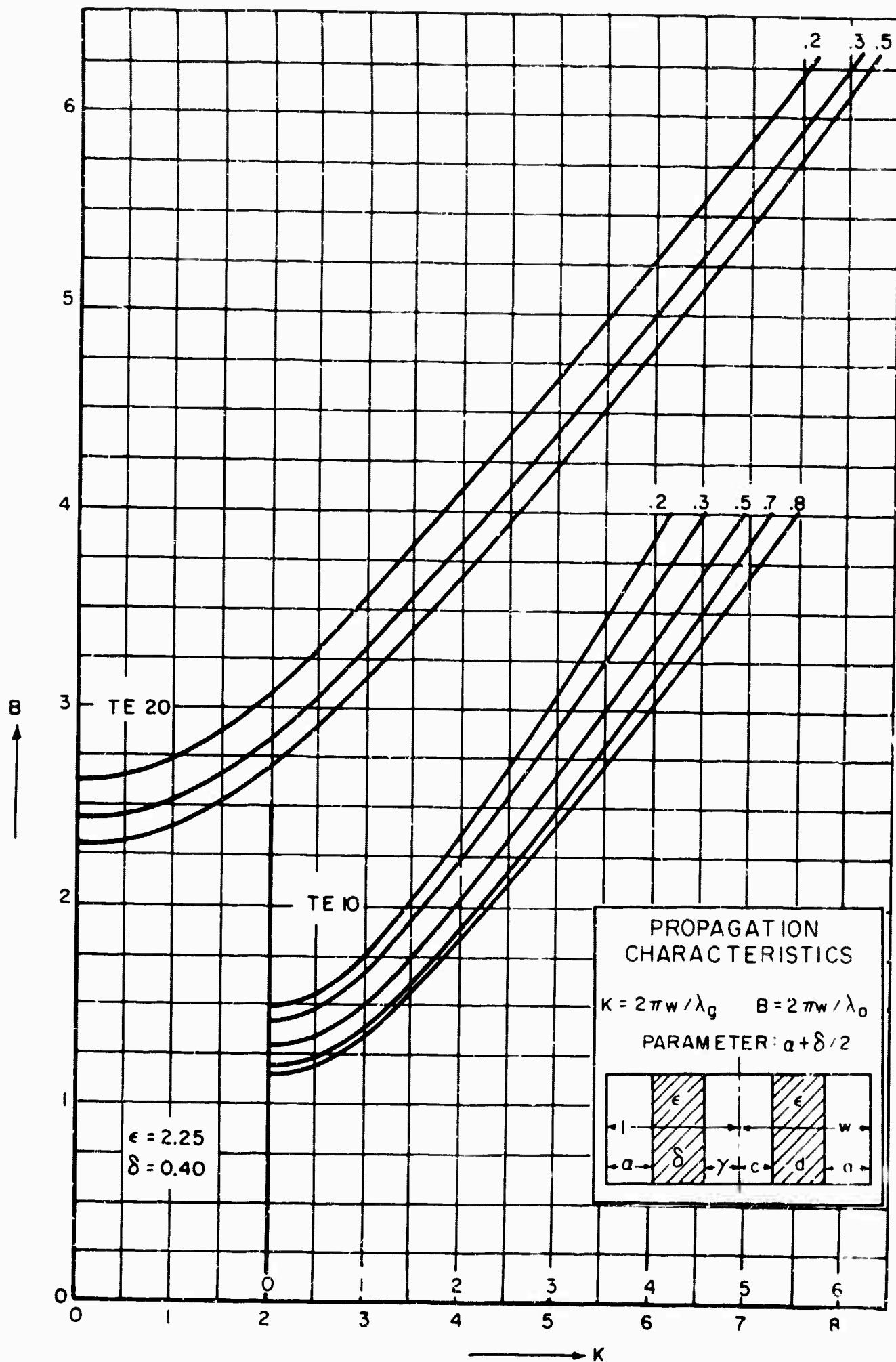


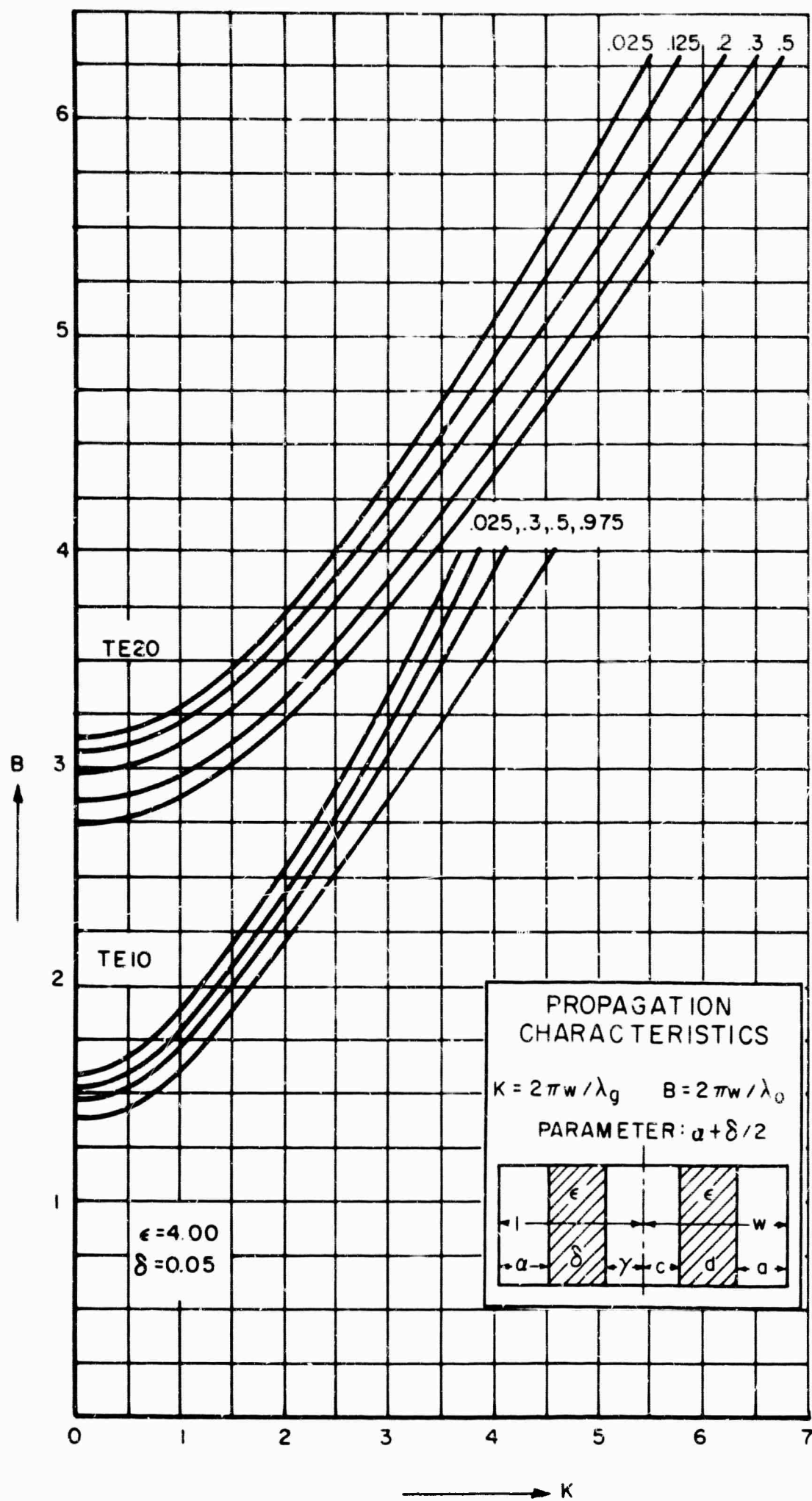


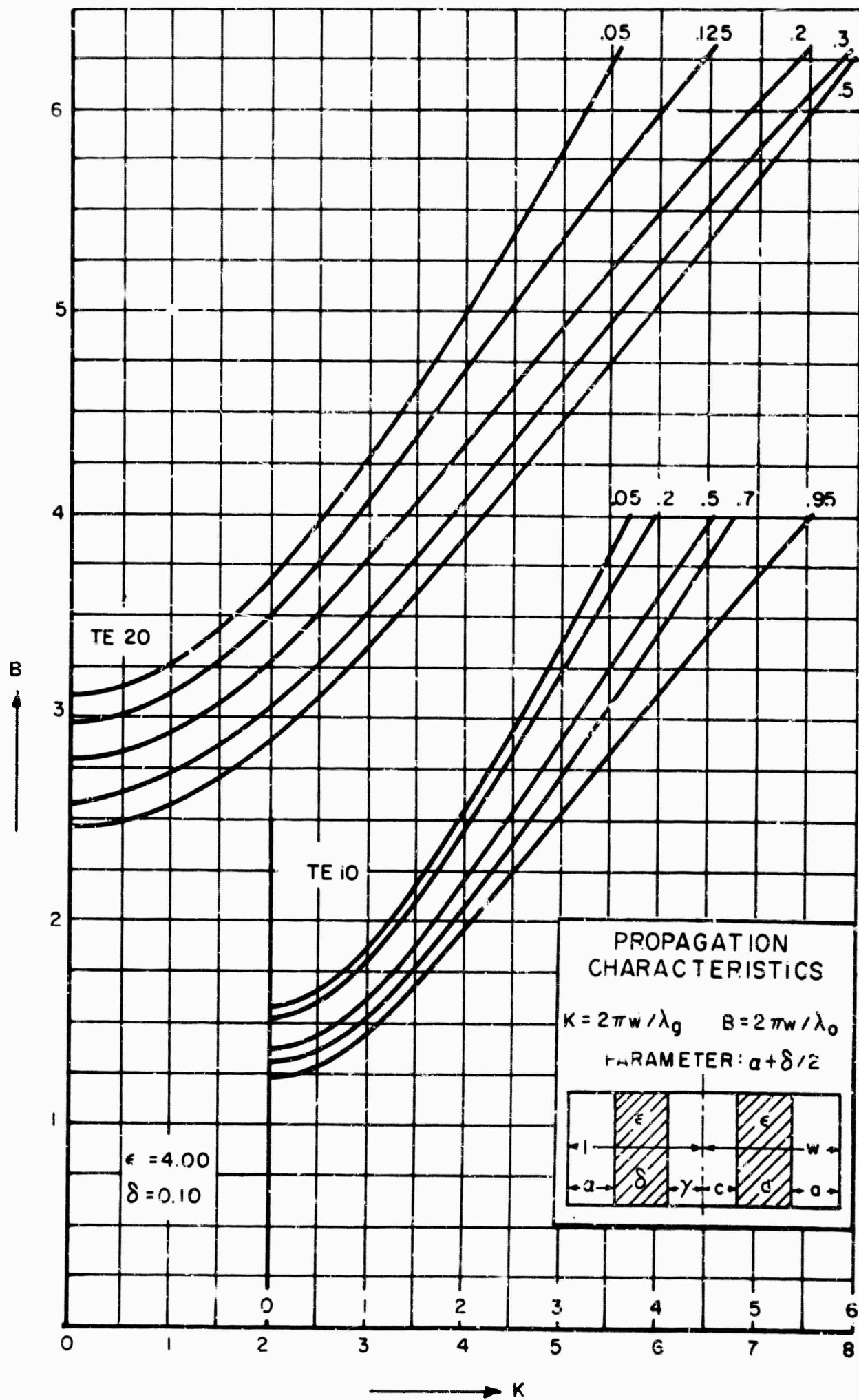


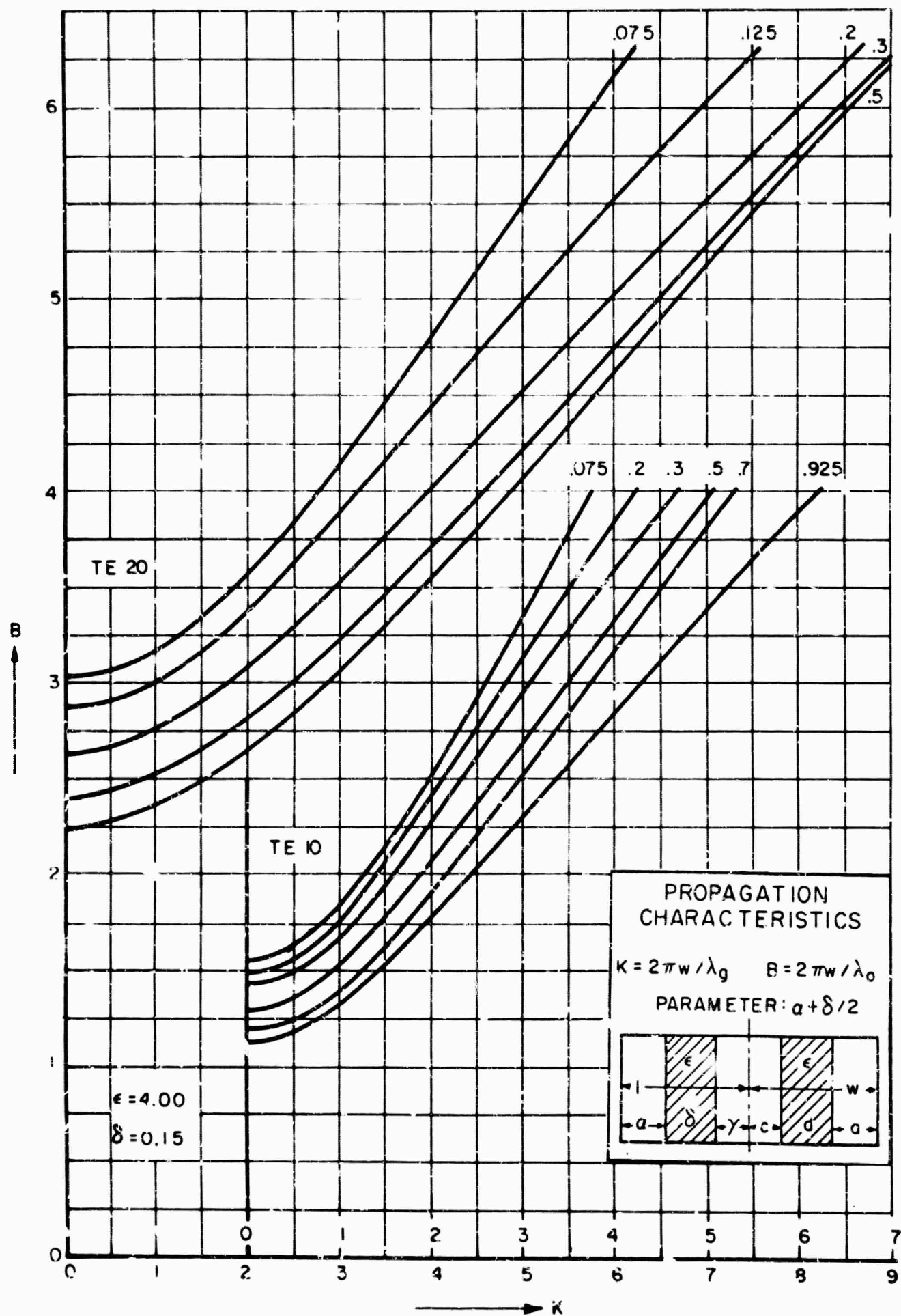




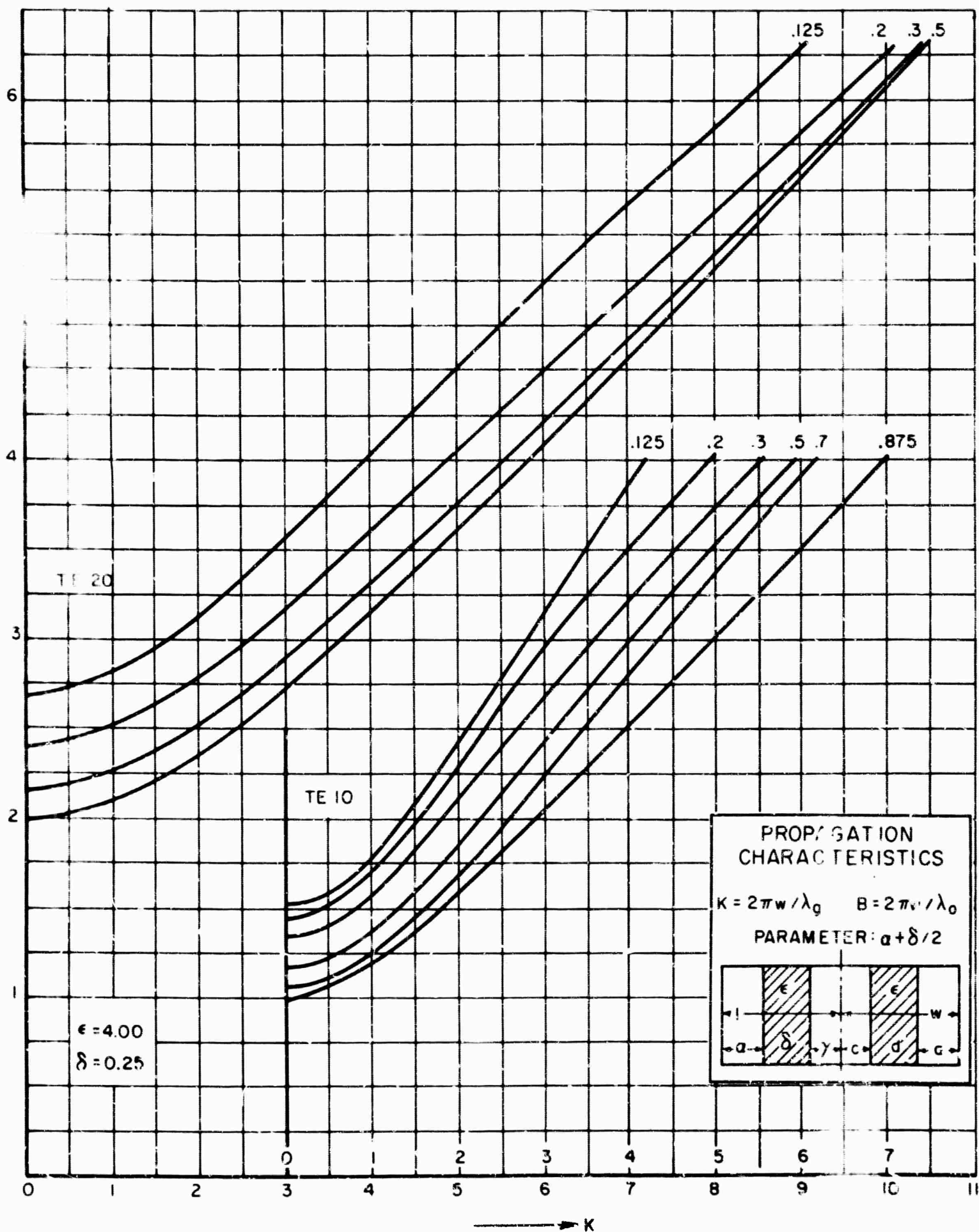


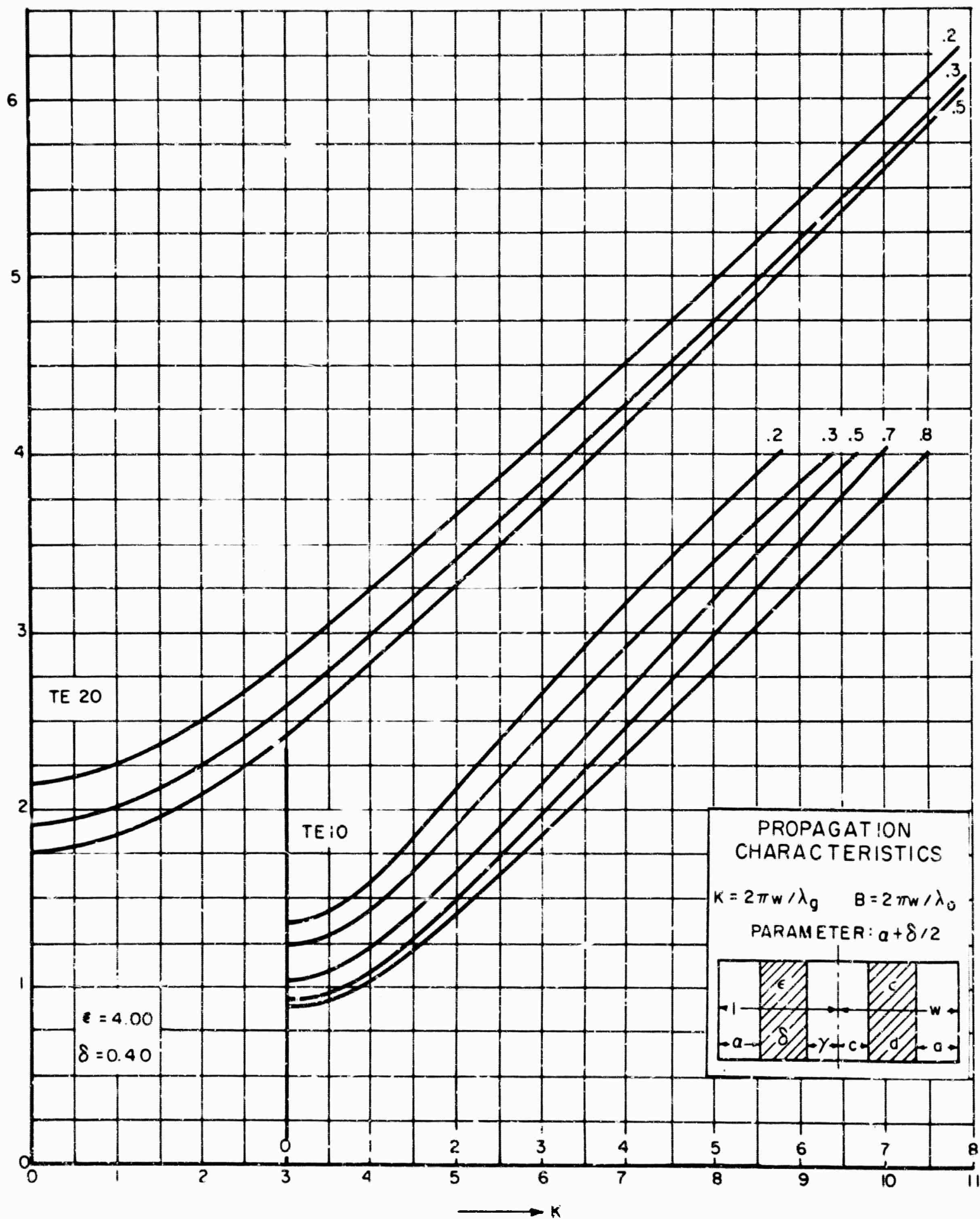




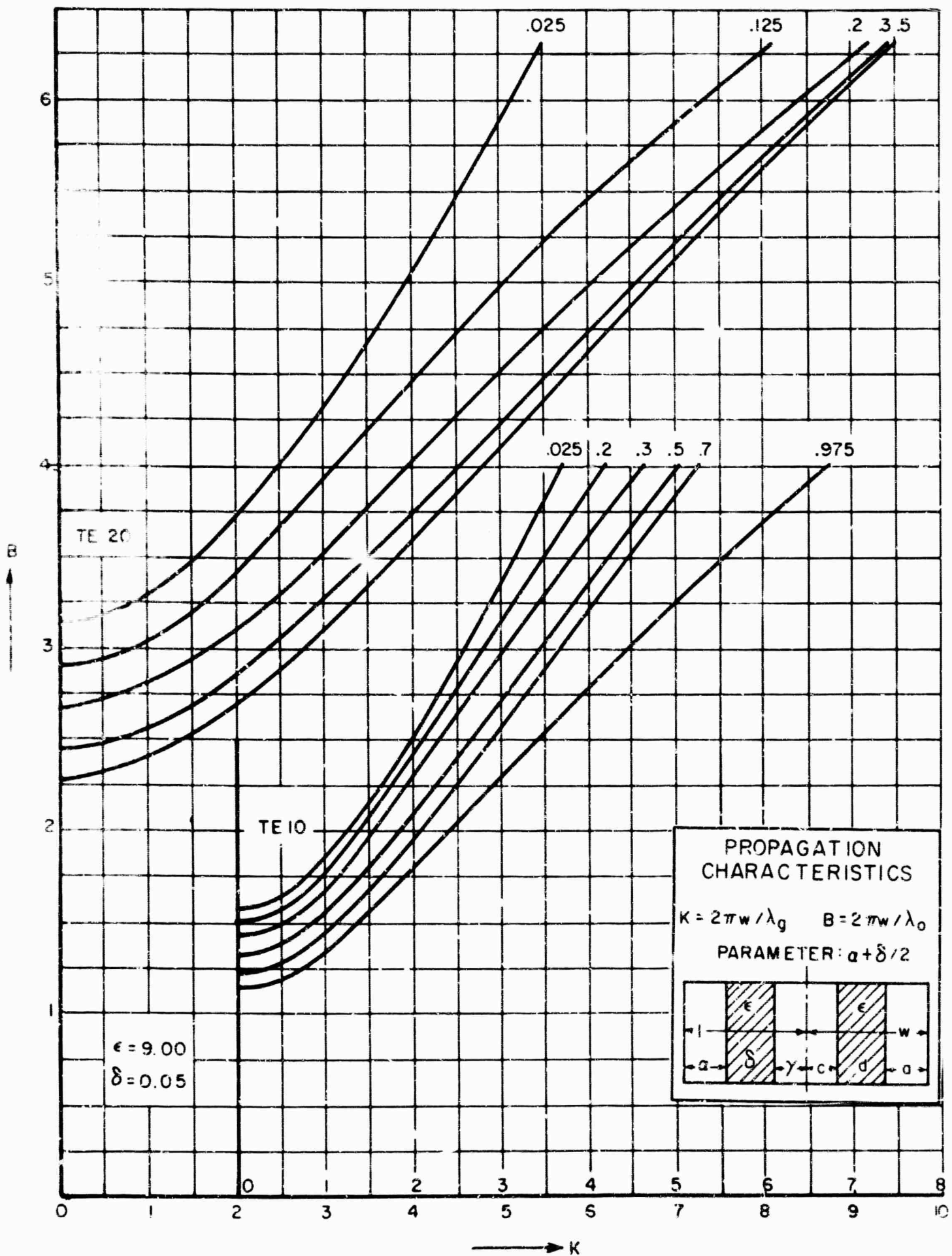


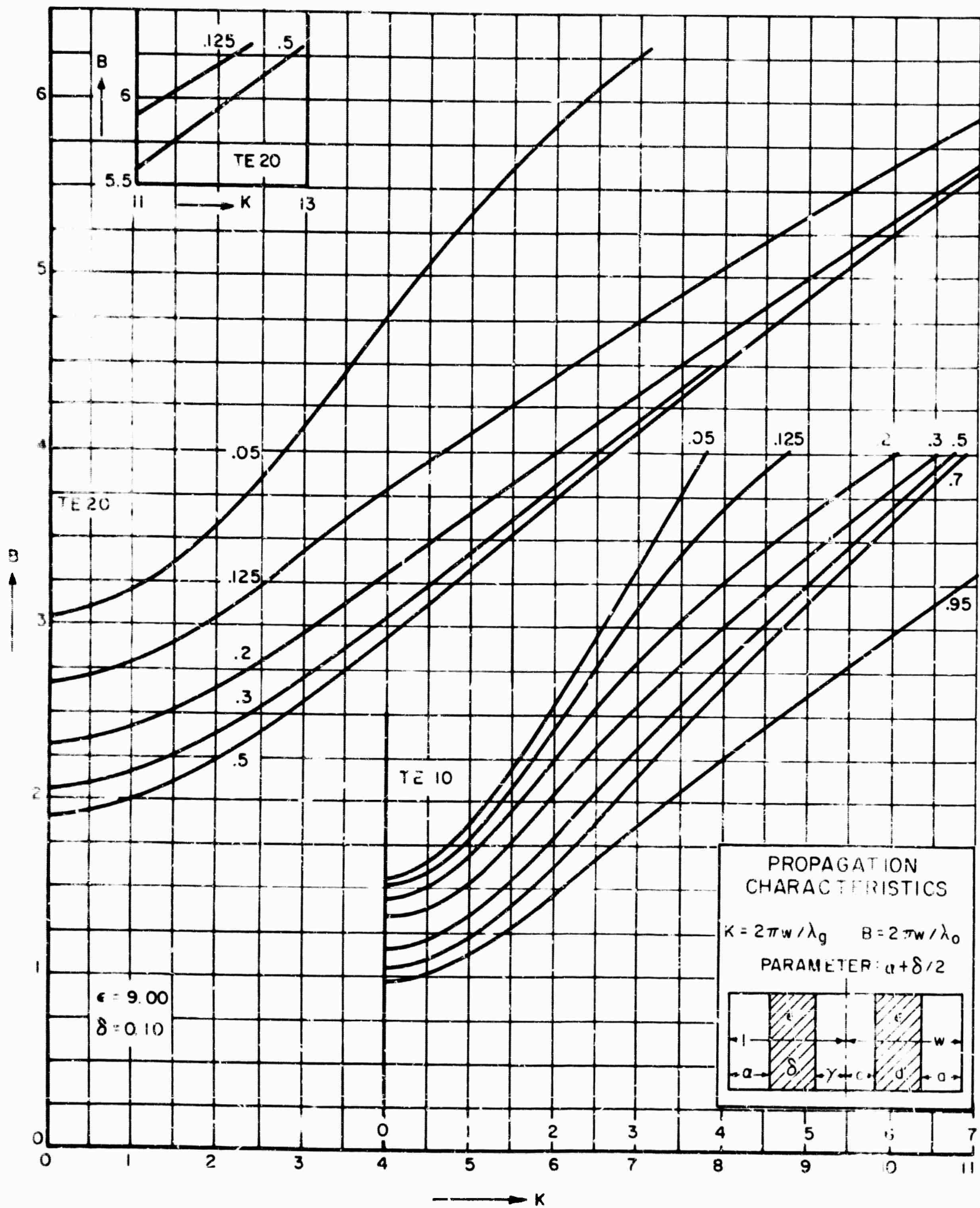


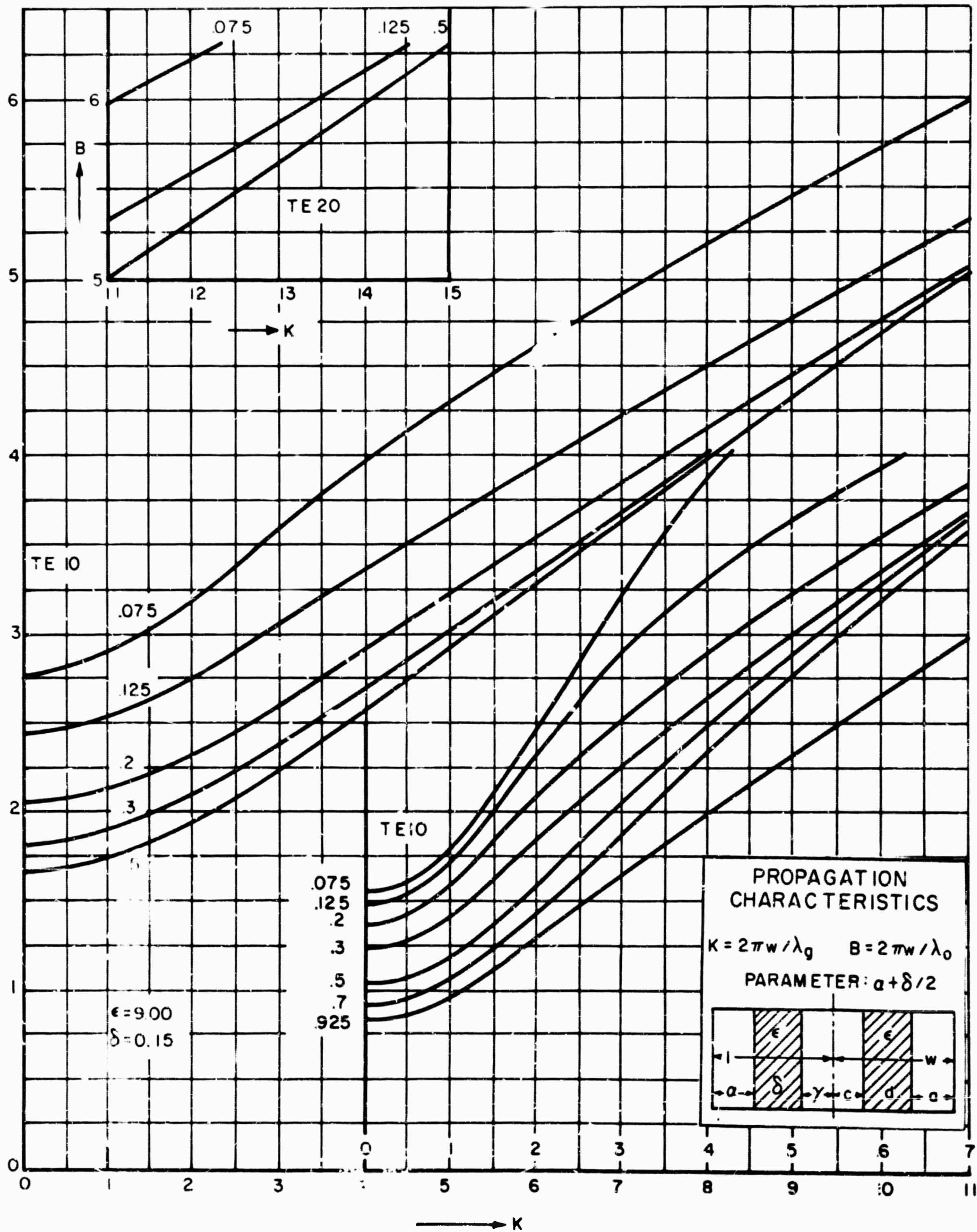


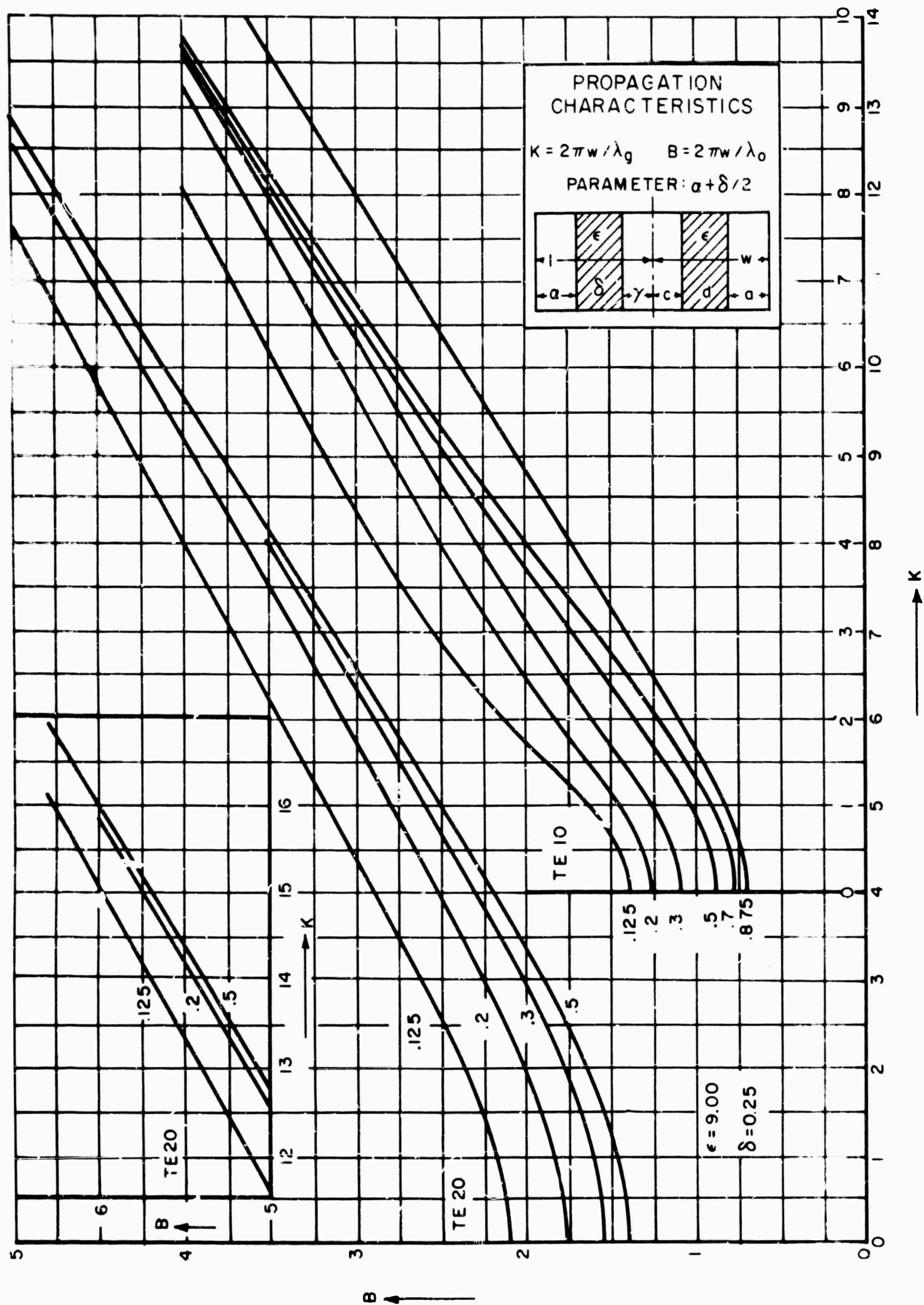




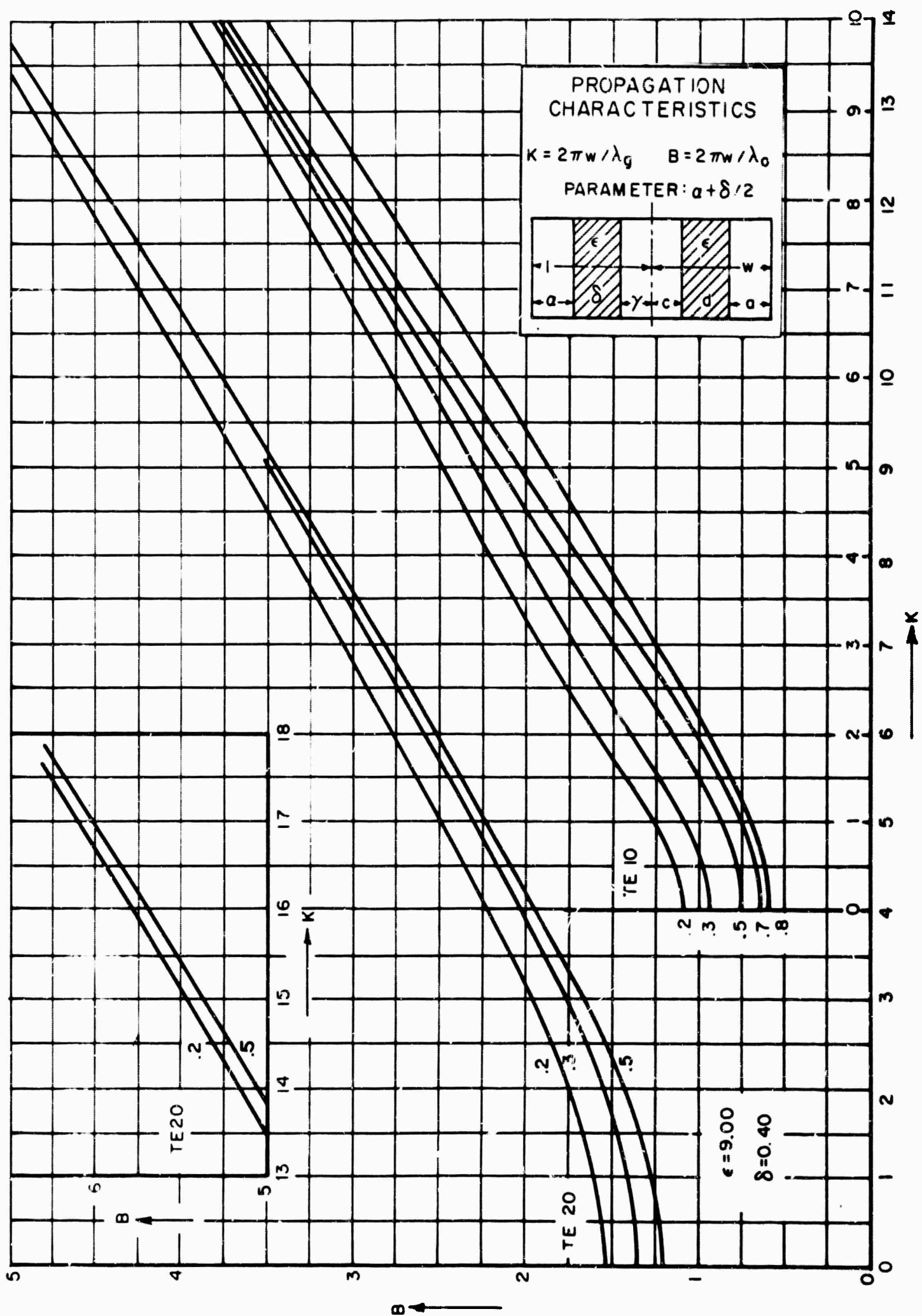


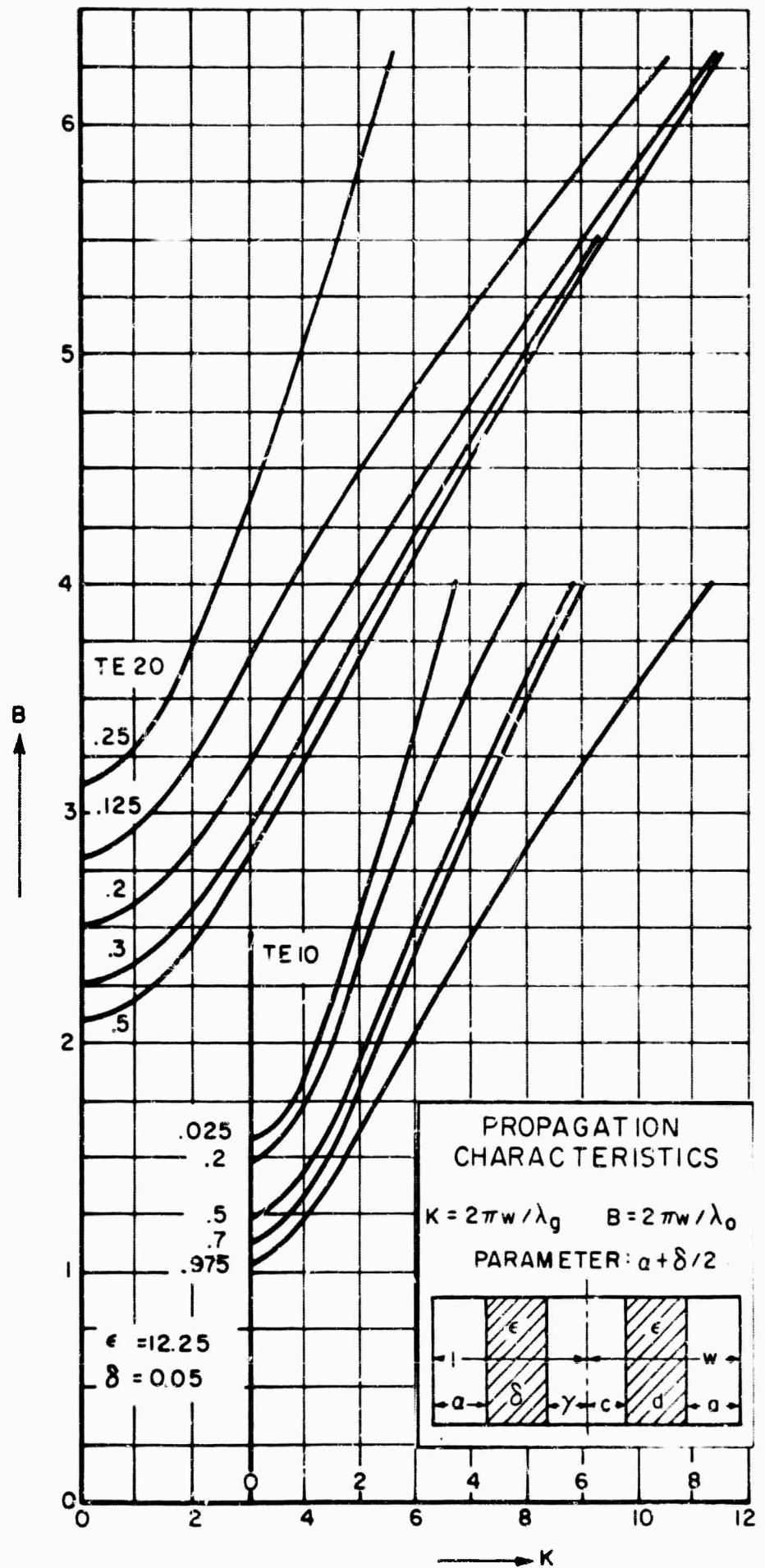


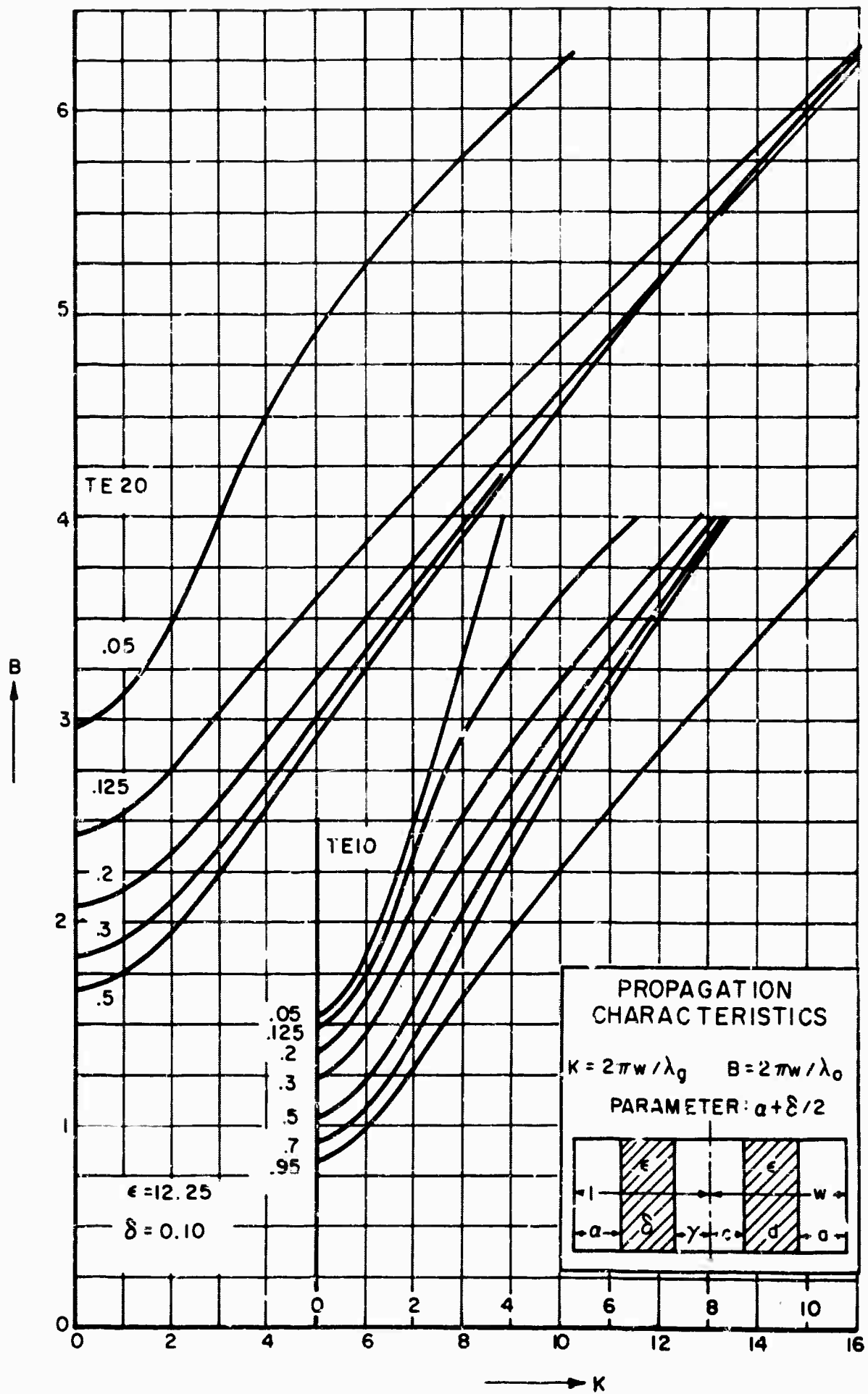


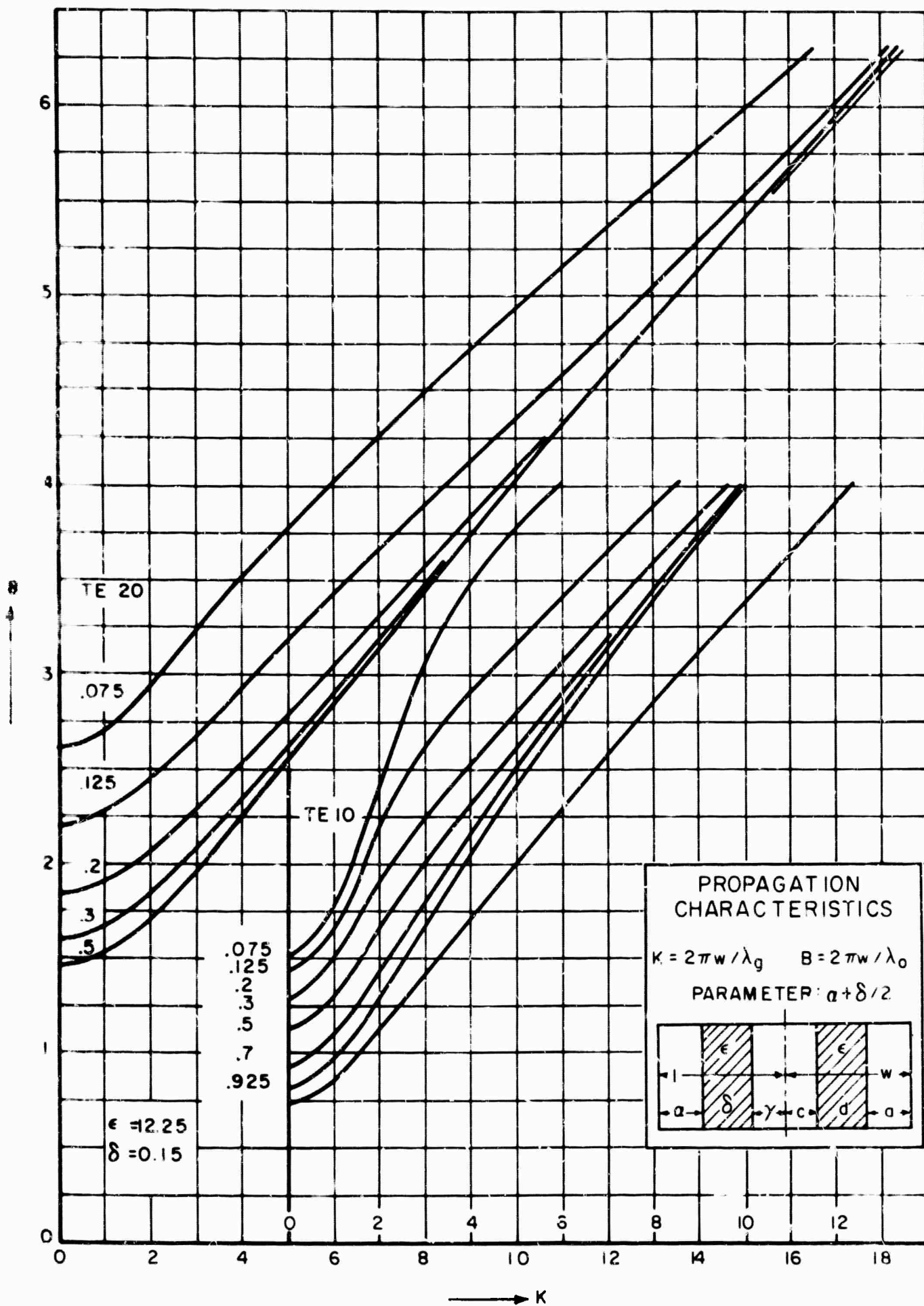




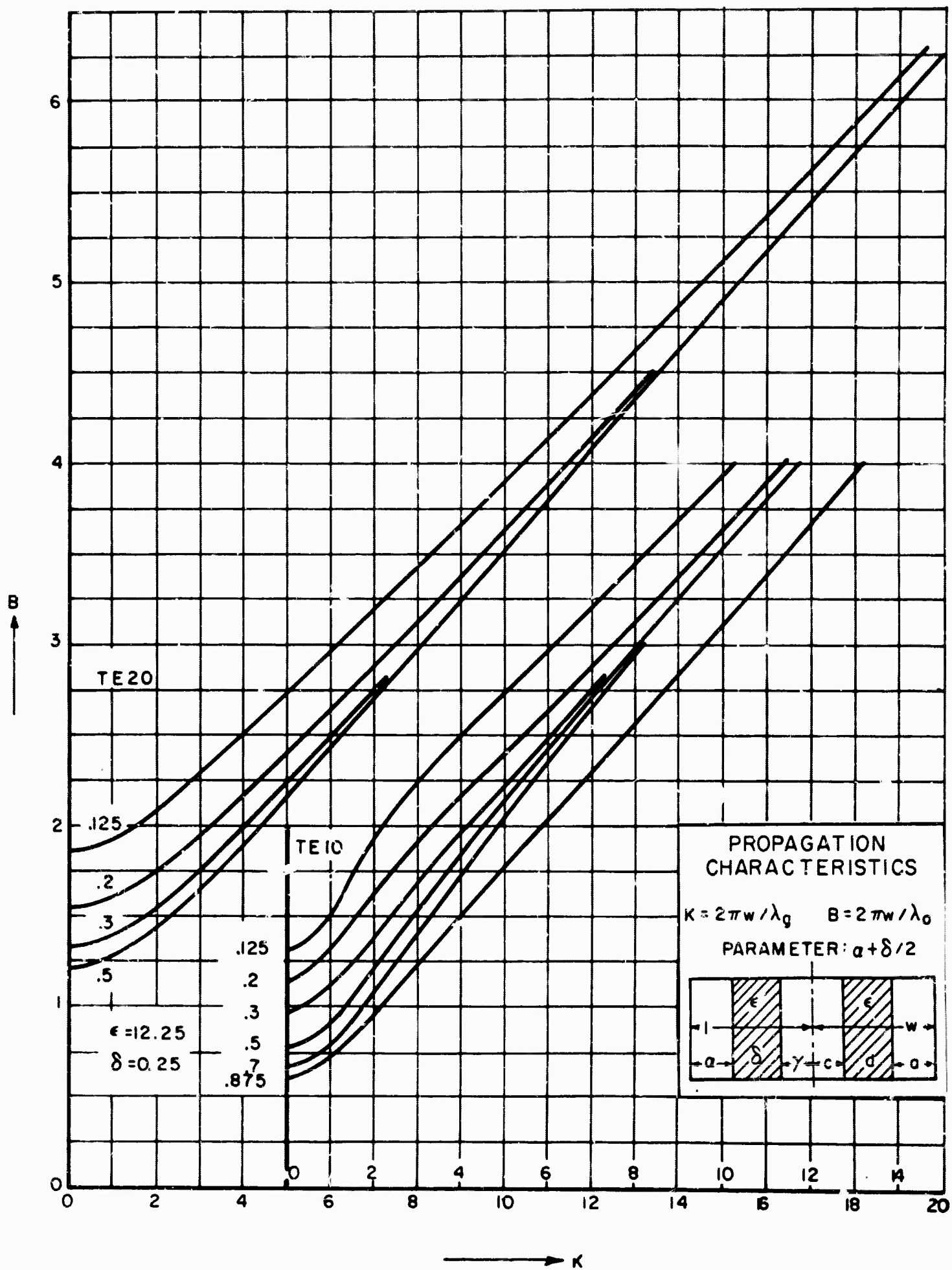


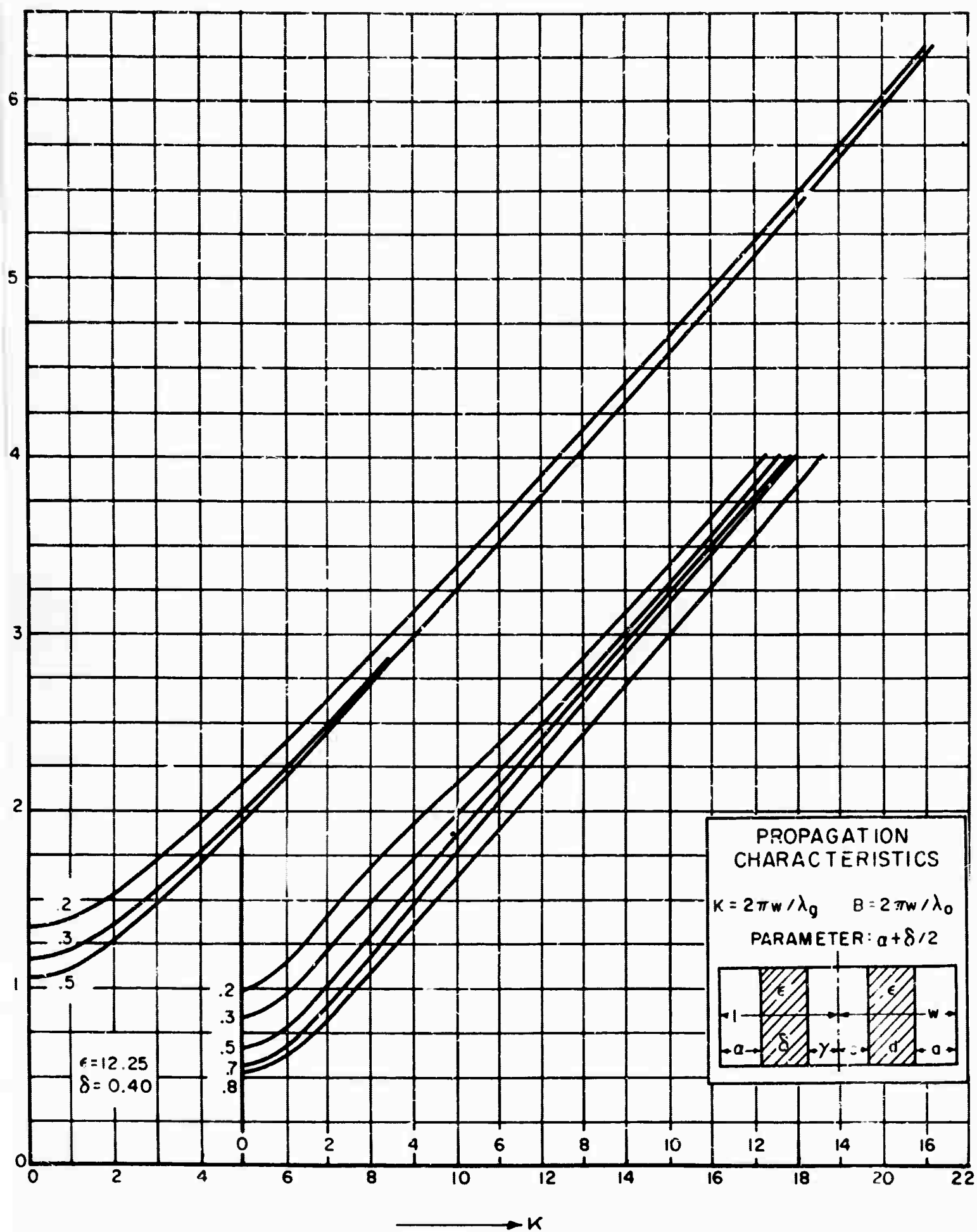


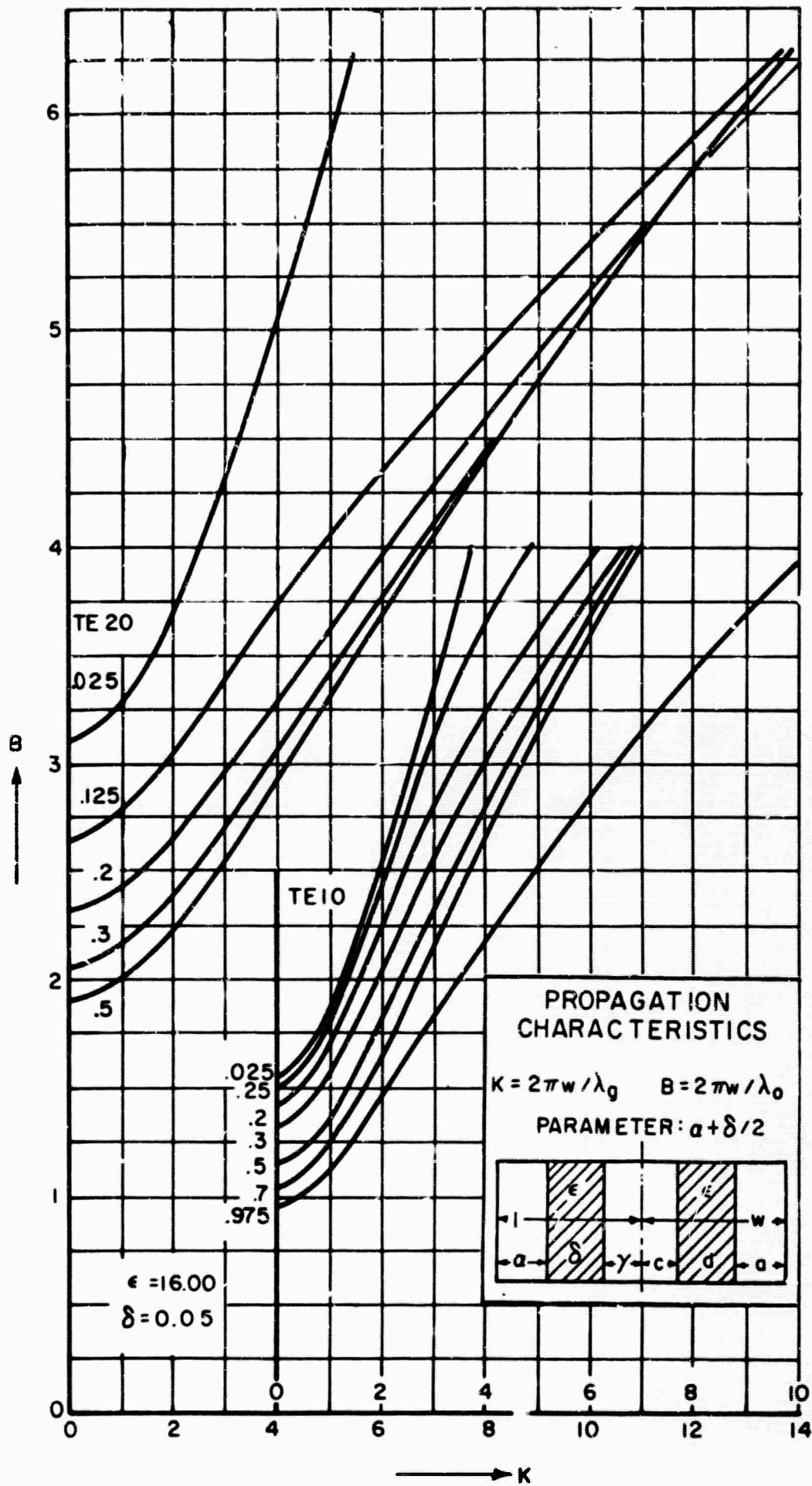


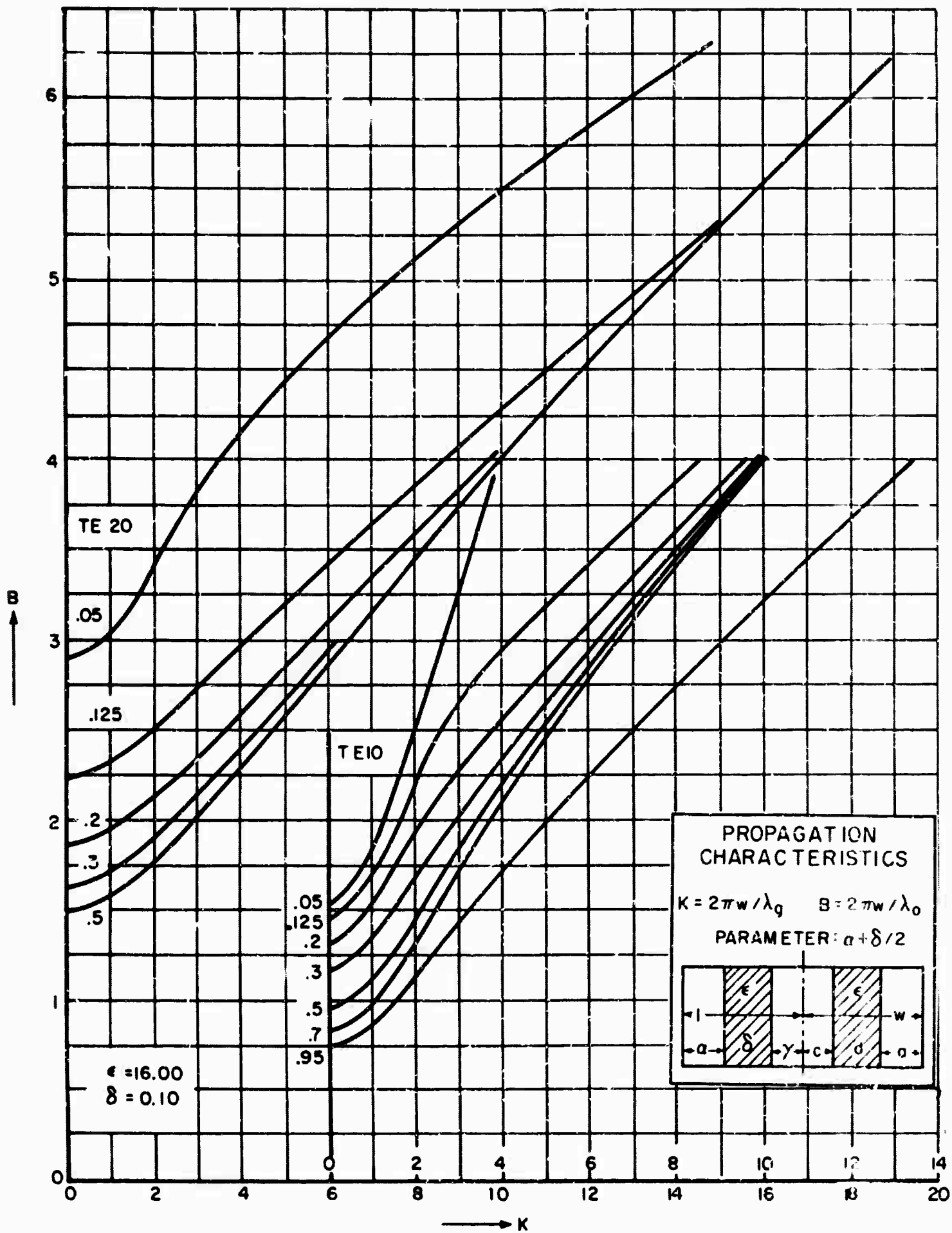


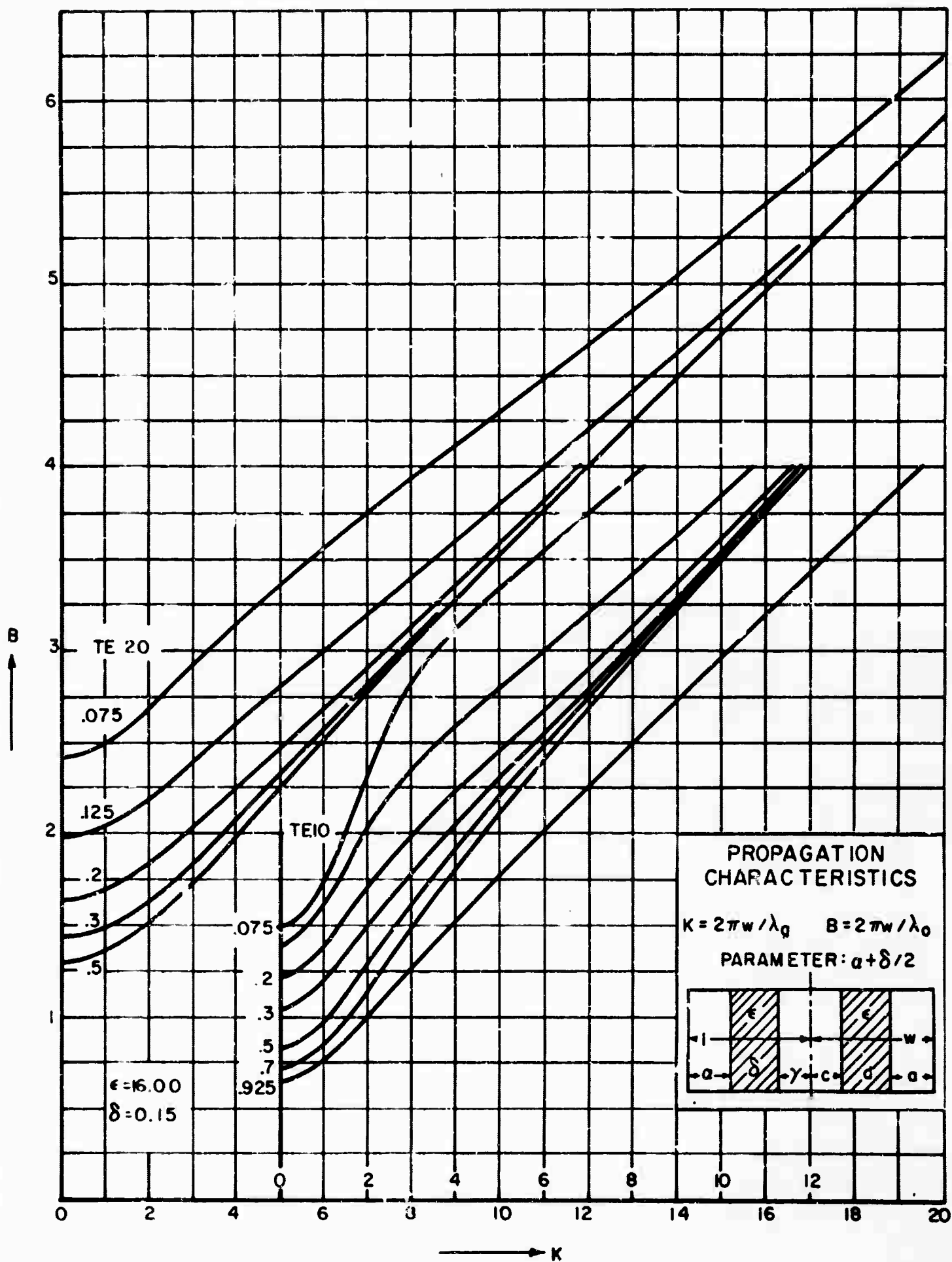


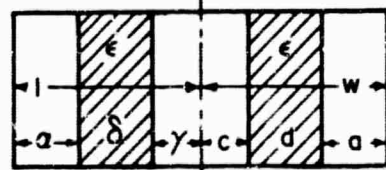
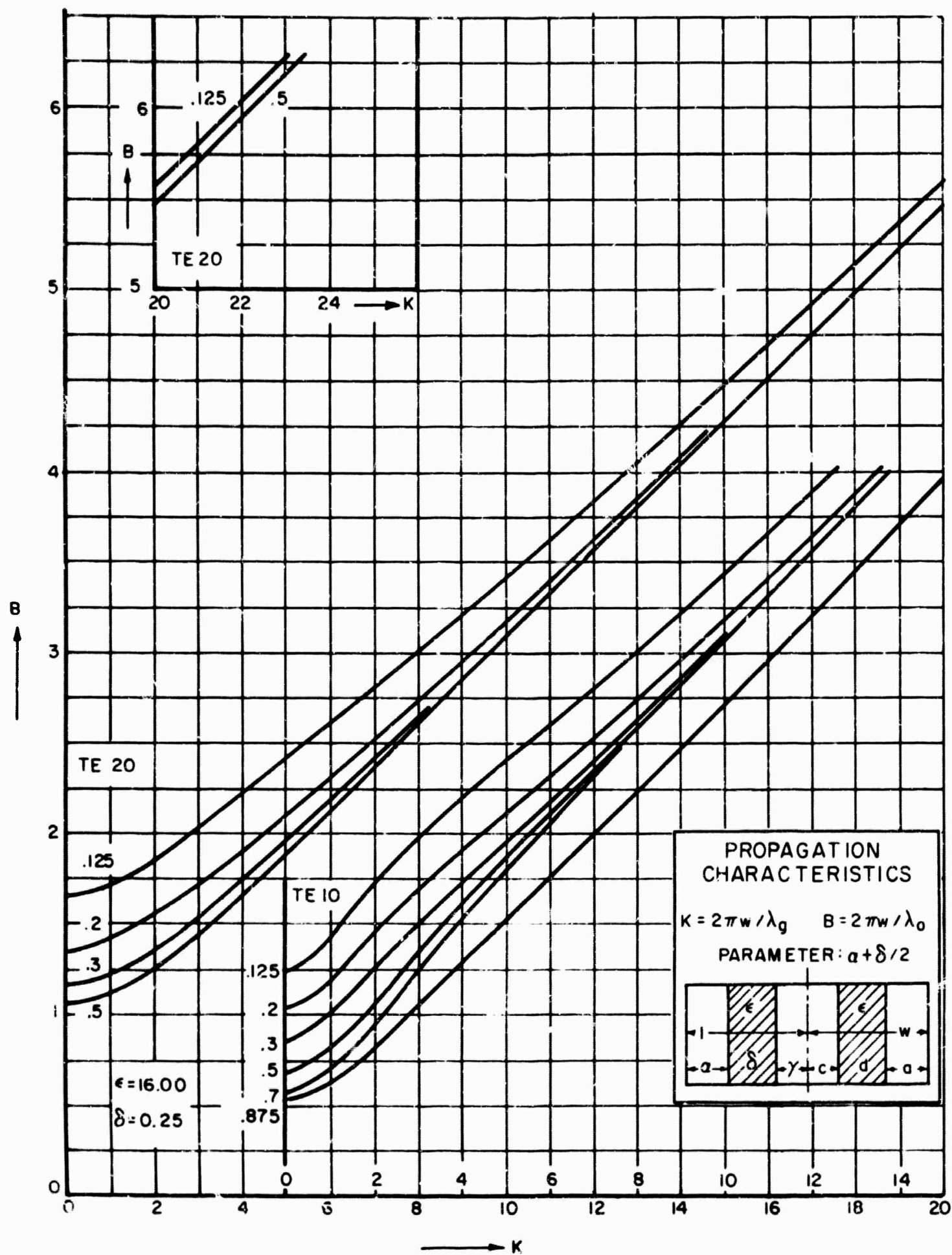




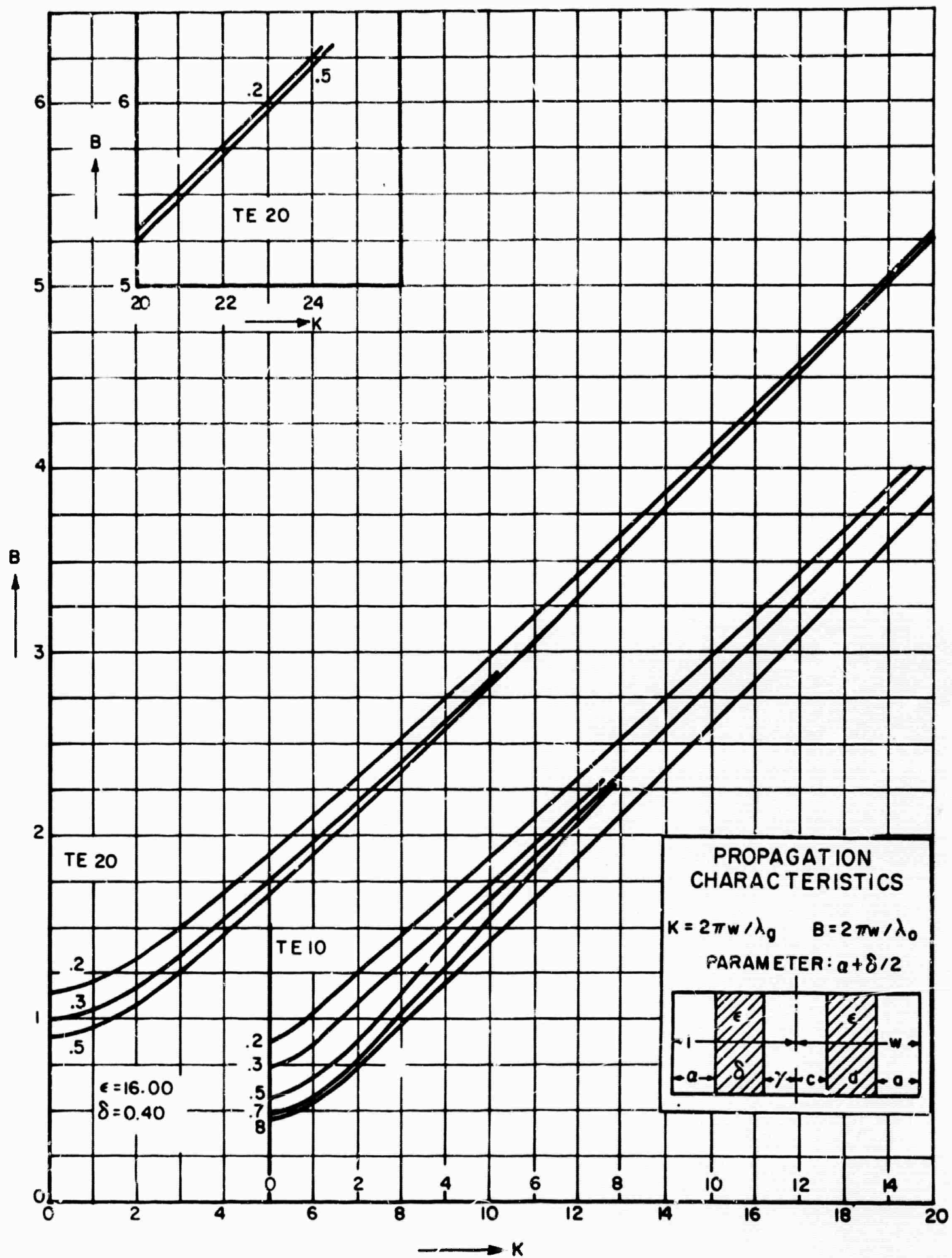


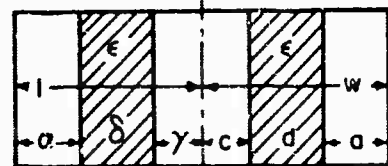
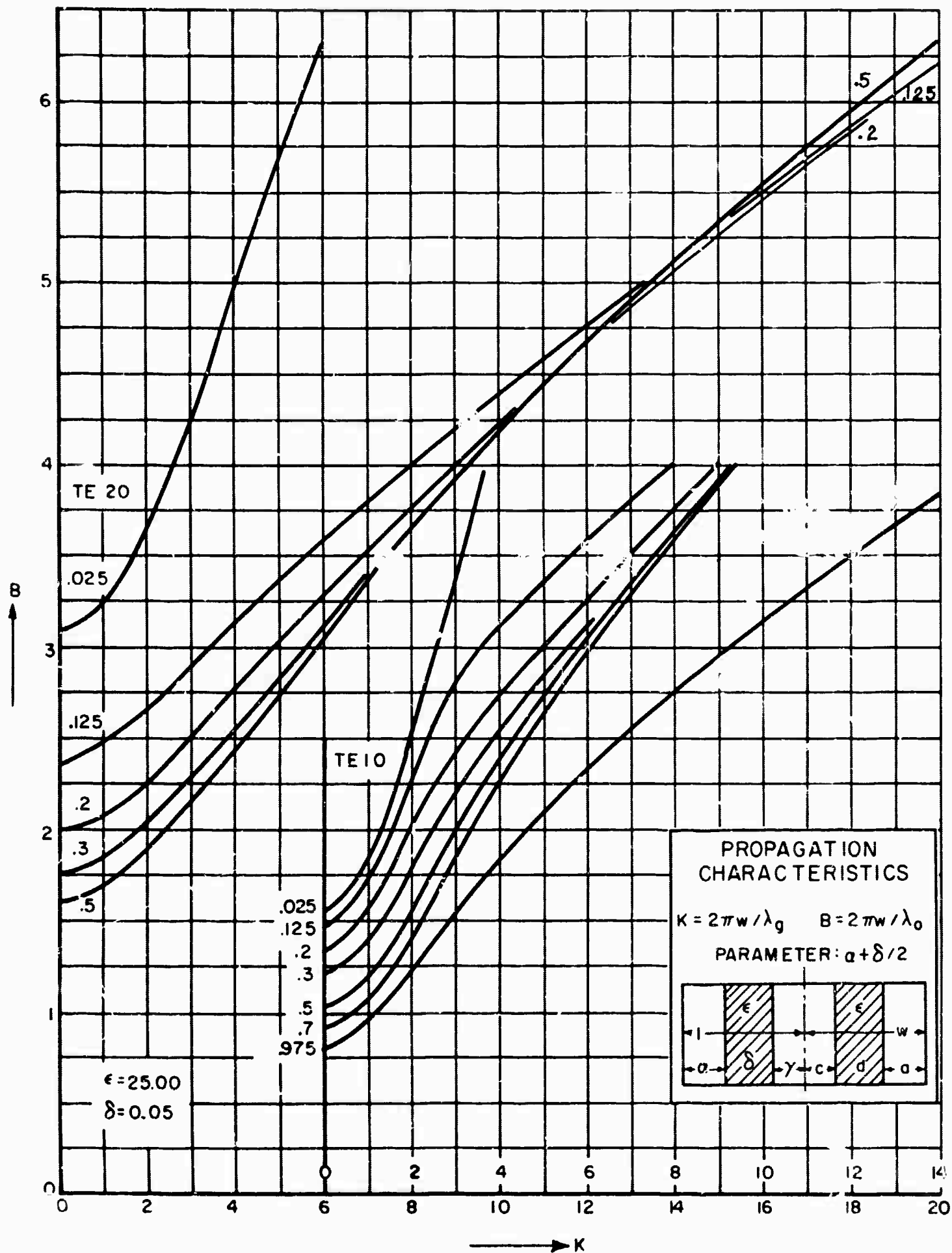




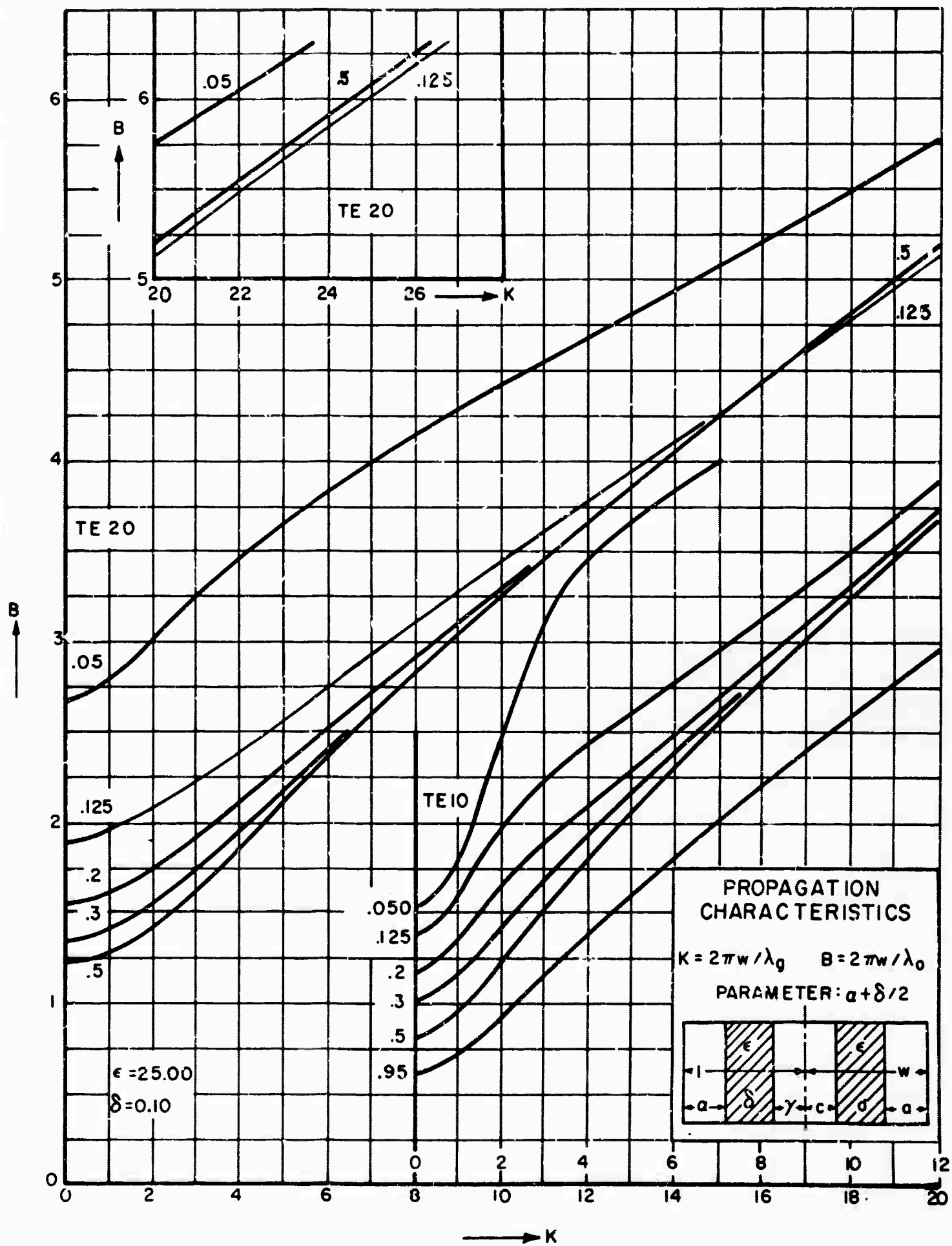


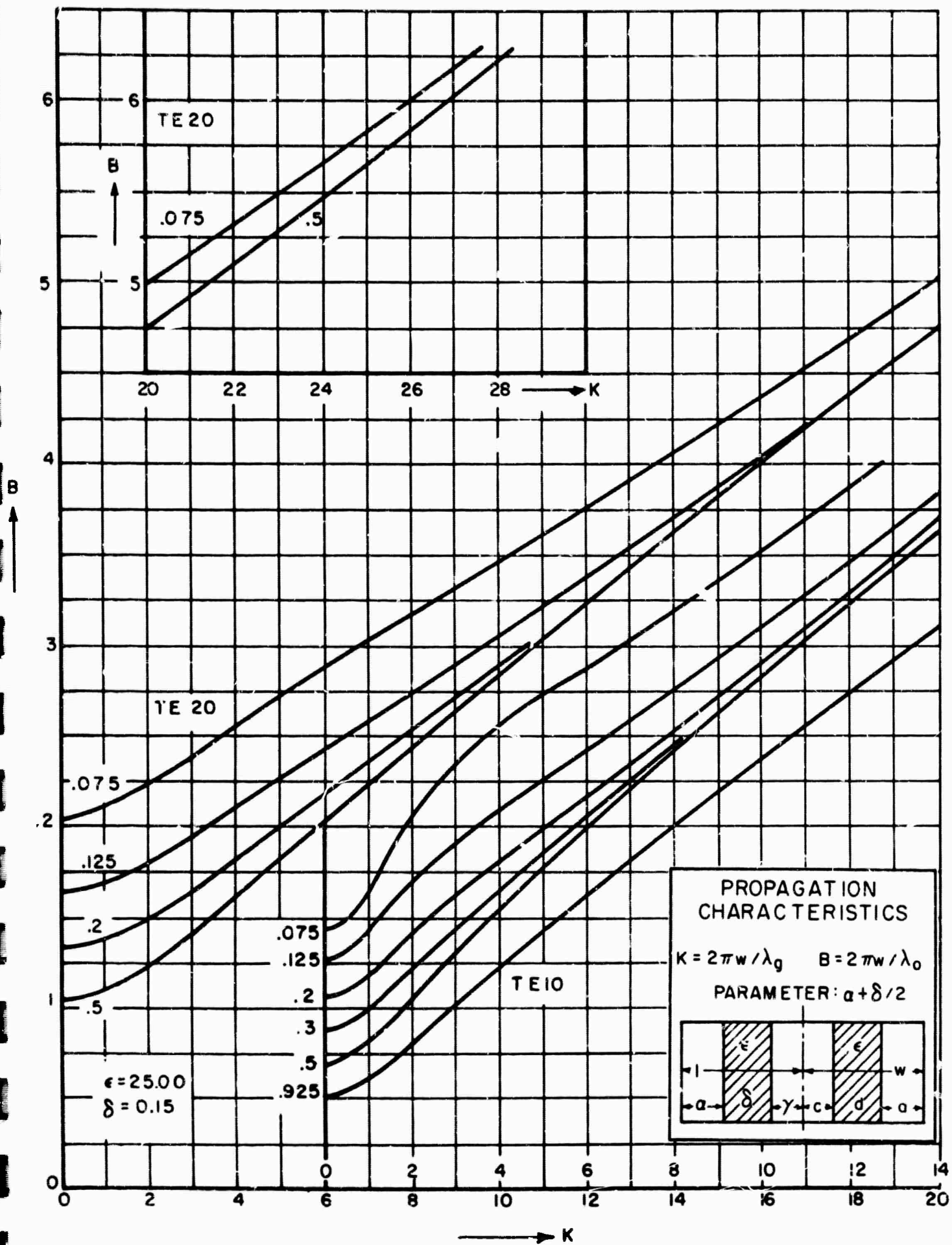


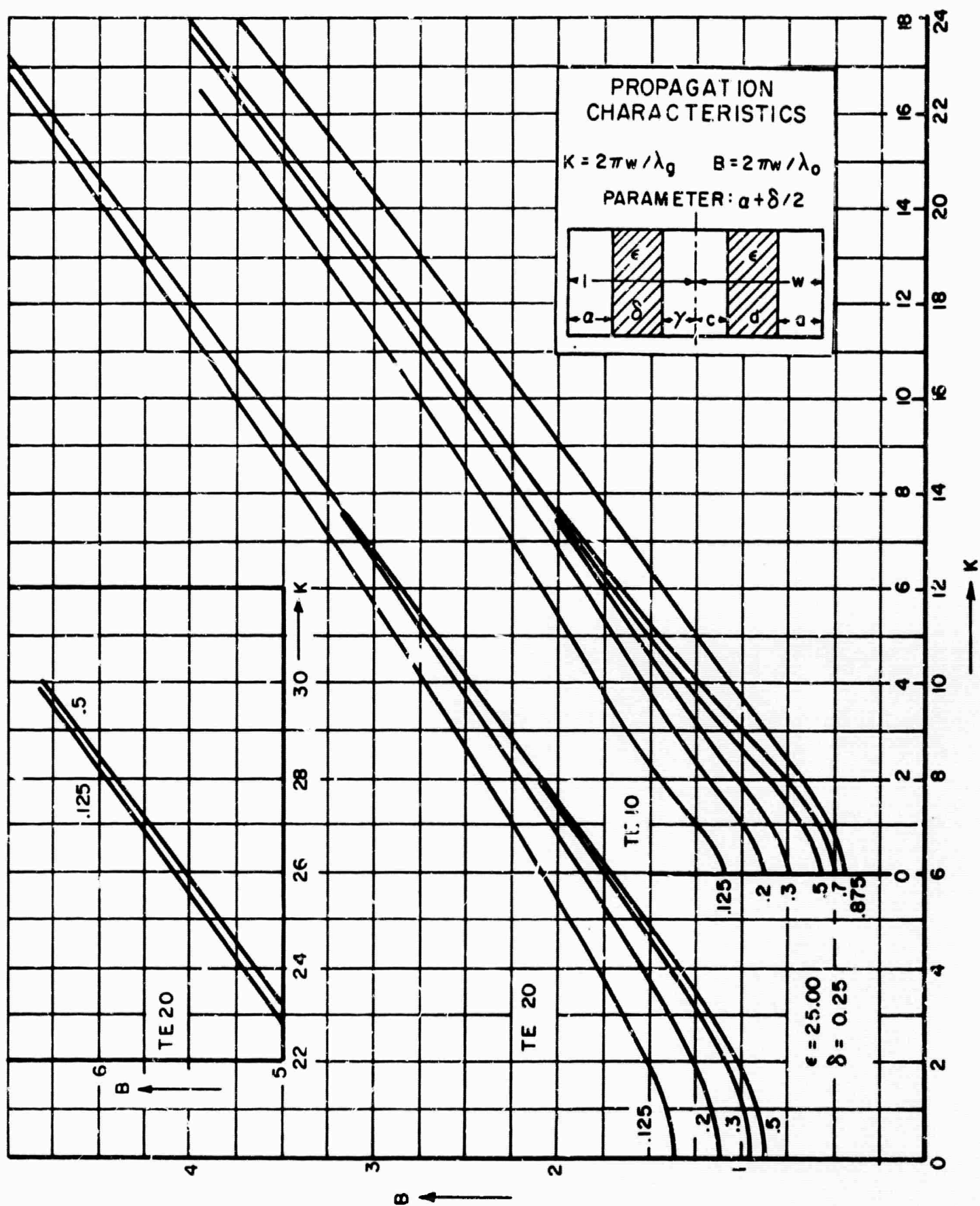














$a+\delta/2$	$N\delta$	0.0000	0.0250	0.0500	0.1000	0.1500	0.2000	0.2500	0.3000	0.4000	0.5000	0.6000	0.7000	0.8000	0.9000	1.0000
$\delta/2$	1	1.5708	1.5708	1.5706	1.5692	1.5653	1.5576	1.5457	1.5282	1.4768	1.4077	1.3299	1.2513	1.1769	1.1087	1.0472
( $a=0$ )	2	3.1416	3.1414	3.1400	3.1285	3.0972	3.0386	2.9524	2.8457	2.6180	2.2762	2.0762	1.8783	1.6817	1.4881	1.2994
	3	4.7124	4.7117	4.7069	4.6674	4.6430	4.5906	4.5188	4.4037	3.7592	3.4725	3.2647	3.0266	2.7905	2.5583	2.3311
	4	6.2832	6.2816	6.2700	6.1781	6.1462	6.0889	6.0232	5.9239	5.2360	4.9115	4.7298	4.5083	4.2834	4.0586	3.8338
	6	9.4248	9.4192	9.3791	9.0708	8.6215	8.4083	8.3776	8.3478	7.8540	7.3849	7.3266	7.1029	6.8771	6.6517	6.4262
0.0500	1	1.5708	1.5705	1.5701	1.5683	1.5655	1.5530	1.5411	1.5282	1.4768	1.4077	1.3299	1.2513	1.1769	1.1087	1.0472
	2	3.1416	3.1391	3.1363	3.1285	3.0972	3.0386	2.9524	2.8457	2.6180	2.2762	2.0762	1.8783	1.6817	1.4881	1.2994
	3	4.7124	4.7040	4.6943	4.6674	4.6430	4.5906	4.5188	4.4037	3.7592	3.4725	3.2647	3.0266	2.7905	2.5583	2.3311
	4	6.2832	6.2632	6.2398	6.1781	6.1462	6.0889	6.0232	5.9239	5.2360	4.9115	4.7298	4.5083	4.2834	4.0586	3.8338
	6	9.4248	9.3268	9.2760	9.0708	8.6215	8.4083	8.3776	8.3478	7.8540	7.3849	7.3266	7.1029	6.8771	6.6517	6.4262
0.1000	1	1.5708	1.5696	1.5683	1.5655	1.5621	1.5500	1.5381	1.5254	1.4768	1.4077	1.3299	1.2513	1.1769	1.1087	1.0472
	2	3.1416	3.1320	3.1218	3.0981	3.0719	2.9898	2.8983	2.7942	2.5432	2.2762	2.0762	1.8783	1.6817	1.4881	1.2994
	3	4.7124	4.6808	4.6465	4.5703	4.4846	4.4016	4.3199	4.2382	3.6679	3.4725	3.2647	3.0266	2.7905	2.5583	2.3311
	4	6.2832	6.2112	6.1322	6.0646	5.9821	5.8966	5.8119	5.7272	5.2360	4.9115	4.7298	4.5083	4.2834	4.0586	3.8338
	6	9.4248	9.2188	9.0074	8.6510	8.4875	8.3668	8.2863	8.2059	7.8540	7.3849	7.3266	7.1029	6.8771	6.6517	6.4262
0.2000	1	1.5708	1.5661	1.5613	1.5514	1.5410	1.5300	1.5181	1.5054	1.4768	1.4077	1.3299	1.2513	1.1769	1.1087	1.0472
	2	3.1416	3.1071	3.0718	3.0302	2.9287	2.8698	2.7942	2.7323	2.4297	2.2762	2.0762	1.8783	1.6817	1.4881	1.2994
	3	4.7124	4.6148	4.5162	4.4302	4.1707	4.0416	3.9406	3.8330	3.6679	3.4725	3.2647	3.0266	2.7905	2.5583	2.3311
	4	6.2832	6.1068	5.9409	5.8574	5.4777	5.3668	5.2863	5.2059	5.2360	4.9115	4.7298	4.5083	4.2834	4.0586	3.8338
	6	9.4248	9.1839	8.9139	8.7707	8.6386	8.5223	8.3776	8.2077	7.8514	7.3806					
0.3000	1	1.5708	1.5607	1.5505	1.5302	1.5098	1.4896	1.4693	1.4491	1.4092	1.3698	1.2523	1.2138	1.1093	1.0772	
	2	3.1416	3.0776	3.0149	2.9382	2.7893	2.6983	2.6137	2.5432	2.2997	2.3436	2.1826	2.1458	2.1018	2.0953	
	3	4.7124	4.6226	4.4432	4.2226	4.0523	3.9280	3.8319	3.7668	3.7135	3.6679	3.5458	3.5761	3.4683	3.4328	
	4	6.2832	6.1159	5.9719	5.7535	5.6082	5.5096	5.4360	5.3715	5.2360	5.0851	4.6601	4.4733	4.2433	4.1962	
	6	9.4248	9.3969	9.3624	9.2731	9.0863	8.7700	8.3776	8.0077	7.8640	7.5800	7.2274	6.9088	6.4453	6.3071	
0.4000	1	1.5708	1.5539	1.5373	1.5080	1.4739	1.4442	1.4169	1.3890	1.3332	1.2939	1.2523	1.2138	1.1442	1.0772	
	2	3.1416	3.0550	2.9733	2.8263	2.7012	2.5957	2.6071	2.4328	2.3182	2.2381	2.1826	2.1458	2.1188	2.0953	
	3	4.7124	4.5846	4.4690	4.2782	4.1298	4.0282	3.9380	3.8753	3.7846	3.7135	3.5458	3.5761	3.4683	3.4328	
	4	6.2832	6.2176	6.1574	6.0624	5.9587	5.8591	5.7376	5.6877	5.2360	4.9049	4.6601	4.4733	4.2433	4.1962	
	6	9.4248	9.3226	9.2242	9.0456	8.8605	8.6311	8.3776	8.1639	7.8640	7.5800	7.2274	6.9088	6.4453	6.3071	
0.6000	1	1.5708	1.5466	1.5233	1.4790	1.4380	1.4002	1.3663	1.3330	1.2767	1.2261	1.1826	1.1442	1.1093	1.0772	
	2	3.1416	3.0466	2.9584	2.8024	2.6719	2.5632	2.4728	2.3974	2.2824	2.2034	2.1510	2.1188	2.1018	2.0953	
	3	4.7124	4.6402	4.5720	4.4806	4.3497	4.2546	4.1888	4.1167	3.9603	3.7886	3.6177	3.4683	3.3380	3.2328	
	4	6.2832	6.2828	6.2759	6.2389	6.1946	6.0776	5.9048	5.6914	5.2360	4.8440	4.5523	4.3565	4.2433	4.1962	
	6	9.4248	9.1451	8.9101	8.5937	8.4398	8.3953	8.3776	8.3700	8.1896	7.7224	7.1813	6.7382	6.4453	6.3071	
0.8000	1	1.5708	1.5398	1.5097	1.4547	1.4063	1.3609	1.3208	1.2845	1.2218	1.1688	1.1240	1.0851	1.0509	1.0177	1.0472
	3	4.7124	4.6982	4.6832	4.6490	4.6077	4.5448	4.4680	4.3641	4.1142	3.8485	3.5983	3.3904	3.2226	3.1538	3.1416
0.7000	1	1.5708	1.5331	1.4978	1.4337	1.3774	1.3277	1.2836	1.2442	1.1770	1.1218	1.0755				
	3	4.7124	4.7087	4.7041	4.6891	4.6602	4.6106	4.5351	4.4321	4.1627	3.8652	3.5907				
0.8000	1	1.5708	1.6280	1.4813	1.4170	1.3562	1.3013	1.2639	1.2120	1.1412						
	3	4.7124	4.6542	4.6211	4.5477	4.4888	4.4274	4.3647	4.2908	4.0948						
1- $\delta/2$	1	1.5708	1.5233	1.4792	1.4012	1.3360	1.2816	1.2364	1.1987	1.1412	1.1017	1.0768	1.0594	1.0509	1.0477	1.0472
( $\gamma=0$ )	3	4.7124	4.5726	4.4861	4.2969	4.2199	4.1926	4.1888	4.1880	4.0948	3.8612	3.5907	3.3691	3.2226	3.1538	3.1416

WAVE GUIDE WIDTH  
2w

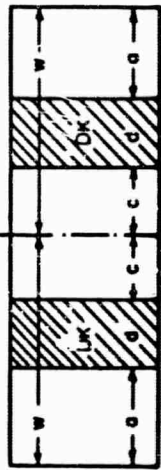
RELATIVE DIELECTRIC CONSTANT  
DK

FREE SPACE PROPAGATION CONSTANT  
 $\beta$

NORMALIZED CUTOFF FREQUENCY  
 $B=\beta_w$

VELOCITY OF LIGHT IN FREE SPACE  
 $v$

CUTOFF FREQUENCY  
 $f=8v/2\pi w$





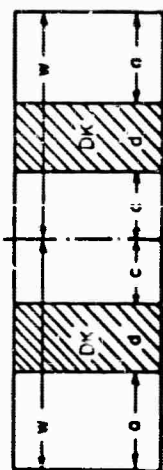
NORMALIZED CUTOFF FREQUENCIES OF A WAVEGUIDE WHICH CONTAINS DIELECTRIC SLABS

DK = 4.00

$\alpha + \delta/2$	$\delta/2$	0.0000	0.0250	0.0500	0.1000	0.1500	0.2000	0.2500	0.3000	0.4000	0.5000	0.6000	0.7000	0.8000	0.9000	1.0000
1	1.5708	1.5707	1.5703	1.5699	1.5694	1.5689	1.5684	1.5679	1.5674	1.5669	1.5664	1.5659	1.5654	1.5649	1.5644	1.5639
2	3.1416	3.1411	3.1376	3.1341	3.1306	3.1271	3.1236	3.1201	3.1166	3.1131	3.1096	3.1061	3.1026	3.0991	3.0956	3.0921
3	4.7124	4.7107	4.6987	4.6867	4.6747	4.6627	4.6507	4.6387	4.6267	4.6147	4.6027	4.5907	4.5787	4.5667	4.5547	4.5427
4	6.2832	6.2792	6.2692	6.2592	6.2492	6.2392	6.2292	6.2192	6.2092	6.1992	6.1892	6.1792	6.1692	6.1592	6.1492	6.1392
6	9.4248	9.4110	9.2992	9.1874	9.0756	8.9638	8.8520	8.7402	8.6284	8.5166	8.4048	8.2930	8.1812	8.0694	7.9576	7.8458
0.0500	1	1.5708	1.5701	1.5692												
	2	3.1416	3.1385	3.1285												
	3	4.7124	4.6918	4.6660												
	4	6.2832	6.2322	6.1666												
	6	9.4248	9.2411	8.9887												
0.1000	1	1.5708	1.5679	1.5648	1.5619	1.5589	1.5559	1.5529	1.5499	1.5469	1.5439	1.5409	1.5379	1.5349	1.5319	1.5289
	2	3.1416	3.1180	3.0918	3.0656	3.0394	3.0132	2.9870	2.9608	2.9346	2.9084	2.8822	2.8560	2.8298	2.8036	2.7774
	3	4.7124	4.6324	4.5392	4.4328	4.3184	4.2040	4.0896	3.9752	3.8608	3.7464	3.6320	3.5176	3.4032	3.2888	3.1744
	4	6.2832	6.0970	5.8790	5.6457	5.4124	5.1791	4.9458	4.7125	4.4792	4.2459	4.0126	3.7793	3.5460	3.3127	3.0794
	6	9.4248	8.9028	8.4433	7.9923	7.6698	7.3473	7.0248	6.7023	6.3798	6.0573	5.7348	5.4123	5.0898	4.7673	4.4448
0.2000	1	1.5708	1.5594	1.5477	1.5353	1.5229	1.5105	1.4981	1.4857	1.4733	1.4609	1.4485	1.4361	1.4237	1.4113	1.3989
	2	3.1416	3.0873	2.9696	2.7989	2.6418	2.4847	2.3276	2.1705	2.0134	1.8563	1.6992	1.5421	1.3850	1.2279	1.0708
	3	4.7124	4.4747	4.2442	3.9682	3.6095	3.2508	2.8921	2.5334	2.1747	1.8160	1.4573	1.0986	0.7399	0.3812	-0.0775
	4	5.2832	5.8706	5.5332	5.1108	4.6918	4.2728	3.8538	3.4348	3.0158	2.5968	2.1778	1.7588	1.3398	0.9208	0.5018
	6	9.4248	8.9236	8.6838	8.4092	8.2126	7.9840	7.7554	7.5268	7.2982	7.0696	6.8410	6.6124	6.3838	6.1552	5.9266
0.3000	1	1.5708	1.5468	1.5221	1.4974	1.4728	1.4481	1.4235	1.3988	1.3742	1.3495	1.3248	1.3001	1.2754	1.2507	1.2260
	2	3.1416	2.9896	2.8468	2.6026	2.4123	2.2649	2.1175	1.9701	1.8227	1.6753	1.5279	1.3805	1.2331	1.0857	0.9383
	3	4.7124	4.3910	4.1268	3.7840	3.5239	3.3792	3.2345	3.0898	2.9451	2.8004	2.6557	2.5110	2.3663	2.2216	2.0769
	4	6.2832	5.9142	5.6882	5.3573	5.2032	5.0946	4.9785	4.8624	4.7463	4.6302	4.5141	4.3980	4.2819	4.1658	4.0497
	6	9.4248	9.3623	9.3026	9.1048	8.6018	7.8840	7.2130	6.7939	6.4388	6.2832	6.1276	5.9720	5.8164	5.6608	5.5052
0.4000	1	1.5708	1.5308	1.4922	1.4206	1.3864	1.2996	1.2492	1.2044	1.1283	1.0687	1.0124	0.9587	0.9049	0.8511	0.7973
	2	3.1416	2.9410	2.7672	2.4924	2.2913	2.1406	2.0246	1.9336	1.8028	1.7164	1.6594	1.6221	1.5848	1.5475	1.5102
	3	4.7124	4.4231	4.1948	3.8814	3.6913	3.5593	3.4638	3.4188	3.2889	3.1416	2.9976	2.8710	2.7544	2.6378	2.5212
	4	6.2832	6.1341	6.0157	5.8481	5.7061	5.5171	5.2468	4.9262	4.3287	3.8942	3.6084	3.4212	3.2340	3.0468	2.8596
	6	9.4248	9.1838	8.9640	8.6878	8.3244	7.8840	7.4774	7.2809	6.8760	6.2832	5.7248	5.3381	4.9255	4.5129	4.1003
0.5000	1	1.5708	1.5141	1.4619	1.3699	1.2926	1.2272	1.1714	1.1232	1.0442	0.9818	0.9306	0.8874	0.8442	0.8010	0.7578
	2	3.1416	2.9239	2.7404	2.4874	2.2840	2.1027	1.9869	1.8962	1.7664	1.6821	1.6261	1.5834	1.5407	1.4980	1.4553
	3	4.7124	4.5446	4.4008	4.1780	4.0392	3.9270	3.8213	3.7047	3.4278	3.1414	2.8560	2.7074	2.5622	2.4170	2.2718
	4	6.2832	6.2822	6.2753	6.2169	6.0820	5.7813	5.3602	4.9892	4.2893	3.8213	3.5064	3.3090	3.1161	2.9232	2.7303
	6	9.4248	8.8038	8.3878	7.9880	7.5694	7.0840	6.6337	6.2132	5.0972	4.2832	3.6266	3.1629	2.7255	2.3081	1.8907
0.6000	1	1.5708	1.4982	1.4334	1.3286	1.2388	1.1678	1.1086	1.0584	0.9779	0.9187	0.8686	0.8244	0.7891	0.7538	0.7185
	3	4.7124	4.6782	4.6430	4.5678	4.4749	4.3476	4.1788	3.9671	3.5282	3.1416	2.8407	2.6109	2.4363	2.2617	2.0871
0.7000	1	1.5708	1.4842	1.4097	1.2887	1.1946	1.1197	1.0582	1.0068	0.9283	0.8633	0.8140				
	3	4.7124	4.7036	4.6933	4.6601	4.5942	4.4746	4.2914	4.0886	3.6636	3.1416	2.8133				
0.8000	1	1.5708	1.4730	1.3908	1.2592	1.1597	1.0814	1.0180	0.9688	0.8832						
	3	4.7124	4.6049	4.5224	4.4067	4.3218	4.2343	4.1191	3.9604	3.5486						
1- $\delta/2$ ( $\gamma=0$ )	1	1.5708	1.4619	1.3702	1.2287	1.1270	1.0513	0.9934	0.9461	0.8632	0.8141	0.7977	0.7891	0.7889	0.7889	0.7884
	3	4.7124	4.4017	4.1939	3.9940	3.9347	3.9270	3.9194	3.8663	3.8486	3.8146	3.8133	2.8185	2.4363	2.3662	2.3862

WAVE GUIDE WIDTH  
RELATIVE DIELECTRIC CONSTANT  
FREE SPACE PROPAGATION CONSTANT  
NORMALIZED CUTOFF FREQUENCY  
VELOCITY OF LIGHT IN FREE SPACE  
CUTOFF FREQUENCY

$$f = 8v/2\pi W$$



$$\alpha = a/W$$

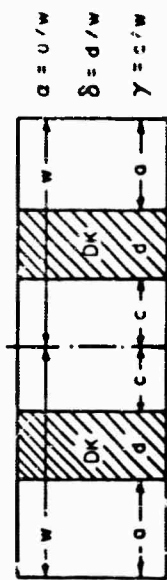
$$\delta = b/H$$

$$\gamma = c/W$$

NORMALIZED CUTOFF FREQUENCIES OF A WAVEGUIDE WHICH CONTAINS DIELECTRIC SLABS DK = 12.25

$\alpha + \delta/2$	$M/S$	0.0000	0.0250	0.0500	0.1000	0.1500	0.2000	0.2500	0.3000	0.4000	0.5000	0.6000	0.7000	0.8000	0.9000	1.0000
$\delta/2$ ( $\alpha=0$ )	1	1.5708	1.5706	1.5689	1.5651	1.5618	1.4353	1.3235	1.2015	0.9858	0.8264	0.7085	0.6192	0.5497	0.4942	0.4488
	2	3.1416	3.1397	3.1256	2.9832	2.5961	2.1745	1.8545	1.6219	1.3186	1.1377	1.0249	0.9584	0.9165	0.9002	0.8976
	3	4.7124	4.7059	4.6512	4.0459	3.2228	2.7657	2.5881	2.4832	2.4687	2.3801	2.0973	1.8518	1.6482	1.4828	1.3464
	4	6.2832	6.2671	6.1085	4.8132	4.1530	4.0419	4.0338	3.9195	3.2807	2.7271	2.3426	2.0778	1.9043	1.8138	1.7952
	6	9.4248	9.3621	8.5416	7.2048	7.1872	6.5122	5.5785	5.0565	4.3739	3.7497	3.2822	2.9491	2.7450	2.6928	
0.0500	1	1.5708	1.5679	1.5645												
	2	3.1416	3.1170	3.0847												
	3	4.7124	4.6184	4.4712												
	4	6.2832	6.0202	5.5812												
	6	9.4248	8.3495	7.8340												
0.1000	1	1.5708	1.5594	1.5468	1	1.4803										
	2	3.1416	3.0414	2.9122	2	3.777										
	3	4.7124	4.3410	3.8956	3	2.9589										
	4	6.2832	5.4329	4.7489	4.2129	4.0728										
	6	9.4248	7.8300	7.4486	7.2764	7.0067										
0.2000	1	1.5708	1.5267	1.4797	1.3840	1.2947	1.2162	1.1481	1.0884							
	2	3.1416	2.8103	2.5028	2.0800	1.8274	1.6613	1.5429	1.4527							
	3	4.7124	3.8697	3.3363	2.8367	2.6312	2.5368	2.4931	2.4750							
	4	6.2832	5.0757	4.6240	4.3518	4.2502	4.1248	3.9091	3.6626							
	6	9.4248	8.3697	8.1746	7.7831	6.4917	5.5413	5.1190	4.9812							
0.3000	1	1.5708	1.4798	1.3929	1.2444	1.1303	1.0426	0.9734	0.9172	0.8301	0.7635					
	2	3.1416	2.6191	2.2548	1.8293	1.5916	1.4390	1.3325	1.2539	1.1461	1.0756					
	3	4.7124	3.7558	3.2962	2.9118	2.7618	2.6846	2.6282	2.5673	2.4048	2.2365					
	4	6.2832	5.3361	5.0348	4.8285	4.6457	4.2882	3.7732	3.3677	2.8306	2.5286					
	6	9.4248	9.2503	9.1303	8.0906	6.4698	5.7361	5.5079	5.3415	4.6663	4.0849					
0.4000	1	1.5708	1.4284	1.3094	1.1327	1.0115	0.9236	0.8566	0.8037	0.7242	0.6661	0.6204	0.5826			
	2	3.1416	2.8106	2.1322	1.7144	1.4859	1.3399	1.2382	1.1638	1.0623	0.9994	0.9590	0.9330			
	3	4.7124	3.8826	3.8067	3.1988	3.0596	2.9748	2.8564	2.6968	2.3453	2.0625	1.8792	1.7478			
	4	6.2832	5.8560	5.6738	5.4664	5.0313	4.3293	3.7222	3.2687	2.6755	2.3252	2.1137	1.9848			
	6	9.4248	8.7035	8.4228	7.7583	6.7402	6.3707	6.1309	5.6500	4.8868	3.8620	3.3984	3.1156			
0.5000	1	1.5708	1.3796	1.2380	1.0472	0.9248	0.8388	0.7745	0.7242	0.6498	0.5965	0.5555	0.5225	0.4947	0.4704	
	2	3.1416	2.4761	2.0951	1.6805	1.4849	1.3109	1.2107	1.1370	1.0375	0.9590	0.9376	0.9148	0.9028	0.8983	
	3	4.7124	4.1888	3.9116	3.6635	3.5203	3.3389	3.0785	2.7919	2.3135	1.9855	1.7609	1.6051	1.4944	1.4125	
	4	6.2832	6.2795	6.2512	5.9664	5.1922	4.3490	3.7091	3.2437	2.6372	2.2755	2.0499	1.9108	1.8329	1.8004	
	6	9.4248	7.8299	7.3772	7.1863	7.1752	7.0180	6.4598	5.7322	4.8705	3.8185	3.3248	3.0005	2.8024	2.7094	
0.6000	1	1.5708	1.3367	1.1793	0.9810	0.8593	0.7755	0.7135	0.6655	0.5949	0.5447	0.5067	0.4765	0.4514		
	3	4.7124	4.5861	4.4786	4.3052	4.0558	3.6583	3.2197	2.8436	2.2971	1.9407	1.6979	1.5260	1.4012		
0.7000	1	1.5708	1.3004	1.1320	0.9291	0.8083	0.7264	0.6664	0.6201	0.5526	0.5048	0.4688				
	3	4.7124	4.6821	4.5523	4.5511	4.2718	3.7777	3.2764	2.8664	2.2887	1.9167	1.6624				
0.8000	1	1.5708	1.2723	1.0943	0.8876	0.7676	0.6873	0.6289	0.5840	0.5187						
	3	4.7124	4.4143	4.2831	4.1558	4.0129	3.7049	3.2707	2.8703	2.2852						
1- $\delta/2$ ( $\gamma=0$ )	1	1.5708	1.2380	1.0476	0.8403	0.7274	0.6555	0.6053	0.5685	0.5187	0.4880	0.4688	0.4574	0.4514	0.4491	0.4488
	3	4.7124	3.7149	3.6886	3.5932	3.5876	3.5075	3.2299	2.8661	2.2852	1.9093	1.6624	1.5003	1.4012	1.3547	1.3464

WAVE GUIDE WIDTH  $2w$   
RELATIVE DIELECTRIC CONSTANT  $DK$   
FREE SPACE PROPAGATION CONSTANT  $\beta$   
NORMALIZED CUTOFF FREQUENCY  $B = \beta w$   
VELOCITY OF LIGHT IN FREE SPACE  $v$   
CUTOFF FREQUENCY  $f = Bv/2\pi w$

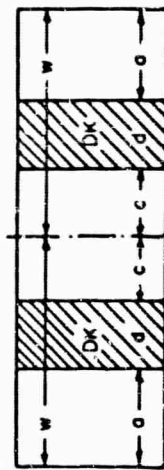


# NORMALIZED CUTOFF FREQUENCIES OF A WAVEGUIDE WHICH CONTAINS DIELECTRIC SLABS

DK = 9.00

$\alpha + \delta/2$	$\delta/2$	$\alpha = 0$	0.0000	0.0250	0.0500	0.1000	0.1500	0.2000	0.2500	0.3000	0.4000	0.5000	0.6000	0.7000	0.8000	0.9000	1.0000
1	1.5708	1.5706	1.5695	1.5685	1.5675	1.5665	1.5655	1.5645	1.5635	1.5625	1.5615	1.5605	1.5595	1.5585	1.5575	1.5565	1.5555
2	3.1416	3.1403	3.1306	3.0391	2.7737	2.4111	2.0944	1.8499	1.6181	1.3963	1.1912	1.0107	0.8517	0.7136	0.5965	0.5001	0.4236
3	4.7124	4.7079	4.6719	4.2896	3.5702	3.0661	2.7925	2.6671	2.6180	2.5732	2.5376	2.5061	2.4785	2.4545	2.4334	2.4152	2.3994
4	6.2832	6.2721	6.1745	5.2360	4.4397	4.2147	4.1888	4.1638	4.1416	4.1213	4.1025	4.0851	4.0691	4.0545	4.0413	4.0293	4.0184
5	7.8540	7.8437	7.7461	6.8076	5.9113	5.2360	4.7138	4.3888	4.2147	4.1888	4.1638	4.1416	4.1213	4.1025	4.0851	4.0691	4.0545
6	9.4248	9.4145	9.3169	8.3784	7.4821	6.7068	6.1846	5.8596	5.6846	5.5596	5.4346	5.3096	5.1846	5.0596	4.9346	4.8096	4.6846

WAVE GUIDE WIDTH  
RELATIVE DIELECTRIC CONSTANT  
FREE SPACE PROPAGATION CONSTANT  
NORMALIZED CUTOFF FREQUENCY  
VELOCITY OF LIGHT IN FREE SPACE  
CUTOFF FREQUENCY



$a = 0/w$   
 $\delta = d/v$   
 $\gamma = c/w$

2w

DK

$\beta$

$B = \beta w$

$v$

$f = 8v/2w$

0.8000	1	1.5708	1.5698	1.5688	1.5678	1.5668	1.5658	1.5648	1.5638	1.5628	1.5618	1.5608	1.5598	1.5588	1.5578	1.5568	1.5558
2	3.1416	3.1406	3.1396	3.1386	3.1376	3.1366	3.1356	3.1346	3.1336	3.1326	3.1316	3.1306	3.1296	3.1286	3.1276	3.1266	3.1256
3	4.7124	4.7114	4.7104	4.7094	4.7084	4.7074	4.7064	4.7054	4.7044	4.7034	4.7024	4.7014	4.7004	4.6994	4.6984	4.6974	4.6964
4	6.2832	6.2822	6.2812	6.2802	6.2792	6.2782	6.2772	6.2762	6.2752	6.2742	6.2732	6.2722	6.2712	6.2702	6.2692	6.2682	6.2672
5	7.8540	7.8530	7.8520	7.8510	7.8500	7.8490	7.8480	7.8470	7.8460	7.8450	7.8440	7.8430	7.8420	7.8410	7.8400	7.8390	7.8380
6	9.4248	9.4238	9.4228	9.4218	9.4208	9.4198	9.4188	9.4178	9.4168	9.4158	9.4148	9.4138	9.4128	9.4118	9.4108	9.4098	9.4088

0.4000	1	1.5708	1.5698	1.5688	1.5678	1.5668	1.5658	1.5648	1.5638	1.5628	1.5618	1.5608	1.5598	1.5588	1.5578	1.5568	1.5558
2	3.1416	3.1406	3.1396	3.1386	3.1376	3.1366	3.1356	3.1346	3.1336	3.1326	3.1316	3.1306	3.1296	3.1286	3.1276	3.1266	3.1256
3	4.7124	4.7114	4.7104	4.7094	4.7084	4.7074	4.7064	4.7054	4.7044	4.7034	4.7024	4.7014	4.7004	4.6994	4.6984	4.6974	4.6964
4	6.2832	6.2822	6.2812	6.2802	6.2792	6.2782	6.2772	6.2762	6.2752	6.2742	6.2732	6.2722	6.2712	6.2702	6.2692	6.2682	6.2672
5	7.8540	7.8530	7.8520	7.8510	7.8500	7.8490	7.8480	7.8470	7.8460	7.8450	7.8440	7.8430	7.8420	7.8410	7.8400	7.8390	7.8380
6	9.4248	9.4238	9.4228	9.4218	9.4208	9.4198	9.4188	9.4178	9.4168	9.4158	9.4148	9.4138	9.4128	9.4118	9.4108	9.4098	9.4088

0.5000	1	1.5708	1.5698	1.5688	1.5678	1.5668	1.5658	1.5648	1.5638	1.5628	1.5618	1.5608	1.5598	1.5588	1.5578	1.5568	1.5558
2	3.1416	3.1406	3.1396	3.1386	3.1376	3.1366	3.1356	3.1346	3.1336	3.1326	3.1316	3.1306	3.1296	3.1286	3.1276	3.1266	3.1256
3	4.7124	4.7114	4.7104	4.7094	4.7084	4.7074	4.7064	4.7054	4.7044	4.7034	4.7024	4.7014	4.7004	4.6994	4.6984	4.6974	4.6964
4	6.2832	6.2822	6.2812	6.2802	6.2792	6.2782	6.2772	6.2762	6.2752	6.2742	6.2732	6.2722	6.2712	6.2702	6.2692	6.2682	6.2672
5	7.8540	7.8530	7.8520	7.8510	7.8500	7.8490	7.8480	7.8470	7.8460	7.8450	7.8440	7.8430	7.8420	7.8410	7.8400	7.8390	7.8380
6	9.4248	9.4238	9.4228	9.4218	9.4208	9.4198	9.4188	9.4178	9.4168	9.4158	9.4148	9.4138	9.4128	9.4118	9.4108	9.4098	9.4088

0.6000	1	1.5708	1.5698	1.5688	1.5678	1.5668	1.5658	1.5648	1.5638	1.5628	1.5618	1.5608	1.5598	1.5588	1.5578	1.5568	1.5558
2	3.1416	3.1406	3.1396	3.1386	3.1376	3.1366	3.1356	3.1346	3.1336	3.1326	3.1316	3.1306	3.1296	3.1286	3.1276	3.1266	3.1256
3	4.7124	4.7114	4.7104	4.7094	4.7084	4.7074	4.7064	4.7054	4.7044	4.7034	4.7024	4.7014	4.7004	4.6994	4.6984	4.6974	4.6964
4	6.2832	6.2822	6.2812	6.2802	6.2792	6.2782	6.2772	6.2762	6.2752	6.2742	6.2732	6.2722	6.2712	6.2702	6.2692	6.2682	6.2672
5	7.8540	7.8530	7.8520	7.8510	7.8500	7.8490	7.8480	7.8470	7.8460	7.8450	7.8440	7.8430	7.8420	7.8410	7.8400	7.8390	7.8380
6	9.4248	9.4238	9.4228	9.4218	9.4208	9.4198	9.4188	9.4178	9.4168	9.4158	9.4148	9.4138	9.4128	9.4118	9.4108	9.4098	9.4088

0.7000	1	1.5708	1.5698	1.5688	1.5678	1.5668	1.5658	1.5648	1.5638	1.5628	1.5618	1.5608	1.5598	1.5588	1.5578	1.5568	1.5558
2	3.1416	3.1406	3.1396	3.1386	3.1376	3.1366	3.1356	3.1346	3.1336	3.1326	3.1316	3.1306	3.1296	3.1286	3.1276	3.1266	3.1256
3	4.7124	4.7114	4.7104	4.7094	4.7084	4.7074	4.7064	4.7054	4.7044	4.7034	4.7024	4.7014	4.7004	4.6994	4.6984	4.6974	4.6964
4	6.2832	6.2822	6.2812	6.2802	6.2792	6.2782	6.2772	6.2762	6.2752	6.2742	6.2732	6.2722	6.2712	6.2702	6.2692	6.2682	6.2672
5	7.8540	7.8530	7.8520	7.8510	7.8500	7.8490	7.8480	7.8470	7.8460	7.8450	7.8440	7.8430	7.8420	7.8410	7.8400	7.8390	7.8380
6	9.4248	9.4238	9.4228	9.4218	9.4208	9.4198	9.4188	9.4178	9.4168	9.4158	9.4148	9.4138	9.4128	9.4118	9.4108	9.4098	9.4088

0.8000	1	1.5708	1.5698	1.5688	1.5678	1.5668	1.5658	1.5648	1.5638	1.5628	1.5618	1.5608	1.5598	1.5588	1.5578	1.5568	1.5558
2	3.1416	3.1406	3.1396	3.1386	3.1376	3.1366	3.1356	3.1346	3.1336	3.1326	3.1316	3.1306	3.1296	3.1286	3.1276	3.1266	3.1256
3	4.7124	4.7114	4.7104	4.7094	4.7084	4.7074	4.7064	4.7054	4.7044	4.7034	4.7024	4.7014	4.7004	4.6994	4.6984	4.6974	4.6964
4	6.2832	6.2822	6.2812	6.2802	6.2792	6.2782	6.2772	6.2762	6.2752	6.2742	6.2732	6.2722	6.2712	6.2702	6.2692	6.2682	6.2672
5	7.8540	7.8530	7.8520	7.8510	7.8500	7.8490	7.8480	7.8470	7.8460	7.8450	7.8440	7.8430	7.8420	7.8410	7.8400	7.8390	7.8380
6	9.4248	9.4238	9.4228	9.4218	9.4208	9.4198	9.4188	9.4178	9.4168	9.4158	9.4148	9.4138	9.4128	9.4118	9.4108	9.4098	9.4088

1- $\delta/2$	1	1.5708	1.5698	1.5688	1.5678	1.5668	1.5658	1.5648	1.5638	1.5628	1.5618	1.5608	1.5598	1.5588	1.5578	1.5568	1.5558
2	3.1416	3.1406	3.1396	3.1386	3.1376	3.1366	3.1356	3.1346	3.1336	3.1326	3.1316	3.1306	3.1296	3.1286	3.1276	3.1266	3.1256
3	4.7124	4.7114	4.7104	4.7094	4.7084	4.7074	4.7064	4.7054	4.7044	4.7034	4.7024	4.7014	4.7004	4.6994	4.6984	4.6974	4.6964
4	6.2832	6.2822	6.2812	6.2802	6.2792	6.2782	6.2772	6.2762	6.2752	6.2742	6.2732	6.2722	6.2712	6.2702	6.2692	6.2682	6.2672
5	7.8540	7.8530	7.8520	7.8510	7.8500	7.8490	7.8480	7.8470	7.8460	7.8450	7.8440	7.8430	7.8420	7.8410	7.8400	7.8390	7.8380
6	9.4248	9.4238	9.4228	9.4218	9.4208	9.4198	9.4188	9.4178	9.4168	9.4158	9.4148	9.4138	9.4128	9.4118	9.4108	9.4098	9.4088

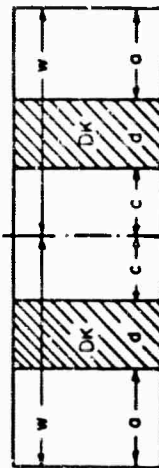
NORMALIZED CUTOFF FREQUENCIES OF A WAVEGUIDE WHICH CONTAINS DIELECTRIC SLABS

DK = 16.00



$a\delta/2$	$N/\delta$	0.0000	0.0250	0.0500	0.1000	0.1500	0.2000	0.2500	0.3000	0.4000	0.5000	0.6000	0.7000	0.8000	0.9000	1.0000
$\delta/2$ ( $\alpha=0$ )	1	1.5708	1.5705	1.5683	1.5492	1.4922	1.3839	1.2450	1.1085	0.8900	0.7373	0.6276	0.5462	0.4833	0.4333	0.3927
	2	3.1416	3.1391	3.1196	2.9077	2.4056	1.9635	1.6559	1.4395	1.1629	1.0000	0.8990	0.8369	0.8022	0.7877	0.7854
	3	4.7124	4.7036	4.6234	3.7679	2.9311	2.5431	2.3966	2.3586	2.3370	2.1416	1.886	1.6356	1.4494	1.2999	1.1781
	4	6.2832	6.2609	6.0096	4.4516	3.9702	3.9270	3.8866	3.8338	2.9169	2.4043	2.0584	1.8222	1.6679	1.5871	1.5708
	6	9.4248	9.3326	8.0794	7.0711	6.9368	5.8908	5.0382	4.7323	4.5732	3.8789	3.2992	2.8801	2.5839	2.4026	2.3562

WAVE GUIDE WIDTH  
RELATIVE DIELECTRIC CONSTANT  
FREE SPACE PROPAGATION CONSTANT  
NORMALIZED CUTOFF FREQUENCY  
VELOCITY OF LIGHT IN FREE SPACE  
CUTOFF FREQUENCY  
 $f = Bv/2\pi W$



$\alpha = a/W$   
 $\delta = d/W$   
 $\gamma = c/W$

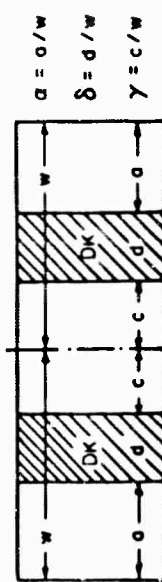
0.0500	1	1.5708	1.5670	1.5624	1.4960	1.4446										
	2	3.1416	3.1076	3.0602	2.4425	2.1691										
	3	4.7124	4.5756	4.3427	2.9923	2.7003										
	4	6.2832	5.8798	5.2476	4.0215	3.9368										
	6	9.4248	7.9353	7.2463	7.1485	6.5444										
0.1000	1	1.5708	1.5554	1.5379	1.4744	1.4218	1.1253	1.0525	0.9911							
	2	3.1416	2.6994	2.3282	1.8805	1.6338	1.4771	1.3673	1.2846							
	3	4.7124	3.6516	3.0958	2.6440	2.4797	2.4100	2.3786	2.3633							
	4	6.2832	4.8633	4.4664	4.2643	4.1448	3.9270	3.5961	3.3029							
	6	9.4248	8.2687	8.1036	7.4405	5.8838	5.0897	4.8362	4.7571							
0.2000	1	1.5708	1.5112	1.4474	1.3219	1.2138	1.1253	1.0525	0.9911							
	2	3.1416	2.6994	2.3282	1.8805	1.6338	1.4771	1.3673	1.2846							
	3	4.7124	3.6516	3.0958	2.6440	2.4797	2.4100	2.3786	2.3633							
	4	6.2832	4.8633	4.4664	4.2643	4.1448	3.9270	3.5961	3.3029							
	6	9.4248	8.2687	8.1036	7.4405	5.8838	5.0897	4.8362	4.7571							
0.3000	1	1.5708	1.4500	1.3387	1.1624	1.0381	0.9476	0.8785	0.8238	0.7409	0.6796					
	2	3.1416	2.4793	2.0739	1.6436	1.4166	1.2748	1.1766	1.1051	1.0076	0.9444					
	3	4.7124	3.5607	3.1061	2.7718	2.6512	2.5838	2.5178	2.4289	2.2026	2.0106					
	4	6.2832	5.1939	4.9269	4.7448	4.4782	3.9270	3.3959	2.9996	2.7996	2.2254					
	6	9.4248	9.2174	9.0757	7.5380	5.9476	5.4519	5.2865	5.0005	4.1618	3.6056					
0.4000	1	1.5708	1.3860	1.2417	1.0448	0.9193	0.8320	0.7673	0.7169	0.6426	0.5890	0.5474	0.5130			
	2	3.1416	2.3620	1.9532	1.5372	1.3205	1.1851	1.0920	1.0241	0.9328	0.8763	0.8402	0.8170			
	3	4.7124	3.7214	3.3634	3.0854	2.9723	2.8646	2.7036	2.4963	2.1098	1.8397	1.6610	1.5400			
	4	6.2832	5.7753	5.6005	5.3669	4.7294	3.9270	3.3289	2.9032	2.3601	2.0449	1.8545	1.7395			
	6	9.4248	8.5724	8.3064	7.3779	6.4040	6.1438	5.7555	5.1148	4.0637	3.4015	2.9847	2.7319			
0.5000	1	1.5708	1.3275	1.1627	0.9582	0.8351	0.7516	0.6904	0.6433	0.5743	0.5257	0.4886	0.4588	0.4338	0.4121	
	2	3.1416	2.3256	1.9171	1.5059	1.2925	1.1691	1.0674	1.0005	0.9106	0.8555	0.8212	0.8007	0.7901	0.7860	
	3	4.7124	4.0723	3.7922	3.5665	3.4147	3.1754	2.8504	2.5367	2.0641	1.7868	1.5516	1.4104	1.3107	1.2372	
	4	6.2832	6.2782	6.2390	5.8154	4.8109	3.9270	3.3118	2.8789	2.3257	2.0000	1.7980	1.6738	1.6044	1.5754	
	6	9.4248	7.5925	7.1986	7.0692	7.0408	6.6949	5.9108	5.1394	4.0435	3.3618	2.9194	2.6302	2.4538	2.3710	
0.6000	1	1.5708	1.2775	1.0996	0.8924	0.7723	0.6921	0.6339	0.5892	0.5244	0.4789	0.4446	0.4174	0.3950		
	3	4.7124	4.5471	4.4210	4.2273	3.8994	3.4022	2.9317	2.5575	2.0405	1.7133	1.4932	1.3386	1.2269		
0.7000	1	1.5708	1.2364	1.0494	0.8411	0.7236	0.6461	0.5902	0.5476	0.4859	0.4429	0.4106				
	3	4.7124	4.6735	4.6377	4.5048	4.0980	3.4897	2.9659	2.5672	2.0279	1.6888	1.4597				
0.8000	1	1.5708	1.2036	1.0094	0.8002	0.6848	0.6095	0.5555	0.5144	0.4553						
	3	4.7124	4.3584	4.2290	4.1051	3.9075	3.4648	2.9697	2.5702	2.0217						
1- $\delta/2$ ( $\gamma=0$ )	1	1.5708	1.1628	0.9586	0.7630	0.6462	0.5796	0.5337	0.5002	0.4653	0.4277	0.4106	0.4004	0.3950	0.3927	
	3	4.7124	3.7962	3.5993	3.5346	3.5204	3.3474	2.9554	2.5647	2.0217	1.6809	1.4597	1.3151	1.2269	1.1855	1.1781

NORMALIZED CUTOFF FREQUENCIES OF A WAVEGUIDE WHICH CONTAINS DIELECTRIC SLABS

DK = 25.00

$a + \delta/2$	$N/\delta$	0.0000	0.0250	0.0500	0.1000	0.1500	0.2000	0.2500	0.3000	0.4000	0.5000	0.6000	0.7000	0.8000	0.9000	1.0000
$\delta/2$	1	1.5708	1.5703	1.5667	1.5338	1.4321	1.2625	1.0888	0.9443	0.7381	0.6035	0.5096	0.4412	0.3888	0.3475	0.3142
( $a=0$ )	2	3.1416	3.1375	3.1030	2.6921	2.0485	1.6261	1.3552	1.1704	0.9388	0.8042	0.7213	0.6704	0.6421	0.6302	0.6283
	3	4.7124	4.6975	4.5347	3.2352	2.5064	2.2607	2.2031	2.1989	2.0847	1.9374	1.8245	1.7225	1.6662	1.6425	1.6425
	4	6.2832	6.2438	5.6606	3.9666	3.7719	3.7529	3.4994	3.0568	2.3711	1.9395	1.6546	1.4616	1.3359	1.2700	1.2566
	5	9.4248	9.2313	7.3029	6.9055	6.1368	4.9095	4.4315	4.3962	3.8381	3.1416	2.6550	2.3115	2.0703	1.9229	1.8850
0.0500	1	1.5708	1.5645	1.5568												
	2	3.1416	3.0818	2.9872												
	3	4.7124	4.4418	3.9692												
	4	6.2832	5.4560	4.5916												
	5	9.4248	7.3190	6.9656												
0.1000	1	1.5708	1.5453	1.5147	1.4396	1.3530										
	2	3.1416	2.8890	2.5637	2.0913	1.8160										
	3	4.7124	3.8068	3.1284	2.5604	2.3504										
	4	6.2832	4.6186	4.0462	3.8142	3.7767										
	5	9.4248	7.3637	7.2019	6.7200	5.5174										
0.2000	1	1.5708	1.4730	1.3693	1.1889	1.0579	0.9626	0.8897	0.8310							
	2	3.1416	2.4567	2.0085	1.5626	1.3398	1.2035	1.1099	1.0403							
	3	4.7124	3.2582	2.7353	2.3988	2.2972	2.2538	2.2234	2.1880							
	4	6.2832	4.5548	4.2752	4.1291	3.9129	3.4347	3.0078	2.7158							
	5	9.4248	8.1379	8.0001	6.4906	5.0187	4.6302	4.5211	4.3274							
0.3000	1	1.5708	1.3110	1.2235	1.0113	0.8812	0.7933	0.7292	0.6797	0.6067	0.5530					
	2	3.1416	2.2074	1.7662	1.3573	1.1564	1.0342	0.9514	0.8915	0.8106	0.7585					
	3	4.7124	3.2337	2.8363	2.5962	2.5111	2.4352	2.3144	2.1505	1.8503	1.6548					
	4	6.2832	4.9869	4.7882	4.6046	4.0324	3.3059	2.7950	2.4451	2.0200	1.7919					
	5	9.4248	8.1625	8.9391	6.4809	5.3413	5.1330	4.7991	4.2397	3.3874	2.9100					
0.4000	1	1.5708	1.2943	1.1101	0.8940	0.7707	0.6896	0.6313	0.5869	0.5226	0.4772	0.4421	0.4134			
	2	3.1416	2.0856	1.6560	1.2666	1.0762	0.9603	0.8818	0.8250	0.7493	0.7028	0.6731	0.6542			
	3	4.7124	3.4548	3.1365	2.9379	2.8308	2.6557	2.3910	2.1281	1.7400	1.4981	1.3442	1.2420			
	4	6.2832	5.6454	5.4953	5.1213	4.0838	3.2631	2.7276	2.3617	1.9055	1.6447	1.4683	1.3942			
	5	9.4248	8.3760	8.1203	6.6201	6.0265	5.6827	4.9114	4.2189	3.2939	2.7410	2.3978	2.1909			
0.5000	1	1.5708	1.2209	1.0243	0.8114	0.6946	0.6189	0.5649	0.5240	0.4652	0.4242	0.3932	0.3685	0.3479	0.3300	
	2	3.1416	2.0488	1.6235	1.2402	1.0529	0.9389	0.8615	0.8057	0.7312	0.6858	0.6576	0.6409	0.6321	0.6288	
	3	4.7124	3.8702	3.6163	3.4289	3.2173	2.8316	2.4335	2.1172	1.6885	1.4243	1.2515	1.1339	1.0514	0.9910	
	4	6.2832	6.2750	6.2060	5.3842	4.0970	3.2522	2.7104	2.3407	1.8775	1.6083	1.4425	1.3409	1.2841	1.2604	
	5	9.4248	7.2444	6.9653	5.9110	6.7411	5.8839	4.9347	4.2144	3.2743	2.7084	2.3450	2.1086	1.9647	1.8971	
0.6000	1	1.5708	1.1611	0.9584	0.7500	0.6386	0.5671	0.5164	0.4781	0.4233	0.3853	0.3569	0.3345	0.3161		
	3	4.7124	4.4641	4.3183	4.0778	3.5456	2.9267	2.4547	2.1117	1.6613	1.3850	1.2017	1.0741	0.9824		
0.7000	1	1.5708	1.1130	0.9067	0.7024	0.5954	0.5273	0.4791	0.4429	0.3912	0.3554	0.3288				
	3	4.7124	4.6554	4.5093	4.3792	3.6709	2.9653	2.4646	2.1088	1.6459	1.3620	1.1725				
0.8000	1	1.5708	1.0747	0.8655	0.6645	0.5609	0.4955	0.4494	0.4148	0.3656						
	3	4.7124	4.2652	4.1499	4.0171	3.6135	2.9696	2.4651	2.1074	1.6371						
1 - $\delta/2$	1	1.5708	1.0244	0.8118	0.6201	0.5264	0.4694	0.4308	0.4028	0.3656	0.3423	0.3288	0.3204	0.3161	0.3144	0.3142
( $\gamma=0$ )	3	4.7124	3.6222	3.4827	3.4555	3.3705	2.9444	2.4674	2.1072	1.6371	1.3842	1.1725	1.0543	0.9824	0.9485	0.9425

WAVE GUIDE WIDTH  
RELATIVE DIELECTRIC CONSTANT  
FREE SPACE PROPAGATION CONSTANT  
NORMALIZED CUTOFF FREQUENCY  
VELOCITY OF LIGHT IN FREE SPACE  
CUTOFF FREQUENCY



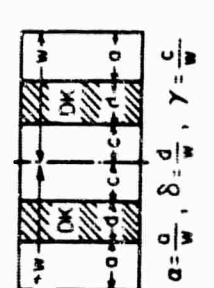
$a = 0/w$   
 $\delta = d/w$   
 $\gamma = c/w$

NORMALIZED PROPAGATION CONSTANTS FOR TE MODES IN WAVEGUIDES WHICH CONTAIN DIELECTRIC SLABS													DELTA = 0.05			
α - d/2 MODE	0.025		0.125		0.200		0.300		0.400		0.500		0.700		0.975	
	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20
BC	1.57059	3.11997	1.56697	3.11146	1.56129	3.07183	1.55051	3.01490	1.53734	2.97328	1.52329	2.95838	1.49777	1.47919		
1.5	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	0.08583	0.26477		
1.6	0.30539	-0.00000	0.32421	-0.00000	0.35200	-0.00000	0.40015	-0.00000	0.45321	-0.00000	0.50498	-0.00000	0.59041	0.64869		
1.7	0.65065	-0.00000	0.66087	-0.00000	0.67678	-0.00000	0.70643	-0.00000	0.74180	-0.00000	0.77865	-0.00000	0.84373	0.89130		
1.8	0.87945	-0.00000	0.88797	-0.00000	0.90135	-0.00000	0.92658	-0.00000	0.95714	-0.00000	0.98942	-0.00000	1.04761	1.09134		
1.9	1.06923	-0.00000	1.07720	-0.00000	1.08985	-0.00000	1.11296	-0.00000	1.14147	-0.00000	1.17174	-0.00000	1.22669	1.26912		
2.0	1.23840	-0.00000	1.24592	-0.00000	1.25780	-0.00000	1.28040	-0.00000	1.30799	-0.00000	1.33735	-0.00000	1.39082	1.43297		
2.1	1.39418	-0.00000	1.40157	-0.00000	1.41327	-0.00000	1.43555	-0.00000	1.46278	-0.00000	1.49471	-0.00000	1.54471	1.58724		
2.2	1.54073	-0.00000	1.54809	-0.00000	1.55976	-0.00000	1.59203	-0.00000	1.60924	-0.00000	1.63322	-0.00000	1.69116	1.73450		
2.3	1.68047	-0.00000	1.68785	-0.00000	1.69962	-0.00000	1.72206	-0.00000	1.74948	-0.00000	1.77865	-0.00000	1.83197	1.87643		
2.4	1.81496	-0.00000	1.82244	-0.00000	1.83435	-0.00000	1.85710	-0.00000	1.88488	-0.00000	1.91440	-0.00000	1.96834	2.01417		
2.5	1.94531	-0.00000	1.95290	-0.00000	1.96802	-0.00000	1.98818	-0.00000	2.01643	-0.00000	2.04640	-0.00000	2.10114	2.14836		
2.6	2.07226	-0.00000	2.08001	-0.00000	2.09238	-0.00000	2.11603	-0.00000	2.14485	-0.00000	2.17535	-0.00000	2.23103	2.28020		
2.7	2.19647	-0.00000	2.20436	-0.00000	2.21702	-0.00000	2.24121	-0.00000	2.27066	-0.00000	2.30176	-0.00000	2.35848	2.40956		
2.8	2.31832	-0.00000	2.32640	-0.00000	2.33936	-0.00000	2.36415	-0.00000	2.39424	-0.00000	2.42603	-0.00000	2.48387	2.53702		
2.9	2.43819	-0.00000	2.44645	-0.00000	2.45975	-0.00000	2.48519	-0.00000	2.51605	-0.00000	2.54848	-0.00000	2.60750	2.66237		
3.0	2.55635	-0.00000	2.56481	-0.00000	2.57847	-0.00000	2.60459	-0.00000	2.63622	-0.00000	2.66937	-0.00000	2.72963	2.78734		
3.1	2.67303	-0.00000	2.68171	-0.00000	2.69573	-0.00000	2.72257	-0.00000	2.75501	-0.00000	2.78891	-0.00000	2.85043	2.91004		
3.2	2.78842	0.61725	2.79731	0.75812	2.81173	0.91828	2.83932	1.12007	2.87259	1.25236	2.90727	1.29738	2.97009	3.03292		
3.3	2.90266	1.01572	2.91179	1.11074	2.92661	1.23507	2.95498	1.40142	2.98912	1.51670	3.02454	1.55631	3.08874	3.15434		
3.4	3.01589	1.30475	3.02526	1.38481	3.04050	1.49292	3.06963	1.64184	3.10472	1.74617	3.14099	1.78261	3.20649	3.27439		
3.5	3.12822	1.54697	3.13783	1.61932	3.15351	1.71858	3.18354	1.85732	3.21949	1.95542	3.25658	1.98982	3.32344	3.39500		
3.6	3.23975	1.76178	3.24961	1.82954	3.26574	1.92343	3.29663	2.05571	3.33353	2.14968	3.37145	2.18268	3.43969	3.51443		
3.7	3.35055	1.95822	3.36067	2.01106	3.37725	2.11352	3.40905	2.24157	3.44691	2.33270	3.48568	2.36470	3.55529	3.63339		
3.8	3.46069	2.14136	3.47107	2.20429	3.48815	2.29255	3.52086	2.41780	3.55970	2.50695	3.59932	2.53824	3.67033	3.75132		
3.9	3.57023	2.31435	3.58089	2.37604	3.59846	2.46293	3.63212	2.58638	3.67197	2.67415	3.71245	2.70491	3.78485	3.87009		
4.0	3.67924	2.47933	3.69017	2.54024	3.70825	2.62635	3.74288	2.74873	3.78376	2.83556	3.82517	2.86594	3.89890	3.98795		
4.1		2.63779		2.69828		2.78405		2.90589		2.99212		3.07222				
4.2		2.79085		2.85117		2.93695		3.05868		3.14456		3.17447				
4.3		2.93934		2.99971		3.08677		3.20774		3.29346		3.32325				
4.4		3.08394		3.14432		3.23109		3.35356		3.43929		3.46900				
4.5		3.22815		3.23609		3.37337		3.49656		3.58243		3.61210				
4.6		3.36342		3.42483		3.52483		3.63709		3.72319		3.75206				
4.7		3.49908		3.56107		3.65022		3.77542		3.86185		3.89153				
4.8		3.63244		3.69509		3.78537		3.91180		3.99863		4.02834				
4.9		3.76373		3.82713		3.91865		4.04643		4.13371		4.16348				
5.0		3.89317		3.95738		4.05026		4.13115		4.26728		4.29711				
5.1		4.02092		4.08602		4.18035		4.31115		4.39947		4.42938				
5.2		4.14716		4.21320		4.30908		4.44152		4.53040		4.56039				
5.3		4.27201		4.33905		4.43656		4.57072		4.66020		4.69028				
5.4		4.39560		4.46368		4.56291		4.69886		4.78896		4.81912				
5.5		4.51802		4.58721		4.68823		4.82604		4.91676		4.94702				
5.6		4.63936		4.70971		4.81260		4.95234		5.04369		5.07404				
5.7		4.75972		4.83127		4.93611		5.07782		5.16982		5.20026				
5.8		4.87916		4.95196		5.05883		5.20251		5.29521		5.32574				
5.9		4.99775		5.07185		5.16082		5.32663		5.41992		5.45053				
6.0		5.11555		5.19099		5.30213		5.45006		5.54399		5.57468				
6.1		5.23261		5.30942		5.42282		5.57292		5.66749		5.69825				
6.2		5.34898		5.42725		5.54296		5.69524		5.79044		5.82128				
6.3		5.46470		5.54445		5.66257		5.81708		5.91290		5.94380				

$a = \frac{d}{W}, \delta = \frac{d}{W}, \gamma = \frac{c}{W}$

NORMALIZED PROPAGATION CONSTANTS FOR TE MODES IN WAVEGUIDES WHICH CONTAIN DIELECTRIC SLABS									
$\alpha \cdot d/2$	MODE	CONSTANTS FOR TE MODES IN		WAVEGUIDES WHICH CONTAIN		DIELECTRIC SLABS		DK = 2.25	
		0.050	0.125	0.200	0.300	0.400	0.500	0.500	0.500
MC	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20	TE 10
		TE 20	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20	TE 10
1.5	1.56917	3.12847	1.56281	3.07807	1.55141	3.00000	1.53016	2.89616	1.50496
1.6	0.31288	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.7	0.65467	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.8	0.89279	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.9	1.07243	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
2.0	1.24134	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
2.1	1.39707	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
2.2	1.54360	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
2.3	1.68335	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
2.4	1.81788	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
2.5	1.94827	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
2.6	2.07529	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
2.7	2.19955	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
2.8	2.32147	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
2.9	2.44140	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
3.0	2.55564	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
3.1	2.67640	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
3.2	2.79187	0.67578	2.80836	0.89482	2.83961	1.17248	2.89490	1.48725	2.96686
3.3	2.90620	1.05471	2.92319	1.21698	2.95553	1.44807	3.01688	1.72946	3.08682
3.4	3.01952	1.37624	3.03704	1.47741	3.07052	1.68598	3.13403	1.94827	3.20601
3.5	3.13194	1.57624	3.15001	1.70464	3.18469	1.90043	3.25046	2.15086	3.32484
3.6	3.24356	1.78911	3.26219	1.91059	3.29813	2.09867	3.36627	2.34148	3.44280
3.7	3.35446	1.98432	3.37368	2.10152	3.41094	2.28495	3.48184	2.52289	3.55998
3.8	3.46470	2.16664	3.48462	2.28122	3.52317	2.46201	3.59635	2.69697	3.67705
3.9	3.57434	2.33910	3.59480	2.45217	3.63489	2.63175	3.71077	2.86509	3.79370
4.0	3.68345	2.50374	3.70456	2.61611	3.74617	2.79584	3.82486	3.02828	3.91021
4.1	2.66200	2.81496	2.92783	2.77428	2.95439	2.98439	3.16732	3.16732	3.33317
4.2	2.81496	2.96345	3.07591	2.92783	3.10910	3.10910	3.34285	3.34285	3.48911
4.3	3.01810	3.10810	3.22267	3.07591	3.26027	3.26027	3.49535	3.49535	3.64024
4.4	3.24944	3.24944	3.36540	3.22267	3.40842	3.40842	3.64523	3.64523	3.78991
4.5	3.38787	3.38787	3.50548	3.36540	3.55396	3.55396	3.79284	3.79284	3.93744
4.6	3.52374	3.52374	3.64322	3.50548	3.69720	3.69720	3.93847	3.93847	4.08307
4.7	3.65733	3.65733	3.77890	3.64322	3.83847	3.83847	4.08238	4.08238	4.22701
4.8	3.78889	3.78889	3.91274	3.77890	3.97799	3.97799	4.22468	4.22468	4.36946
4.9	3.91862	3.91862	4.04495	3.91274	4.11598	4.11598	4.36566	4.36566	4.51057
5.0	4.04670	4.04670	4.17568	4.04495	4.25262	4.25262	4.50542	4.50542	4.65049
5.1	4.17329	4.17329	4.30809	4.17568	4.38806	4.38806	4.64413	4.64413	4.78933
5.2	4.29851	4.29851	4.43531	4.30809	4.52246	4.52246	4.78188	4.78188	4.92721
5.3	4.42248	4.42248	4.56046	4.43531	4.65594	4.65594	4.91879	4.91879	5.06422
5.4	4.54530	4.54530	4.68664	4.56046	4.78861	4.78861	5.05196	5.05196	5.20046
5.5	4.66707	4.66707	4.81194	4.68664	4.92087	4.92087	5.19046	5.19046	5.33899
5.6	4.78787	4.78787	4.93646	4.81194	5.05192	5.05192	5.32539	5.32539	5.47089
5.7	4.90777	4.90777	5.06026	4.93646	5.18274	5.18274	5.45979	5.45979	5.60622
5.8	5.02683	5.02683	5.18341	5.06026	5.31311	5.31311	5.59375	5.59375	5.73904
5.9	5.14512	5.14512	5.30599	5.18341	5.44310	5.44310	5.72731	5.72731	5.87240
6.0	5.26269	5.26269	5.42806	5.30599	5.57278	5.57278	5.86052	5.86052	6.00838
6.1	5.37958	5.37958	5.54966	5.42806	5.70221	5.70221	5.99344	5.99344	6.13793
6.2	5.49584	5.49584	5.67076	5.54966	5.83148	5.83148	6.12611	6.12611	6.27018
6.3				5.67076	5.96084	5.96084	6.25856	6.25856	6.40218

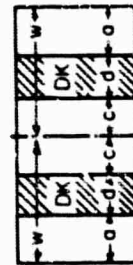
2W GUIDE WIDTH  
 DK RELATIVE DIELEC-  
 TRIC CONSTANT  
 $\lambda_0$  FREE SPACE WAVE  
 $\lambda_g$  GUIDE WAVELENGTH  
 $B = 2\pi w/\lambda_0$   
 NORMALIZED FREQ.  
 BC CUTOFF-FREQ.  
 $K = 2\pi/\lambda_g$   
 NORMALIZED  
 PROPAGATION  
 CONSTANT



NORMALIZED PROPAGATION CONSTANTS FOR TE MODES IN									
MODE	0.075		0.125		0.200		0.300		DIELECTRIC SLABS
	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20	
BC	1.56531	3.09723	1.55805	3.04078	1.54101	2.92867	1.50980	2.78926	1.47589
B	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.4	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.5	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.6	0.33253	-0.00000	0.36707	-0.00000	0.43920	-0.00000	0.55226	-0.00000	0.66557
1.7	0.66555	-0.00000	0.68586	-0.00000	0.73253	-0.00000	0.81499	-0.00000	0.90597
1.8	0.89191	-0.00000	0.90910	-0.00000	0.94939	-0.00000	1.02259	-0.00000	1.10551
1.9	1.08083	-0.00000	1.09679	-0.00000	1.13453	-0.00000	1.20326	-0.00000	1.28345
2.0	1.24943	-0.00000	1.26486	-0.00000	1.30158	-0.00000	1.36950	-0.00000	1.44783
2.1	1.40503	-0.00000	1.42028	-0.00000	1.45678	-0.00000	1.52469	-0.00000	1.60286
2.2	1.55155	-0.00000	1.56686	-0.00000	1.60361	-0.00000	1.67220	-0.00000	1.75109
2.3	1.69136	-0.00000	1.70684	-0.00000	1.74419	-0.00000	1.81402	-0.00000	1.89415
2.4	1.82600	-0.00000	1.84175	-0.00000	1.87993	-0.00000	1.95142	-0.00000	2.03319
2.5	1.95653	-0.00000	1.97264	-0.00000	2.01185	-0.00000	2.08533	-0.00000	2.16902
2.6	2.08374	-0.00000	2.10027	-0.00000	2.14066	-0.00000	2.21642	-0.00000	2.30227
2.7	2.20819	-0.00000	2.22519	-0.00000	2.26691	-0.00000	2.34519	-0.00000	2.43340
2.8	2.33033	-0.00000	2.34785	-0.00000	2.39103	-0.00000	2.47206	-0.00000	2.56279
2.9	2.45051	-0.00000	2.46859	-0.00000	2.51334	-0.00000	2.59734	-0.00000	2.69074
3.0	2.56900	-0.00000	2.58767	-0.00000	2.63413	-0.00000	2.72130	-0.00000	2.81750
3.1	2.68603	0.13303	2.70534	0.62560	2.75362	1.10142	2.84417	1.54318	2.94327
3.2	2.80178	0.81704	2.82176	1.03445	2.87199	1.39856	2.96612	1.79073	3.06323
3.3	2.91640	1.15685	2.93710	1.33087	2.98939	1.65069	3.08732	2.01525	3.19252
3.4	3.03003	1.42475	3.05148	1.57952	3.10597	1.85612	3.20790	2.22378	3.31628
3.5	3.14277	1.65604	3.16501	1.80033	3.22183	2.08359	3.32800	2.42053	3.43962
3.6	3.25472	1.86440	3.27778	2.00257	3.33709	2.27810	3.44771	2.60823	3.56264
3.7	3.36595	2.05621	3.38989	2.19141	3.45183	2.46278	3.56713	2.78876	3.68543
3.8	3.47655	2.23740	3.50140	2.37008	3.56613	2.63976	3.68637	2.96349	3.80807
3.9	3.58656	2.40884	3.61237	2.54076	3.68007	2.81056	3.80551	3.13344	3.93063
4.0	3.69604	2.57297	3.72286	2.70499	3.79372	2.97632	3.92461	3.29941	4.05319
4.1	2.73111	2.86332	2.86332	2.86332	3.13789	3.13789	3.45203	3.45203	3.64530
4.2	2.86425	3.01840	3.01840	3.01840	3.29596	3.29596	3.62179	3.62179	3.80393
4.3	3.03316	3.16912	3.16912	3.16912	3.45107	3.45107	3.77911	3.77911	3.96027
4.4	3.17844	3.31653	3.31653	3.31653	3.60369	3.60369	3.93432	3.93432	4.11460
4.5	3.32058	3.45136	3.45136	3.45136	3.75417	3.75417	4.08773	4.08773	4.26714
4.6	3.45997	3.60369	3.60369	3.60369	3.90283	3.90283	4.23956	4.23956	4.41822
4.7	3.59693	3.74392	3.74392	3.74392	4.04994	4.04994	4.39003	4.39003	4.56789
4.8	3.73175	3.88231	3.88231	3.88231	4.19574	4.19574	4.53932	4.53932	4.71637
4.9	3.86465	4.01910	4.01910	4.01910	4.34042	4.34042	4.68758	4.68758	4.86378
5.0	3.99804	4.15448	4.15448	4.15448	4.48417	4.48417	4.83494	4.83494	5.01024
5.1	4.12849	4.28862	4.28862	4.28862	4.62713	4.62713	4.98194	4.98194	5.15586
5.2	4.25374	4.42167	4.42167	4.42167	4.76946	4.76946	5.12747	5.12747	5.30075
5.3	4.38073	4.55378	4.55378	4.55378	4.91126	4.91126	5.27282	5.27282	5.44497
5.4	4.50658	4.68507	4.68507	4.68507	5.05269	5.05269	5.41768	5.41768	5.58862
5.5	4.63139	4.81565	4.81565	4.81565	5.19380	5.19380	5.56212	5.56212	5.73175
5.6	4.75525	4.94563	4.94563	4.94563	5.33471	5.33471	5.70620	5.70620	5.87443
5.7	4.87824	5.07510	5.07510	5.07510	5.47550	5.47550	5.84999	5.84999	6.01671
5.8	5.00045	5.20416	5.20416	5.20416	5.61624	5.61624	5.99353	5.99353	6.15865
5.9	5.12193	5.33289	5.33289	5.33289	5.75700	5.75700	6.13688	6.13688	6.30028
6.0	5.24275	5.46138	5.46138	5.46138	5.89784	5.89784	6.28006	6.28006	6.44166
6.1	5.36297	5.58969	5.58969	5.58969	6.03882	6.03882	6.42312	6.42312	6.58282
6.2	5.48265	5.71791	5.71791	5.71791	6.17998	6.17998	6.56610	6.56610	6.72379
6.3	5.60183	5.84611	5.84611	5.84611	6.32138	6.32138	6.70901	6.70901	6.86460

2W GUIDE WIDTH  
DK RELATIVE DIELEC-  
TRIC CONSTANT  
 $\lambda_0$  FREE SPACE WAVE  
 $\lambda_g$  GUIDE WAVELENGTH  
 $B = 2\pi w/\lambda_0$   
NORMALIZED FREQ.  
BC CUTOFF-FREQ.

$K = 2\pi w/\lambda_g$   
NORMALIZED  
PROPAGATION  
CONSTANT



$$\alpha = \frac{a}{w}, \delta = \frac{d}{w}, \gamma = \frac{c}{w}$$



NORMALIZED PROPAGATION CONSTANTS FOR TE MODES IN WAVEGUIDES WHICH CONTAIN DIELECTRIC SLABS									
$\alpha \cdot d/2$	MODE	TE 10		TE 20		TE 10		TE 20	
		TE 10	TE 20	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20
BC		1.54571	2.95240	1.51811	2.79424	1.46926	2.61369	1.41591	2.60711
DELTA = 0.25									
1.3		-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.4		-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.5		-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.6		0.42034	-0.00000	0.52414	-0.00000	0.68081	-0.00000	0.83215	-0.00000
1.7		0.71986	-0.00000	0.79386	-0.00000	0.91972	-0.00000	1.05164	-0.00000
1.8		0.93839	-0.00000	1.00389	-0.00000	1.11924	-0.00000	1.24361	-0.00000
1.9		1.12422	-0.00000	1.18642	-0.00000	1.29775	-0.00000	1.41903	-0.00000
2.0		1.29156	-0.00000	1.35267	-0.00000	1.46308	-0.00000	1.58377	-0.00000
2.1		1.44685	-0.00000	1.50809	-0.00000	1.61935	-0.00000	1.74092	-0.00000
2.2		1.59363	-0.00000	1.65577	-0.00000	1.76907	-0.00000	1.89246	-0.00000
2.3		1.73409	-0.00000	1.79766	-0.00000	1.91388	-0.00000	2.03977	-0.00000
2.4		1.86964	-0.00000	1.93610	-0.00000	2.05491	-0.00000	2.18379	-0.00000
2.5		2.00132	-0.00000	2.06902	-0.00000	2.19300	-0.00000	2.32526	-0.00000
2.6		2.12986	-0.00000	2.20012	-0.00000	2.32878	-0.00000	2.46471	-0.00000
2.7		2.25581	-0.00000	2.32893	-0.00000	2.46274	-0.00000	2.60259	-0.00000
2.8		2.37960	-0.00000	2.45587	-0.00000	2.59628	-0.00000	2.73924	-0.00000
2.9		2.50156	-0.00000	2.58127	-0.00000	2.72673	-0.00000	2.87495	-0.00000
3.0		2.62197	0.57190	2.70542	1.28188	2.85737	1.80703	3.00996	2.09166
3.1		2.74106	1.01622	2.82855	1.34100	2.98743	2.04828	3.14446	2.31672
3.2		2.85901	1.32804	2.95087	1.79286	3.11711	2.27186	3.27864	2.52890
3.3		2.97598	1.58756	3.07255	2.02140	3.24661	2.48263	3.41265	2.73123
3.4		3.09210	1.81739	3.19376	2.23391	3.37609	2.68373	3.54661	2.92523
3.5		3.20749	2.02776	3.31463	2.43474	3.50671	2.87727	3.68064	3.11413
3.6		3.32226	2.22424	3.43530	2.62668	3.63659	3.06476	3.81485	3.29741
3.7		3.43648	2.41034	3.55591	2.81168	3.76688	3.24732	3.94931	3.47636
3.8		3.55026	2.58836	3.67657	2.99116	3.89664	3.42883	4.08412	3.65171
3.9		3.66366	2.75994	3.79741	3.16618	4.02813	3.60098	4.21932	3.82399
4.0		3.77674	2.92630	3.91854	3.33754	4.16030	3.77322	4.35499	4.00000
4.1		3.88836	3.08836	4.03961	3.50691	4.29309	3.94309	4.49102	4.16102
4.2		3.99884	3.24684	4.16181	3.67181	4.41090	4.11090	4.62642	4.32642
4.3		4.10733	3.40233	4.28567	3.83567	4.52795	4.27695	4.76101	4.49010
4.4		4.21444	3.55630	4.40978	3.99785	4.64150	4.44150	4.89627	4.65227
4.5		4.32015	3.70615	4.53415	4.15865	4.75475	4.60475	5.03132	4.81312
4.6		4.42422	3.85822	4.65822	4.31833	4.86888	4.76888	5.16144	4.97279
4.7		4.52726	4.00280	4.78111	4.47711	4.98208	4.88208	5.29144	5.13144
4.8		4.62913	4.14913	4.90317	4.63517	5.09538	5.09538	5.41917	5.28917
4.9		4.72945	4.29445	5.02267	4.79267	5.20799	5.20799	5.54610	5.44610
5.0		4.82894	4.43894	5.14977	4.94977	5.31938	5.31938	5.69231	5.60231
5.1		4.92679	4.58279	5.27656	5.10656	5.43042	5.43042	5.83789	5.75789
5.2		5.02314	4.72614	5.40317	5.26317	5.54097	5.54097	5.98293	5.91292
5.3		5.11895	4.86918	5.52918	5.41966	5.65097	5.65097	6.12744	6.06744
5.4		5.21466	5.01195	5.65612	5.57612	5.76042	5.76042	6.27153	6.21653
5.5		5.30938	5.15466	5.78262	5.73262	5.86917	5.86917	6.41511	6.36011
5.6		5.40319	5.29738	5.90819	5.88919	5.97744	5.97744	6.55859	6.50359
5.7		5.49622	5.44022	6.03588	6.04588	6.08616	6.08616	6.70361	6.64861
5.8		5.58838	5.58328	6.16273	6.20273	6.23444	6.23444	6.84844	6.79344
5.9		5.67936	5.72664	6.28936	6.35976	6.39701	6.39701	6.99301	6.93801
6.0		5.76936	5.87036	6.41656	6.51656	7.13937	7.13937	7.13937	7.13937
6.1		5.85936	6.01464	6.54444	6.67444	7.29166	7.29166	7.29166	7.29166
6.2		5.94936	6.15922	6.67208	6.83208	7.44360	7.44360	7.44360	7.44360
6.3		6.03936	6.30446	6.80996	6.98996	7.59582	7.59582	7.59582	7.59582

2W GUIDE WIDTH

DK RELATIVE DIELECTRIC CONSTANT

$\lambda_0$  FREE SPACE WAVELENGTH

$\lambda_g$  GUIDE WAVELENGTH

$B = 2\pi w/\lambda_0$

NORMALIZED FREQ.

BC CUTOFF-FREQ.

$K = 2\pi w/\lambda_g$

NORMALIZED PROPAGATION CONSTANT

$\alpha = \frac{\pi}{w}, \delta = \frac{d}{w}, \gamma = \frac{c}{w}$

# $\alpha + \delta/2$ NORMALIZED PROPAGATION CONSTANTS FOR TE MODES IN WAVEGUIDES WHICH CONTAIN DIELECTRIC SLABS DK = 2.25 0.500 0.400 0.300 0.200 0.100 0.000 TE 10 TE 20 TE 10 TE 20 TE 10 TE 20 BC 1.47680 2.61799 1.40918 2.42967 1.33917 2.31824 1.27565 2.28240 1.17699 1.14120 B 0.31266 0.51303 0.73820 0.86156 1.01415 1.12308 1.24442 1.34928 1.45097 1.55677 1.64282 1.74963 1.82454 1.93476 1.99923 2.11353 2.16837 2.28752 2.33327 2.45782 2.49482 2.62521 2.65366 2.79027 2.81029 2.95346 3.11511 3.11836 3.27552 3.27035 3.43488 3.42126 3.59338 3.57126 3.75117 3.72047 3.90836 3.86903 4.06505 4.01703 4.22132 4.16455 4.37723 4.31168 4.53285 4.45846 4.68622 4.60496 4.84337 4.75124 4.99834 4.89752 5.15316 5.04325 5.30785 5.18906 5.46242 4.74616 4.91564 5.08367 5.25042 5.41600 5.58054 5.74415 5.90392 6.06892 6.23023 6.39091 6.55103 6.71063 6.86976 7.02846 7.18676 7.34471 7.50234 7.65966 7.81672 7.97352 8.13009 8.28646

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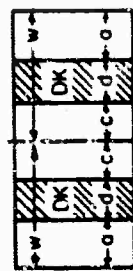
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NORMALIZED PROPAGATION CONSTANTS FOR TE MODES IN WAVEGUIDES WHICH CONTAIN DIELECTRIC SLABS									
$\alpha \cdot d/2$	$0.025$	$0.125$	$0.225$	$0.300$	$0.400$	$0.500$	$0.600$	$0.700$	$0.800$
$\beta \cdot d/2$	$0.025$	$0.125$	$0.225$	$0.300$	$0.400$	$0.500$	$0.600$	$0.700$	$0.800$
BC	1.57031	1.56149	1.54768	1.52212	1.49224	1.46188	1.42705	1.39222	1.35739
1.5	0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.6	0.30650	-0.00000	0.35108	0.41219	0.48972	0.58072	0.69097	0.82491	0.98498
1.7	0.65145	-0.00000	0.76249	0.91444	1.08275	1.26364	1.45854	1.66999	1.89899
1.8	0.88012	-0.00000	0.99099	1.23366	1.49364	1.76364	2.04364	2.33364	2.63364
1.9	1.05997	-0.00000	1.08926	1.31975	1.57633	1.84633	2.12633	2.41633	2.71633
2.0	1.23999	-0.00000	1.26787	1.52717	1.80442	2.08442	2.36442	2.64442	2.92442
2.1	1.39476	-0.00000	1.41309	1.64245	1.93446	2.21446	2.49446	2.77446	3.05446
2.2	1.54131	-0.00000	1.55964	1.78916	2.14461	2.42461	2.70461	2.98461	3.26461
2.3	1.68104	-0.00000	1.69937	1.92950	2.27476	2.55476	2.83476	3.11476	3.39476
2.4	1.81555	-0.00000	1.83388	2.06491	2.39500	2.67500	2.95500	3.23500	3.51500
2.5	1.94590	-0.00000	1.96423	2.24526	2.57529	2.85529	3.13529	3.41529	3.69529
2.6	2.07288	-0.00000	2.09121	2.42561	2.75564	3.03564	3.31564	3.59564	3.87564
2.7	2.19709	-0.00000	2.21542	2.60596	2.93599	3.21599	3.49599	3.77599	4.05599
2.8	2.31895	-0.00000	2.33728	2.78631	3.11634	3.39634	3.67634	3.95634	4.23634
2.9	2.43883	-0.00000	2.45716	2.96666	3.29669	3.57669	3.85669	4.13669	4.41669
3.0	2.55701	-0.00000	2.57534	3.14701	3.47704	3.75704	4.03704	4.31704	4.59704
3.1	2.67370	-0.00000	2.69203	3.32736	3.65739	3.93739	4.21739	4.49739	4.77739
3.2	2.78910	0.62950	2.80743	3.50771	3.83774	4.11774	4.39774	4.67774	4.95774
3.3	2.90336	1.02370	2.92169	3.68806	4.01809	4.29809	4.57809	4.85809	5.13809
3.4	3.01661	1.31138	3.03494	3.86841	4.19844	4.47844	4.75844	5.03844	5.31844
3.5	3.12896	1.55291	3.14724	4.04876	4.37876	4.65876	4.93876	5.21876	5.49876
3.6	3.24081	1.76732	3.25860	4.22909	4.55909	4.83909	5.11909	5.39909	5.67909
3.7	3.35132	1.93390	3.36919	4.40942	4.73942	5.01942	5.29942	5.57942	5.85942
3.8	3.46148	2.14647	3.47934	4.58975	4.91975	5.19975	5.47975	5.75975	6.03975
3.9	3.57105	2.31935	3.58892	4.77008	5.10008	5.38008	5.66008	5.94008	6.22008
4.0	3.68067	2.48245	3.69852	4.95041	5.28041	5.56041	5.84041	6.12041	6.40041
4.1		2.64267	3.80113	5.13074	5.46074	5.74074	6.02074	6.30074	6.58074
4.2		2.79870	3.96124	5.31107	5.64107	5.92107	6.20107	6.48107	6.76107
4.3		2.94419	4.11140	5.49140	5.82140	6.10140	6.38140	6.66140	6.94140
4.4		3.08879	4.25819	5.67173	6.00173	6.28173	6.56173	6.84173	7.12173
4.5		3.23063	4.40208	5.85206	6.18206	6.46206	6.74206	7.02206	7.30206
4.6		3.36832	4.54343	6.03239	6.36239	6.64239	6.92239	7.20239	7.48239
4.7		3.50402	4.68258	6.21272	6.54272	6.82272	7.10272	7.38272	7.66272
4.8		3.63742	4.81980	6.39305	6.72305	7.00305	7.28305	7.56305	7.84305
4.9		3.76876	4.95531	6.57338	6.90338	7.18338	7.46338	7.74338	8.02338
5.0		3.89828	5.09000	6.75371	7.08371	7.36371	7.64371	7.92371	8.20371
5.1		4.02607	5.22400	6.93404	7.26404	7.54404	7.82404	8.10404	8.38404
5.2		4.15237	5.35832	7.11437	7.44437	7.72437	8.00437	8.28437	8.56437
5.3		4.27729	5.49200	7.29470	7.62470	7.90470	8.18470	8.46470	8.74470
5.4		4.40095	5.62530	7.47503	7.80503	8.08503	8.36503	8.64503	8.92503
5.5		4.52344	5.75860	7.65536	7.98536	8.26536	8.54536	8.82536	9.10536
5.6		4.64486	5.89190	7.83569	8.16569	8.44569	8.72569	9.00569	9.28569
5.7		4.76530	6.02520	8.01602	8.34602	8.62602	8.90602	9.18602	9.46602
5.8		4.88574	6.15850	8.19635	8.52635	8.80635	9.08635	9.36635	9.64635
5.9		5.00618	6.29180	8.37668	8.70668	8.98668	9.26668	9.54668	9.82668
6.0		5.12662	6.42510	8.55701	8.88701	9.16701	9.44701	9.72701	10.00701
6.1		5.24706	6.55840	8.73734	9.06734	9.34734	9.62734	9.90734	10.18734
6.2		5.36750	6.69170	8.91767	9.24767	9.52767	9.80767	10.08767	10.36767
6.3		5.48794	6.82500	9.09800	9.42800	9.70800	9.98800	10.26800	10.54800

NORMALIZED PROPAGATION CONSTANTS FOR TE MODES IN WAVEGUIDES WHICH CONTAIN DIELECTRIC SLABS									
$\alpha \cdot d/2$	0.025	0.125	0.225	0.300	0.400	0.500	0.600	0.700	0.800
$\beta \cdot d/2$	0.025	0.125	0.225	0.300	0.400	0.500	0.600	0.700	0.800
BC	1.57031	1.56149	1.54768	1.52212	1.49224	1.46188	1.42705	1.39222	1.35739
1.5	0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.6	0.30650	-0.00000	0.35108	0.41219	0.48972	0.58072	0.69097	0.82491	0.98498
1.7	0.65145	-0.00000	0.76249	0.91444	1.08275	1.26364	1.45854	1.66999	1.89899
1.8	0.88012	-0.00000	0.99099	1.23366	1.49364	1.76364	2.04364	2.33364	2.63364
1.9	1.05997	-0.00000	1.08926	1.31975	1.57633	1.84633	2.12633	2.41633	2.71633
2.0	1.23999	-0.00000	1.26787	1.52717	1.80442	2.08442	2.36442	2.64442	2.92442
2.1	1.39476	-0.00000	1.41309	1.64245	1.93446	2.21446	2.49446	2.77446	3.05446
2.2	1.54131	-0.00000	1.55964	1.78916	2.14461	2.42461	2.70461	2.98461	3.26461
2.3	1.68104	-0.00000	1.69937	1.92950	2.27476	2.55476	2.83476	3.11476	3.39476
2.4	1.81555	-0.00000	1.83388	2.06491	2.39500	2.67500	2.95500	3.23500	3.51500
2.5	1.94590	-0.00000	1.96423	2.24526	2.57529	2.85529	3.13529	3.41529	3.69529
2.6	2.07288	-0.00000	2.09121	2.42561	2.75564	3.03564	3.31564	3.59564	3.87564
2.7	2.19709	-0.00000	2.21542	2.60596	2.93599	3.21599	3.49599	3.77599	4.05599
2.8	2.31895	-0.00000	2.33728	2.78631	3.11634	3.39634	3.67634	3.95634	4.23634
2.9	2.43883	-0.00000	2.45716	2.96666	3.29669	3.57669	3.85669	4.13669	4.41669
3.0	2.55701	-0.00000	2.57534	3.14701	3.47704	3.75704	4.03704	4.31704	4.59704
3.1	2.67370	-0.00000	2.69203	3.32736	3.65739	3.93739	4.21739	4.49739	4.77739
3.2	2.78910	0.62950	2.80743	3.50771	3.83774	4.11774	4.39774	4.67774	4.95774
3.3	2.90336	1.02370	2.92169	3.68806	4.01809	4.29809	4.57809	4.85809	5.13809
3.4	3.01661	1.31138	3.03494	3.86841	4.19844	4.47844	4.75844	5.03844	5.31844
3.5	3.12896	1.55291	3.14724	4.04876	4.37876	4.65876	4.93876	5.21876	5.49876
3.6	3.24081	1.76732	3.25860	4.22909	4.55909	4.83909	5.11909	5.39909	5.67909
3.7	3.35132	1.93390	3.36919	4.40942	4.73942	5.01942	5.29942	5.57942	5.85942
3.8	3.46148	2.14647	3.47934	4.58975	4.91975	5.19975	5.47975	5.75975	6.03975
3.9	3.57105	2.31935	3.58892	4.77008	5.10008	5.38008	5.66008	5.94008	6.22008
4.0	3.68067	2.48245	3.69852	4.95041	5.28041	5.56041	5.84041	6.12041	6.40041
4.1		2.64267	3.80113	5.13074	5.46074	5.74074	6.02074	6.30074	6.58074
4.2		2.79870	3.96124	5.31107	5.64107	5.92107	6.20107	6.48107	6.76107
4.3		2.94419	4.11140	5.49140	5.82140	6.10140	6.38140	6.66140	6.94140
4.4		3.08879	4.25819	5.67173	6.00173	6.28173	6.56173	6.84173	7.12173
4.5		3.23063	4.40208	5.85206	6.18206	6.46206	6.74206	7.02206	7.30206
4.6		3.36832	4.54343	6.03239	6.36239	6.64239	6.92239	7.20239	7.48239
4.7		3.50402	4.68258	6.21272	6.54272	6.82272	7.10272	7.38272	7.66272
4.8		3.63742	4.81980	6.39305	6.72305	7.00305	7.28305	7.56305	7.84305
4.9		3.76876	4.95531	6.57338	6.90338	7.18338	7.46338	7.74338	8.02338
5.0		3.89828	5.09000	6.75371	7.08371	7.36371	7.64371	7.92371	8.20371
5.1		4.02607	5.22400	6.93404	7.26404	7.54404	7.82404	8.10404	8.38404
5.2		4.15237	5.35832	7.11437	7.44437	7.72437	8.00437	8.28437	8.56437
5.3		4.27729	5.49200	7.29470	7.62470	7.90470	8.18470	8.46470	8.74470
5.4		4.40095	5.62530	7.47503	7.80503	8.08503	8.36503	8.64503	8.92503
5.5		4.52344	5.75860	7.65536	7.98536	8.26536	8.54536	8.82536	9.10536
5.6		4.64486	5.89190	7.83569	8.16569	8.44569	8.72569	9.00569	9.28569
5.7		4.76530	6.02520	8.01602	8.34602	8.62602	8.90602	9.18602	9.46602
5.8		4.88574	6.15850	8.19635	8.52635	8.80635	9.08635	9.36635	9.64635
5.9		5.00618	6.29180	8.37668	8.70668	8.98668	9.26668	9.54668	9.82668
6.0		5.12662	6.42510	8.55701	8.88701	9.16701	9.44701	9.72701	10.00701



MODE	NORMALIZED PROPAGATION CONSTANTS FOR TE MODES IN WAVEGUIDES WHICH CONTAIN DIELECTRIC SLABS				DIELECTRIC SLABS				DK				DELTA			
	0.050	0.125	0.200	0.300	0.400	0.500	0.600	0.700	0.800	0.900	1.000	1.100	0.700	0.800	0.900	1.000
BC	1.56485	1.55115	1.52333	1.47410	1.42056	1.35994	1.28868	1.20868	1.12017	1.02610	0.92610	0.82010	0.700	0.500	0.300	0.100
1.4	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.5	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.6	0.32483	0.33763	0.35043	0.36323	0.37603	0.38883	0.40163	0.41443	0.42723	0.44003	0.45283	0.46563	0.47843	0.49123	0.50403	0.51683
1.7	0.66123	0.67403	0.68683	0.69963	0.71243	0.72523	0.73803	0.75083	0.76363	0.77643	0.78923	0.80203	0.81483	0.82763	0.84043	0.85323
1.8	0.88929	0.90209	0.91489	0.92769	0.94049	0.95329	0.96609	0.97889	0.99169	1.00449	1.01729	1.03009	1.04289	1.05569	1.06849	1.08129
1.9	1.07750	1.09030	1.10310	1.11590	1.12870	1.14150	1.15430	1.16710	1.17990	1.19270	1.20550	1.21830	1.23110	1.24390	1.25670	1.26950
2.0	1.24623	1.25903	1.27183	1.28463	1.29743	1.31023	1.32303	1.33583	1.34863	1.36143	1.37423	1.38703	1.40000	1.41280	1.42560	1.43840
2.1	1.40189	1.41469	1.42749	1.44029	1.45309	1.46589	1.47869	1.49149	1.50429	1.51709	1.52989	1.54269	1.55549	1.56829	1.58109	1.59389
2.2	1.54843	1.56123	1.57403	1.58683	1.59963	1.61243	1.62523	1.63803	1.65083	1.66363	1.67643	1.68923	1.70203	1.71483	1.72763	1.74043
2.3	1.68822	1.70102	1.71382	1.72662	1.73942	1.75222	1.76502	1.77782	1.79062	1.80342	1.81622	1.82902	1.84182	1.85462	1.86742	1.88022
2.4	1.82282	1.83562	1.84842	1.86122	1.87402	1.88682	1.89962	1.91242	1.92522	1.93802	1.95082	1.96362	1.97642	1.98922	2.00202	2.01482
2.5	1.95331	1.96611	1.97891	1.99171	2.00451	2.01731	2.03011	2.04291	2.05571	2.06851	2.08131	2.09411	2.10691	2.11971	2.13251	2.14531
2.6	2.08046	2.09326	2.10606	2.11886	2.13166	2.14446	2.15726	2.17006	2.18286	2.19566	2.20846	2.22126	2.23406	2.24686	2.25966	2.27246
2.7	2.20484	2.21764	2.23044	2.24324	2.25604	2.26884	2.28164	2.29444	2.30724	2.32004	2.33284	2.34564	2.35844	2.37124	2.38404	2.39684
2.8	2.32691	2.33971	2.35251	2.36531	2.37811	2.39091	2.40371	2.41651	2.42931	2.44211	2.45491	2.46771	2.48051	2.49331	2.50611	2.51891
2.9	2.44701	2.45981	2.47261	2.48541	2.49821	2.51101	2.52381	2.53661	2.54941	2.56221	2.57501	2.58781	2.60061	2.61341	2.62621	2.63901
3.0	2.56542	2.57822	2.59102	2.60382	2.61662	2.62942	2.64222	2.65502	2.66782	2.68062	2.69342	2.70622	2.71902	2.73182	2.74462	2.75742
3.1	2.68236	2.69516	2.70796	2.72076	2.73356	2.74636	2.75916	2.77196	2.78476	2.79756	2.81036	2.82316	2.83596	2.84876	2.86156	2.87436
3.2	2.75601	2.76881	2.78161	2.79441	2.80721	2.82001	2.83281	2.84561	2.85841	2.87121	2.88401	2.89681	2.90961	2.92241	2.93521	2.94801
3.3	2.91254	2.92534	2.93814	2.95094	2.96374	2.97654	2.98934	3.00214	3.01494	3.02774	3.04054	3.05334	3.06614	3.07894	3.09174	3.10454
3.4	3.02807	3.04087	3.05367	3.06647	3.07927	3.09207	3.10487	3.11767	3.13047	3.14327	3.15607	3.16887	3.18167	3.19447	3.20727	3.22007
3.5	3.13872	3.15152	3.16432	3.17712	3.18992	3.20272	3.21552	3.22832	3.24112	3.25392	3.26672	3.27952	3.29232	3.30512	3.31792	3.33072
3.6	3.25055	3.26335	3.27615	3.28895	3.30175	3.31455	3.32735	3.34015	3.35295	3.36575	3.37855	3.39135	3.40415	3.41695	3.42975	3.44255
3.7	3.36169	3.37449	3.38729	3.40009	3.41289	3.42569	3.43849	3.45129	3.46409	3.47689	3.48969	3.50249	3.51529	3.52809	3.54089	3.55369
3.8	3.47218	3.48498	3.49778	3.51058	3.52338	3.53618	3.54898	3.56178	3.57458	3.58738	3.60018	3.61298	3.62578	3.63858	3.65138	3.66418
3.9	3.58208	3.59488	3.60768	3.62048	3.63328	3.64608	3.65888	3.67168	3.68448	3.69728	3.71008	3.72288	3.73568	3.74848	3.76128	3.77408
4.0	3.59146	3.60426	3.61706	3.62986	3.64266	3.65546	3.66826	3.68106	3.69386	3.70666	3.71946	3.73226	3.74506	3.75786	3.77066	3.78346
4.1	2.70775	2.72055	2.73335	2.74615	2.75895	2.77175	2.78455	2.79735	2.81015	2.82295	2.83575	2.84855	2.86135	2.87415	2.88695	2.89975
4.2	2.66102	2.67382	2.68662	2.69942	2.71222	2.72502	2.73782	2.75062	2.76342	2.77622	2.78902	2.80182	2.81462	2.82742	2.84022	2.85302
4.3	3.00998	3.02278	3.03558	3.04838	3.06118	3.07398	3.08678	3.09958	3.11238	3.12518	3.13798	3.15078	3.16358	3.17638	3.18918	3.20198
4.4	3.15525	3.16805	3.18085	3.19365	3.20645	3.21925	3.23205	3.24485	3.25765	3.27045	3.28325	3.29605	3.30885	3.32165	3.33445	3.34725
4.5	3.29734	3.31014	3.32294	3.33574	3.34854	3.36134	3.37414	3.38694	3.39974	3.41254	3.42534	3.43814	3.45094	3.46374	3.47654	3.48934
4.6	3.43664	3.44944	3.46224	3.47504	3.48784	3.50064	3.51344	3.52624	3.53904	3.55184	3.56464	3.57744	3.59024	3.60304	3.61584	3.62864
4.7	3.57349	3.58629	3.59909	3.61189	3.62469	3.63749	3.65029	3.66309	3.67589	3.68869	3.70149	3.71429	3.72709	3.73989	3.75269	3.76549
4.8	3.70816	3.72096	3.73376	3.74656	3.75936	3.77216	3.78496	3.79776	3.81056	3.82336	3.83616	3.84896	3.86176	3.87456	3.88736	3.90016
4.9	3.94393	3.95673	3.96953	3.98233	3.99513	4.00793	4.02073	4.03353	4.04633	4.05913	4.07193	4.08473	4.09753	4.11033	4.12313	4.13593
5.0	3.97195	3.98475	3.99755	4.01035	4.02315	4.03595	4.04875	4.06155	4.07435	4.08715	4.10000	4.11280	4.12560	4.13840	4.15120	4.16400
5.1	4.10142	4.11422	4.12702	4.13982	4.15262	4.16542	4.17822	4.19102	4.20382	4.21662	4.22942	4.24222	4.25502	4.26782	4.28062	4.29342
5.2	4.22949	4.24229	4.25509	4.26789	4.28069	4.29349	4.30629	4.31909	4.33189	4.34469	4.35749	4.37029	4.38309	4.39589	4.40869	4.42149
5.3	4.35630	4.36910	4.38190	4.39470	4.40750	4.42030	4.43310	4.44590	4.45870	4.47150	4.48430	4.49710	4.50990	4.52270	4.53550	4.54830
5.4	4.48196	4.49476	4.50756	4.52036	4.53316	4.54596	4.55876	4.57156	4.58436	4.59716	4.60996	4.62276	4.63556	4.64836	4.66116	4.67396
5.5	4.60659	4.61939	4.63219	4.64499	4.65779	4.67059	4.68339	4.69619	4.70899	4.72179	4.73459	4.74739	4.76019	4.77299	4.78579	4.79859
5.6	4.73028	4.74308	4.75588	4.76868	4.78148	4.79428	4.80708	4.81988	4.83268	4.84548	4.85828	4.87108	4.88388	4.89668	4.90948	4.92228
5.7	4.85311	4.86591	4.87871	4.89151	4.90431	4.91711	4.92991	4.94271	4.95551	4.96831	4.98111	4.99391	5.00671	5.01951	5.03231	5.04511
5.8	4.97516	4.98796	4.99976	5.01256	5.02536	5.03816	5.05096	5.06376	5.07656	5.08936	5.10216	5.11496	5.12776	5.14056	5.15336	5.16616
5.9	5.09651	5.10931	5.12211	5.13491	5.14771	5.16051	5.17331	5.18611	5.19891	5.21171	5.22451	5.23731	5.25011	5.26291	5.27571	5.28851
6.0	5.21722	5.23002	5.24282	5.25562	5.26842	5.28122	5.29402	5.30682	5.31962	5.33242	5.34522	5.35802	5.37082	5.38362	5.39642	5.40922
6.1	5.33736	5.35016	5.36296	5.37576	5.38856	5.40136	5.41416	5.42696	5.43976	5.45256	5.46536	5.47816	5.49096	5.50376	5.51656	5.52936
6.2	5.45696	5.46976	5.48256	5.49536	5.50816	5.52096	5.53376	5.54656	5.55936	5.57216	5.58496	5.59776	5.61056	5.62336	5.63616	5.64896
6.3	5.57614	5.58894	5.60174	5.61454	5.62734	5.64014	5.65294	5.66574	5.67854	5.69134	5.70414	5.71694	5.72974	5.74254	5.75534	5.76814



$$B = \frac{2\pi W}{\lambda_0}$$

$$BC = \frac{2\pi w}{\lambda_0}$$

$$K = \frac{2\pi w}{\lambda_0}$$

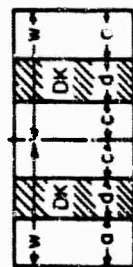
$$\alpha = \frac{\omega}{c}, \delta = \frac{d}{W}, \gamma = \frac{c}{W}$$

W GUIDE WIDTH  
DK RELATIVE DIELEC-  
TRIC CONSTANT  
 $\lambda_0$  FREE SPACE WAVE  
LENGTH  
 $\lambda_g$  GUIDE WAVELENGTH  
 $B = 2\pi W/\lambda_0$   
NORMALIZED FREQ.  
BC CUTOFF-FREQ.  
 $K = 2\pi w/\lambda_0$   
NORMALIZED  
PROPAGATION  
CONSTANT

$\alpha \cdot \delta/2$ MODE	NORMALIZED PROPAGATION CONSTANTS FOR TE MODES IN				WAVEGUIDES WHICH CONTAIN				DIELECTRIC SLABS				DK = 4.00				DELTA = 0.15			
	0.075 TE 10	0.125 TE 20	0.175 TE 30	0.225 TE 40	0.275 TE 50	0.325 TE 60	0.375 TE 70	0.425 TE 80	0.475 TE 90	0.525 TE 100	0.575 TE 110	0.625 TE 120	0.675 TE 130	0.725 TE 140	0.775 TE 150	0.825 TE 160	0.875 TE 170	0.925 TE 180	0.975 TE 190	1.025 TE 200
1.3	1.55733	3.02599	1.53926	2.88139	1.49781	2.64154	1.42780	2.41232	1.35644	2.29134	1.29261	2.25396	1.19478	1.12698	1.06742	0.92107	0.7742	0.62742	0.48061	0.33381
1.4	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.5	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.6	0.37042	-0.00000	0.44644	-0.00000	0.51937	-0.00000	0.59472	-0.00000	0.67152	-0.00000	0.75064	-0.00000	0.83204	-0.00000	0.91572	-0.00000	1.00176	-0.00000	1.09016	-0.00000
1.7	0.68800	-0.00000	0.73793	-0.00000	0.79320	-0.00000	0.85432	-0.00000	0.92124	-0.00000	0.99404	-0.00000	1.07272	-0.00000	1.15728	-0.00000	1.24872	-0.00000	1.34704	-0.00000
1.8	0.91102	-0.00000	0.95452	-0.00000	1.00320	-0.00000	1.05712	-0.00000	1.11632	-0.00000	1.18072	-0.00000	1.24644	-0.00000	1.31352	-0.00000	1.38296	-0.00000	1.45472	-0.00000
1.9	1.09886	-0.00000	1.13979	-0.00000	1.18712	-0.00000	1.24096	-0.00000	1.29924	-0.00000	1.36304	-0.00000	1.43232	-0.00000	1.50712	-0.00000	1.58752	-0.00000	1.67364	-0.00000
2.0	1.26875	-0.00000	1.30715	-0.00000	1.35312	-0.00000	1.40672	-0.00000	1.46704	-0.00000	1.53424	-0.00000	1.60752	-0.00000	1.68696	-0.00000	1.77264	-0.00000	1.86472	-0.00000
2.1	1.42236	-0.00000	1.46121	-0.00000	1.51167	-0.00000	1.56484	-0.00000	1.62992	-0.00000	1.69704	-0.00000	1.77632	-0.00000	1.86084	-0.00000	1.95176	-0.00000	2.04912	-0.00000
2.2	1.58893	-0.00000	1.63147	-0.00000	1.68712	-0.00000	1.74604	-0.00000	1.80832	-0.00000	1.87404	-0.00000	1.94332	-0.00000	2.01632	-0.00000	2.09312	-0.00000	2.17364	-0.00000
2.3	1.70906	-0.00000	1.75147	-0.00000	1.80712	-0.00000	1.86604	-0.00000	1.92832	-0.00000	2.00404	-0.00000	2.08332	-0.00000	2.16632	-0.00000	2.25312	-0.00000	2.34364	-0.00000
2.4	1.84414	-0.00000	1.88802	-0.00000	1.94432	-0.00000	2.00324	-0.00000	2.06572	-0.00000	2.13184	-0.00000	2.20164	-0.00000	2.27512	-0.00000	2.35232	-0.00000	2.43336	-0.00000
2.5	1.97821	-0.00000	2.02086	-0.00000	2.07612	-0.00000	2.13444	-0.00000	2.19584	-0.00000	2.26032	-0.00000	2.32792	-0.00000	2.39872	-0.00000	2.47284	-0.00000	2.55032	-0.00000
2.6	2.10305	-0.00000	2.15073	-0.00000	2.20204	-0.00000	2.25704	-0.00000	2.31572	-0.00000	2.37812	-0.00000	2.44424	-0.00000	2.51412	-0.00000	2.58784	-0.00000	2.66536	-0.00000
2.7	2.22821	-0.00000	2.27818	-0.00000	2.33144	-0.00000	2.38804	-0.00000	2.44804	-0.00000	2.51144	-0.00000	2.57832	-0.00000	2.64872	-0.00000	2.72272	-0.00000	2.80032	-0.00000
2.8	2.35113	-0.00000	2.40367	-0.00000	2.45992	-0.00000	2.51992	-0.00000	2.58364	-0.00000	2.65096	-0.00000	2.72192	-0.00000	2.79652	-0.00000	2.87484	-0.00000	2.95696	-0.00000
2.9	2.47216	-0.00000	2.52785	-0.00000	2.58765	-0.00000	2.65164	-0.00000	2.71984	-0.00000	2.79232	-0.00000	2.86912	-0.00000	2.94936	-0.00000	3.03312	-0.00000	3.12044	-0.00000
3.0	2.59188	-0.00000	2.65011	-0.00000	2.71216	-0.00000	2.77804	-0.00000	2.84784	-0.00000	2.92164	-0.00000	2.99944	-0.00000	3.08132	-0.00000	3.16736	-0.00000	3.25764	-0.00000
3.1	2.70960	-0.00000	2.77162	-0.00000	2.83764	-0.00000	2.90784	-0.00000	2.98224	-0.00000	3.06084	-0.00000	3.14372	-0.00000	3.23096	-0.00000	3.32264	-0.00000	3.41884	-0.00000
3.2	2.82643	-0.00000	2.89228	-0.00000	2.96204	-0.00000	3.03584	-0.00000	3.11372	-0.00000	3.19584	-0.00000	3.28224	-0.00000	3.37304	-0.00000	3.46752	-0.00000	3.56584	-0.00000
3.3	2.94220	-0.00000	3.01231	-0.00000	3.08656	-0.00000	3.16496	-0.00000	3.24752	-0.00000	3.33432	-0.00000	3.42544	-0.00000	3.52096	-0.00000	3.62096	-0.00000	3.72552	-0.00000
3.4	3.05707	-0.00000	3.13187	-0.00000	3.21116	-0.00000	3.29496	-0.00000	3.38336	-0.00000	3.47644	-0.00000	3.57432	-0.00000	3.67712	-0.00000	3.78484	-0.00000	3.89752	-0.00000
3.5	3.17114	-0.00000	3.25034	-0.00000	3.33416	-0.00000	3.42264	-0.00000	3.51584	-0.00000	3.61392	-0.00000	3.71696	-0.00000	3.82504	-0.00000	3.93824	-0.00000	4.05656	-0.00000
3.6	3.28452	-0.00000	3.36832	-0.00000	3.45764	-0.00000	3.55164	-0.00000	3.65044	-0.00000	3.75416	-0.00000	3.86284	-0.00000	3.97656	-0.00000	4.09536	-0.00000	4.21924	-0.00000
3.7	3.39729	-0.00000	3.48632	-0.00000	3.58064	-0.00000	3.67944	-0.00000	3.78284	-0.00000	3.89096	-0.00000	4.00384	-0.00000	4.12164	-0.00000	4.24456	-0.00000	4.37264	-0.00000
3.8	3.50984	-0.00000	3.59944	-0.00000	3.69872	-0.00000	3.80784	-0.00000	3.92684	-0.00000	4.04584	-0.00000	4.17484	-0.00000	4.30384	-0.00000	4.43792	-0.00000	4.57712	-0.00000
3.9	3.62133	-0.00000	3.71294	-0.00000	3.81432	-0.00000	3.92564	-0.00000	4.03696	-0.00000	4.14832	-0.00000	4.26976	-0.00000	4.39124	-0.00000	4.51284	-0.00000	4.63456	-0.00000
4.0	3.73273	-0.00000	3.82693	-0.00000	3.93112	-0.00000	4.04536	-0.00000	4.15964	-0.00000	4.27396	-0.00000	4.38832	-0.00000	4.50272	-0.00000	4.61716	-0.00000	4.73168	-0.00000
4.1	3.84413	-0.00000	3.94033	-0.00000	4.04656	-0.00000	4.15284	-0.00000	4.25916	-0.00000	4.36552	-0.00000	4.47192	-0.00000	4.57836	-0.00000	4.68484	-0.00000	4.79136	-0.00000
4.2	3.95553	-0.00000	4.05373	-0.00000	4.16004	-0.00000	4.26644	-0.00000	4.37284	-0.00000	4.47924	-0.00000	4.58564	-0.00000	4.69204	-0.00000	4.79844	-0.00000	4.90484	-0.00000
4.3	4.06693	-0.00000	4.16713	-0.00000	4.27352	-0.00000	4.37992	-0.00000	4.48632	-0.00000	4.59272	-0.00000	4.69912	-0.00000	4.80552	-0.00000	4.91192	-0.00000	5.01832	-0.00000
4.4	4.17833	-0.00000	4.27853	-0.00000	4.38492	-0.00000	4.49132	-0.00000	4.59772	-0.00000	4.70412	-0.00000	4.81052	-0.00000	4.91692	-0.00000	5.02332	-0.00000	5.12972	-0.00000
4.5	4.28973	-0.00000	4.38993	-0.00000	4.49632	-0.00000	4.60272	-0.00000	4.70912	-0.00000	4.81552	-0.00000	4.92192	-0.00000	5.02832	-0.00000	5.13472	-0.00000	5.24112	-0.00000
4.6	4.40113	-0.00000	4.50133	-0.00000	4.60772	-0.00000	4.71412	-0.00000	4.82052	-0.00000	4.92692	-0.00000	5.03332	-0.00000	5.13972	-0.00000	5.24612	-0.00000	5.35252	-0.00000
4.7	4.51253	-0.00000	4.61273	-0.00000	4.71912	-0.00000	4.82552	-0.00000	4.93192	-0.00000	5.03832	-0.00000	5.14472	-0.00000	5.25112	-0.00000	5.35752	-0.00000	5.46392	-0.00000
4.8	4.62393	-0.00000	4.72413	-0.00000	4.83052	-0.00000	4.93692	-0.00000	5.04332	-0.00000	5.14972	-0.00000	5.25612	-0.00000	5.36252	-0.00000	5.46892	-0.00000	5.57532	-0.00000
4.9	4.73533	-0.00000	4.83553	-0.00000	4.94192	-0.00000	5.04832	-0.00000	5.15472	-0.00000	5.26112	-0.00000	5.36752	-0.00000	5.47392	-0.00000	5.58032	-0.00000	5.68672	-0.00000
5.0	4.84673	-0.00000	4.94693	-0.00000	5.05332	-0.00000	5.15972	-0.00000	5.26612	-0.00000	5.37252	-0.00000	5.47892	-0.00000	5.58532	-0.00000	5.69172	-0.00000	5.79812	-0.00000
5.1	4.95813	-0.00000	5.05833	-0.00000	5.16472	-0.00000	5.27112	-0.00000	5.37752	-0.00000	5.48392	-0.00000	5.59032	-0.00000	5.69672	-0.00000	5.80312	-0.00000	5.90952	-0.00000
5.2	5.06953	-0.00000	5.16973	-0.00000	5.27612	-0.00000	5.38252	-0.00000	5.48892	-0.00000	5.59532	-0.00000	5.70172	-0.00000	5.80812	-0.00000	5.91452	-0.00000	6.02092	-0.00000
5.3	5.18093	-0.00000	5.28113	-0.00000	5.38752	-0.00000	5.49392	-0.00000	5.60032	-0.00000	5.70672	-0.00000	5.81312	-0.00000	5.91952	-0.00000	6.02592	-0.00000	6.13232	-0.00000
5.4	5.29233	-0.00000	5.39253	-0.00000	5.49892	-0.00000	5.60532	-0.00000	5.71172	-0.00000	5.81812	-0.00000	5.92452	-0.00000	6.03092	-0.00000	6.13732	-0.00000	6.24372	-0.00000
5.5	5.40373	-0.00000	5.50393	-0.00000	5.61032	-0.00000	5.71672	-0.00000	5.82312	-0.00000	5.92952	-0.00000	6.03592	-0.00000	6.14232	-0.00000	6.24872	-0.00000	6.35512	-0.00000
5.6	5.51513	-0.00000	5.61533	-0.00000	5.72172	-0.00000	5.82812	-0.00000	5.93452	-0.00000	6.04092	-0.00000	6.14732	-0.00000	6.25372	-0.00000	6.36012	-0.00000	6.46652	-0.00000
5.7	5.62653	-0.00000	5.72673	-0.00000	5.83312	-0.00000	5.93952	-0.00000	6.04592	-0.00000	6.15232	-0.00000	6.25872	-0.00000	6.36512	-0.00000	6.47152	-0.00000	6.57792	-0.00000
5.8	5.73793	-0.00000	5.83813	-0.00000	5.94452	-0.00000	6.05092	-0.00000	6.15732	-0.00000	6.26372	-0.00000	6.37012	-0.00000	6.47652	-0.00000	6.58292	-0.00000	6.68932	-0.00000
5.9	5.84933	-0.00000	5.94953	-0.00000																

NORMALIZED PROPAGATION CONSTANTS FOR TE MODES IN WAVEGUIDES WHICH CONTAIN DIELECTRIC SLABS									
MODE	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20	TE 10
	0.125	0.125	0.125	0.125	0.125	0.125	0.125	0.125	0.125
BC	1.50871	2.68008	1.44356	2.59707	1.34258	2.14923	1.24918	2.02463	1.17137
1.1	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.2	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.3	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.4	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.5	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.6	0.55766	-0.00000	0.76206	-0.00000	1.34784	-0.00000	1.29788	-0.00000	1.49917
1.7	0.92077	-0.00000	0.99401	-0.00000	1.26092	-0.00000	1.80343	-0.00000	1.70026
1.8	1.02940	-0.00000	1.19382	-0.00000	1.46545	-0.00000	1.69772	-0.00000	1.89772
1.9	1.21215	-0.00000	1.37821	-0.00000	1.64148	-0.00000	1.88484	-0.00000	2.07938
2.0	1.37948	-0.00000	1.54581	-0.00000	1.81994	-0.00000	2.06737	-0.00000	2.26205
2.1	1.53656	-0.00000	1.70911	-0.00000	1.99439	-0.00000	2.24702	-0.00000	2.44200
2.2	1.68643	-0.00000	1.86785	-0.00000	2.16858	-0.00000	2.42802	-0.00000	2.62017
2.3	1.83100	-0.00000	2.02282	-0.00000	2.34273	-0.00000	2.60230	-0.00000	2.79725
2.4	1.97163	-0.00000	2.17622	-0.00000	2.51692	-0.00000	2.77953	-0.00000	2.97378
2.5	2.10928	-0.00000	2.32878	-0.00000	2.69100	-0.00000	2.95726	-0.00000	3.15018
2.6	2.24472	-0.00000	2.48143	-0.00000	2.86503	-0.00000	3.13591	-0.00000	3.32678
2.7	2.37856	0.40172	2.63498	-0.00000	3.03918	-0.00000	3.31876	-0.00000	3.50383
2.8	2.51133	0.99858	2.79019	-0.00000	3.21025	-0.00000	3.49706	-0.00000	3.68155
2.9	2.64349	1.37116	2.94781	-0.00000	3.38042	-0.00000	3.67994	-0.00000	3.86010
3.0	2.77547	1.67642	3.10883	-0.00000	3.55061	-0.00000	3.86480	-0.00000	4.03960
3.1	2.90769	1.94719	3.27303	-0.00000	3.72453	-0.00000	4.05078	-0.00000	4.22014
3.2	3.04059	2.19894	3.44190	-0.00000	3.89858	-0.00000	4.23828	-0.00000	4.40179
3.3	3.17460	2.43201	3.61566	-0.00000	4.07245	-0.00000	4.42828	-0.00000	4.58458
3.4	3.31020	2.65916	3.79468	-0.00000	4.24650	-0.00000	4.61936	-0.00000	4.76855
3.5	3.44791	2.87888	3.97918	-0.00000	4.42050	-0.00000	4.81192	-0.00000	4.95369
3.6	3.58830	3.09401	4.16909	-0.00000	4.59454	-0.00000	5.00583	-0.00000	5.14000
3.7	3.73198	3.30608	4.36431	-0.00000	4.77854	-0.00000	5.20096	-0.00000	5.32746
3.8	3.87960	3.51624	4.56447	-0.00000	4.95113	-0.00000	5.39720	-0.00000	5.51604
3.9	4.03184	3.72540	4.76907	-0.00000	5.12335	-0.00000	5.59444	-0.00000	5.70571
4.0	4.18937	3.93426	4.97787	-0.00000	5.29594	-0.00000	5.79257	-0.00000	5.89643
4.1	4.35316	4.14337	5.14533	-0.00000	5.46812	-0.00000	5.99202	-0.00000	6.09202
4.2	4.52316	4.35316	5.37032	-0.00000	5.64024	-0.00000	6.19615	-0.00000	6.29615
4.3	4.69896	4.56395	5.59519	-0.00000	5.81244	-0.00000	6.39990	-0.00000	6.49990
4.4	4.87896	4.77896	5.81996	-0.00000	6.00530	-0.00000	6.60337	-0.00000	6.70337
4.5	5.06333	4.98933	6.04465	-0.00000	6.19761	-0.00000	6.80663	-0.00000	6.90663
4.6	5.25243	5.20413	6.26926	-0.00000	6.39021	-0.00000	7.00975	-0.00000	7.10975
4.7	5.44643	5.42043	6.49377	-0.00000	6.58286	-0.00000	7.21279	-0.00000	7.31279
4.8	5.64581	5.63815	6.71816	-0.00000	6.77518	-0.00000	7.41580	-0.00000	7.51580
4.9	5.85026	5.85726	6.94242	-0.00000	6.96752	-0.00000	7.61883	-0.00000	7.71883
5.0	6.05976	6.07764	7.16652	-0.00000	7.16000	-0.00000	7.82190	-0.00000	7.92190
5.1	6.27223	6.29923	7.39045	-0.00000	7.35244	-0.00000	8.02504	-0.00000	8.12504
5.2	6.48819	6.52189	7.61419	-0.00000	7.54458	-0.00000	8.22829	-0.00000	8.32829
5.3	6.70682	6.74682	7.83772	-0.00000	7.73686	-0.00000	8.43164	-0.00000	8.53164
5.4	6.92900	6.97000	8.06104	-0.00000	7.92918	-0.00000	8.63513	-0.00000	8.73513
5.5	7.15522	7.19522	8.28412	-0.00000	8.12129	-0.00000	8.83876	-0.00000	8.93876
5.6	7.38609	7.42109	8.50697	-0.00000	8.31354	-0.00000	9.04254	-0.00000	9.14254
5.7	7.62151	7.64751	8.72951	-0.00000	8.50598	-0.00000	9.24648	-0.00000	9.34648
5.8	7.86140	7.87440	8.95194	-0.00000	8.69818	-0.00000	9.45059	-0.00000	9.55059
5.9	8.10176	8.10176	9.17408	-0.00000	8.89042	-0.00000	9.65485	-0.00000	9.75485
6.0	8.34293	8.32933	9.39596	-0.00000	9.08266	-0.00000	9.85928	-0.00000	9.95928
6.1	8.58526	8.55726	9.61760	-0.00000	9.27486	-0.00000	10.06388	-0.00000	10.16388
6.2	8.82843	8.78843	9.83900	-0.00000	9.46606	-0.00000	10.26863	-0.00000	10.36863
6.3	9.07381	9.01381	10.06017	-0.00000	9.65726	-0.00000	10.47355	-0.00000	10.57355

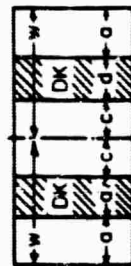
$2w$  GUIDE WIDTH  
 $DK$  RELATIVE DIELECTRIC CONSTANT  
 $\lambda_0$  FREE SPACE WAVELENGTH  
 $\lambda_g$  GUIDE WAVELENGTH  
 $B = 2\pi w/\lambda_0$   
 NORMALIZED FREQ.  
 $BC$  CUTOFF-FREQ.  
 $K = 2\pi w/\lambda_g$   
 NORMALIZED PROPAGATION CONSTANT



$$a = \frac{w}{2}, \delta = \frac{d}{w}, \gamma = \frac{c}{w}$$



MODE	NORMALIZED PROPAGATION CONSTANTS FOR TE MODES IN WAVEGUIDES WHICH CONTAIN DIELECTRIC SLABS										DK =		DELTA = 0.05	
	0.025		0.125		0.200		0.300		0.400		0.500		0.700	
	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20
BC	1.56948	3.1056	1.54495	2.91776	1.50710	2.67576	1.44244	2.45053	1.37556	2.33167	1.31502	2.29492	1.22134	1.14746
B	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	0.57659	0.85922
1	-0.00070	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	0.30086	-0.00000	0.58014	-0.00000	0.86700	1.13293
1.1	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	0.69233	-0.00000	0.87302	-0.00000	1.13003	1.37110
1.5	-0.31127	-0.00000	-0.42375	-0.00000	-0.56311	-0.00000	-0.76313	-0.00000	0.94772	-0.00000	1.10470	-0.00000	1.53093	1.59093
1.6	0.65300	-0.00000	0.72255	-0.00000	0.82507	-0.00000	0.99335	-0.00000	1.16096	-0.00000	1.30639	-0.00000	1.53850	1.79554
1.7	0.68207	-0.00000	0.94113	-0.00000	1.03335	-0.00000	1.19119	-0.00000	1.35247	-0.00000	1.49584	-0.00000	1.72277	2.00153
1.8	1.07177	-0.00000	1.12720	-0.00000	1.21599	-0.00000	1.37101	-0.00000	1.53074	-0.00000	1.67276	-0.00000	1.89922	2.19945
2.0	1.24071	-0.00000	1.29459	-0.00000	1.38332	-0.00000	1.53944	-0.00000	1.70026	-0.00000	1.84247	-0.00000	2.07002	2.39513
2.1	1.39646	-0.00000	1.48062	-0.00000	1.54049	-0.00000	1.70017	-0.00000	1.86373	-0.00000	2.00704	-0.00000	2.23662	2.58945
2.2	1.54300	-0.00000	1.59793	-0.00000	1.69348	-0.00000	1.85551	-0.00000	2.02797	-0.00000	2.16791	-0.00000	2.40007	2.78459
2.2	1.68275	-0.00000	1.73900	-0.00000	1.83522	-0.00000	2.00705	-0.00000	2.17922	-0.00000	2.32610	-0.00000	2.56113	2.98014
2.3	1.81727	-0.00000	1.87529	-0.00000	1.97604	-0.00000	2.15596	-0.00000	2.33344	0.78763	2.48239	0.98269	2.72040	3.17709
2.4	1.94766	-0.00000	2.00780	-0.00000	2.11391	-0.00000	2.30312	0.66074	2.48638	1.28220	2.63739	1.39155	2.87834	3.37593
2.5	2.07468	-0.00000	2.13730	-0.00000	2.24960	-0.00000	2.44926	1.16410	2.63861	1.60125	2.79159	1.71343	3.03534	3.57706
2.6	2.19892	-0.00000	2.26435	-0.00000	2.38371	0.44354	2.59500	1.52425	2.79065	1.89988	2.94538	2.00462	3.19171	3.78080
2.7	2.32084	-0.00000	2.38939	-0.00000	2.51677	1.01756	2.74088	1.82839	2.94291	2.16933	3.09911	2.26585	3.34770	3.98740
2.8	2.44077	-0.00000	2.51280	-0.00000	2.64924	1.38550	2.88740	2.10145	3.09575	2.41965	3.25307	2.51018	3.50355	4.19706
2.9	2.55899	-0.00000	2.62487	0.79599	2.78155	1.68869	3.03501	2.35468	3.24951	2.65653	3.40750	2.74236	3.65947	4.40993
3.0	2.67574	-0.00000	2.75586	1.17887	2.91411	1.95813	3.18413	2.59429	3.40447	2.80353	3.56265	2.96548	3.81562	4.62613
3.1	2.79120	0.66855	2.87599	1.47613	3.04736	2.20686	3.33518	2.82414	3.56089	3.10309	3.71871	3.18167	3.97216	4.84574
3.2	2.90553	1.												



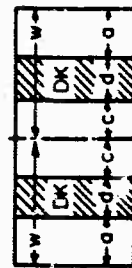


α-δ/2 MODE	NORMALIZED PROPAGATION CONSTANTS FOR TE MODES IN				WAVEGUIDES WHICH CONTAIN				DIELECTRIC SLABS				DK =				DELTA = 0.10			
	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20
BC	1.55988	3.03913	1.51483	2.65034	1.43926	2.30767	1.32674	2.05081	1.22742	1.92908	1.14707	1.89276	0.58444	0.97011	0.700	0.980	0.700	0.980	0.700	0.980
B	1.55988	3.03913	1.51483	2.65034	1.43926	2.30767	1.32674	2.05081	1.22742	1.92908	1.14707	1.89276	0.58444	0.97011	0.700	0.980	0.700	0.980	0.700	0.980
1.1	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.2	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.3	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.4	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.5	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.6	0.35872	-0.00000	0.53785	-0.00000	0.78060	-0.00000	1.10640	-0.00000	1.37337	-0.00000	1.59166	-0.00000	1.90633	-0.00000	2.23320	-0.00000	2.61318	-0.00000	3.04566	-0.00000
1.7	0.68098	-0.00000	0.80656	-0.00000	1.01426	-0.00000	1.32414	-0.00000	1.59190	-0.00000	1.80031	-0.00000	2.11318	-0.00000	2.58466	-0.00000	3.11318	-0.00000	3.73379	-0.00000
1.8	0.90510	-0.00000	1.01762	-0.00000	1.21706	-0.00000	1.52714	-0.00000	1.79558	-0.00000	2.00224	-0.00000	2.31498	-0.00000	2.73379	-0.00000	3.25340	-0.00000	3.86424	-0.00000
1.9	1.09322	-0.00000	1.20200	-0.00000	1.40330	-0.00000	1.72205	-0.00000	1.99429	-0.00000	2.20012	-0.00000	2.51340	-0.00000	2.93377	-0.00000	3.46242	-0.00000	4.07377	-0.00000
2.0	1.26155	-0.00000	1.37075	-0.00000	1.57977	-0.00000	1.91286	-0.00000	2.19057	-0.00000	2.39574	-0.00000	2.70963	-0.00000	3.13322	-0.00000	3.66261	-0.00000	4.25356	-0.00000
2.1	1.41718	-0.00000	1.52928	-0.00000	1.75080	-0.00000	2.10230	-0.00000	2.38620	-0.00000	2.59041	-0.00000	2.90456	-0.00000	3.33736	-0.00000	3.87758	-0.00000	4.48048	-0.00000
2.2	1.66391	-0.00000	1.68076	-0.00000	1.91821	-0.00000	2.29245	-0.00000	2.56252	-0.00000	2.78511	-0.00000	3.09891	-0.00000	3.53352	-0.00000	4.07898	-0.00000	4.73398	-0.00000
2.3	1.70405	-0.00000	1.82725	-0.00000	2.06506	-0.00000	2.48193	-0.00000	2.78052	-0.00000	2.98060	-0.00000	3.29323	-0.00000	3.73322	-0.00000	4.28440	-0.00000	4.95075	-0.00000
2.4	1.83912	-0.00000	1.97023	-0.00000	2.28288	-0.00000	2.68110	-0.00000	2.96100	-0.00000	3.17749	-0.00000	3.48798	-0.00000	3.93352	-0.00000	4.50076	-0.00000	5.18076	-0.00000
2.5	1.97018	-0.00000	2.11081	-0.00000	2.42339	-0.00000	2.88205	-0.00000	3.18454	-0.00000	3.37628	-0.00000	3.68356	-0.00000	4.13352	-0.00000	4.68031	-0.00000	5.36029	-0.00000
2.6	2.09799	-0.00000	2.24995	-0.00000	2.59831	-0.00000	3.08661	-0.00000	3.39186	-0.00000	3.57736	-0.00000	3.88031	-0.00000	4.33352	-0.00000	4.87798	-0.00000	5.56029	-0.00000
2.7	2.22312	-0.00000	2.38849	-0.00000	2.77939	-0.00000	3.30135	-0.00000	3.60233	-0.00000	3.78104	-0.00000	4.07898	-0.00000	4.53352	-0.00000	5.03352	-0.00000	5.71029	-0.00000
2.8	2.34802	-0.00000	2.52726	-0.00000	2.96636	-0.00000	3.52051	-0.00000	3.81701	-0.00000	3.98758	-0.00000	4.28440	-0.00000	4.73352	-0.00000	5.23352	-0.00000	5.91029	-0.00000
2.9	2.46704	-0.00000	2.66714	-0.00000	3.16888	-0.00000	3.74610	-0.00000	4.03563	-0.00000	4.19715	-0.00000	4.48031	-0.00000	4.93352	-0.00000	5.38031	-0.00000	6.06029	-0.00000
3.0	2.58845	-0.00000	2.80912	-0.00000	3.37633	-0.00000	3.97784	-0.00000	4.25815	-0.00000	4.40990	-0.00000	4.68440	-0.00000	5.13352	-0.00000	5.58031	-0.00000	6.31029	-0.00000
3.1	2.70448	-0.00000	2.98433	-0.00000	3.59760	-0.00000	4.21527	-0.00000	4.48443	-0.00000	4.62591	-0.00000	4.89081	-0.00000	5.33352	-0.00000	5.78031	-0.00000	6.51029	-0.00000
3.2	2.82134	-0.00000	3.10420	-0.00000	3.83092	-0.00000	4.45781	-0.00000	4.71432	-0.00000	4.84825	-0.00000	5.09972	-0.00000	5.53352	-0.00000	5.98031	-0.00000	6.76029	-0.00000
3.3	2.93718	-0.00000	3.28045	-0.00000	4.07879	-0.00000	4.70462	-0.00000	4.94762	-0.00000	5.06792	-0.00000	5.31129	-0.00000	5.75352	-0.00000	6.14076	-0.00000	6.93034	-0.00000
3.4	3.05213	-0.00000	3.42521	-0.00000	4.33110	-0.00000	4.98570	-0.00000	5.18414	-0.00000	5.29392	-0.00000	5.56576	-0.00000	5.99758	-0.00000	6.38644	-0.00000	7.16937	-0.00000
3.5	3.16634	-0.00000	3.60102	-0.00000	4.58536	-0.00000	5.20957	-0.00000	5.42366	-0.00000	5.52322	-0.00000	5.78576	-0.00000	6.21477	-0.00000	6.60644	-0.00000	7.38034	-0.00000
3.6	3.27989	-0.00000	3.79064	-0.00000	4.86695	-0.00000	5.46586	-0.00000	5.66608	-0.00000	5.75876	-0.00000	6.02141	-0.00000	6.43352	-0.00000	6.81477	-0.00000	7.59034	-0.00000
3.7	3.39290	-0.00000	3.99671	-0.00000	5.14435	-0.00000	5.72626	-0.00000	5.91166	-0.00000	6.00146	-0.00000	6.26390	-0.00000	6.66644	-0.00000	7.04825	-0.00000	7.80034	-0.00000
3.8	3.50845	-0.00000	4.22115	-0.00000	5.42623	-0.00000	5.96772	-0.00000	6.15862	-0.00000	6.23030	-0.00000	6.49243	-0.00000	6.95352	-0.00000	7.34076	-0.00000	8.09034	-0.00000
3.9	3.61784	-0.00000	4.46452	-0.00000	5.71155	-0.00000	6.25099	-0.00000	6.44054	-0.00000	6.50850	-0.00000	6.75213	-0.00000	7.21477	-0.00000	7.59758	-0.00000	8.33034	-0.00000
4.0	3.72954	-0.00000	4.72886	-0.00000	6.00000	-0.00000	6.51586	-0.00000	6.70000	-0.00000	6.75850	-0.00000	7.00000	-0.00000	7.46644	-0.00000	7.85253	-0.00000	8.59034	-0.00000
4.1	3.84691	-0.00000	4.96430	-0.00000	6.26556	-0.00000	6.77795	-0.00000	6.96000	-0.00000	7.00000	-0.00000	7.24644	-0.00000	7.65253	-0.00000	8.04076	-0.00000	8.78034	-0.00000
4.2	3.96931	-0.00000	5.28904	-0.00000	6.57913	-0.00000	7.04749	-0.00000	7.23000	-0.00000	7.26644	-0.00000	7.50000	-0.00000	7.89253	-0.00000	8.28076	-0.00000	9.02034	-0.00000
4.3	4.09290	-0.00000	5.67678	-0.00000	6.87350	-0.00000	7.31605	-0.00000	7.50000	-0.00000	7.53644	-0.00000	7.77000	-0.00000	8.16253	-0.00000	8.55076	-0.00000	9.25034	-0.00000
4.4	4.21733	-0.00000	6.06293	-0.00000	7.16669	-0.00000	7.58971	-0.00000	7.77000	-0.00000	7.80644	-0.00000	8.04000	-0.00000	8.42644	-0.00000	8.81676	-0.00000	9.54034	-0.00000
4.5	4.34242	-0.00000	6.45091	-0.00000	7.46471	-0.00000	7.86282	-0.00000	8.04000	-0.00000	8.07644	-0.00000	8.31000	-0.00000	8.69644	-0.00000	9.08276	-0.00000	9.83034	-0.00000
4.6	4.46751	-0.00000	6.83720	-0.00000	7.76157	-0.00000	8.13656	-0.00000	8.31000	-0.00000	8.34644	-0.00000	8.58000	-0.00000	8.96644	-0.00000	9.35276	-0.00000	10.09034	-0.00000
4.7	4.59260	-0.00000	7.22037	-0.00000	8.05932	-0.00000	8.41186	-0.00000	8.58000	-0.00000	8.61644	-0.00000	8.85000	-0.00000	9.23644	-0.00000	9.62276	-0.00000	10.38034	-0.00000
4.8	4.71769	-0.00000	7.63637	-0.00000	8.36602	-0.00000	8.68680	-0.00000	8.85000	-0.00000	8.88644	-0.00000	9.12000	-0.00000	9.50644	-0.00000	9.89276	-0.00000	10.67034	-0.00000
4.9	4.84278	-0.00000	8.01312	-0.00000	8.68773	-0.00000	8.99772	-0.00000	9.16000	-0.00000	9.19644	-0.00000	9.43000	-0.00000	9.81644	-0.00000	10.20276	-0.00000	10.95034	-0.00000
5.0	4.96787	-0.00000	8.39908	-0.00000	9.00000	-0.00000	9.28000	-0.00000	9.44000	-0.00000	9.47644	-0.00000	9.71000	-0.00000	10.09644	-0.00000	10.48276	-0.00000	11.23034	-0.00000
5.1	5.09296	-0.00000	8.78016	-0.00000	9.26045	-0.00000	9.55920	-0.00000	9.71000	-0.00000	9.74644	-0.00000	9.98000	-0.00000	10.36644	-0.00000	10.75276	-0.00000	11.51034	-0.00000
5.2	5.21805	-0.00000	9.16124	-0.00000	9.53539	-0.00000	9.83400	-0.00000	9.98000	-0.00000	10.01644	-0.00000	10.25000	-0.00000	10.63644	-0.00000	11.02276	-0.00000	11.76034	-0.00000
5.3	5.34314	-0.00000	9.54277	-0.00000	9.81000	-0.00000	10.10922	-0.00000	10.25000	-0.00000	10.28644	-0.00000	10.52000	-0.00000	10.90644	-0.00000	11.29276	-0.00000	12.03034	-0.00000
5.4	5.46823	-0.00000	9.92466	-0.00000	10.17374	-0.00000	10.47231	-0.00000	10.61000	-0.00000	10.64644	-0.00000	10.88000	-0.00000	11.26644	-0.00000	11.65276	-0.00000	12.32034	-0.00000
5.5	5.59332	-0.00000	10.30659	-0.00000	10.46064	-0.00000	10.76077	-0.00000	10.89000	-0.00000	10.92644	-0.00000	11.16000	-0.00000	11.54644	-0.00000	11.932			

NORMALIZED PROPAGATION CONSTANTS FOR TE MODES IN WAVEGUIDES WHICH CONTAIN DIELECTRIC SLABS									
$\alpha \cdot \delta/2$	MODE	0.075		0.125		0.175		0.225	
		TE 10	TE 20	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20
BC	1.53240	2.77371	1.48006	2.42592	1.37206	2.06103	1.22574	1.80887	1.11671
8									
1.0	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.1	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.2	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.3	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.4	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.5	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.6	0.47362	-0.00000	0.65565	-0.00000	0.99098	-0.00000	1.42408	-0.00000	1.51731
1.7	0.75806	-0.00000	0.90464	-0.00000	1.21932	-0.00000	1.65563	-0.00000	1.75553
1.8	0.97327	-0.00000	1.11158	-0.00000	1.43000	-0.00000	1.86244	-0.00000	1.98763
1.9	1.15861	-0.00000	1.29740	-0.00000	1.63238	-0.00000	2.10912	-0.00000	2.21725
2.0	1.32678	-0.00000	1.47085	-0.00000	1.83213	-0.00000	2.38860	-0.00000	2.47928
2.1	1.48376	-0.00000	1.63670	-0.00000	2.03341	-0.00000	2.67373	-0.00000	2.74428
2.2	1.63291	0.00000	1.79806	-0.00000	2.23972	1.34486	2.81538	2.42118	3.15362
2.3	1.77637	0.00000	1.95723	-0.00000	2.45420	1.80851	3.06459	3.39709	3.64508
2.4	1.91559	0.00000	2.11619	-0.00000	2.67968	2.20923	3.32111	3.64508	3.92097
2.5	2.05151	0.00000	2.27685	0.89222	2.91844	2.57885	3.58595	3.92097	4.20977
2.6	2.18533	-0.00000	2.44123	1.40183	3.17192	2.93107	3.85708	4.15403	4.50604
2.7	2.31736	-0.00000	2.61168	1.80451	3.44036	3.27307	4.13398	4.41457	4.80683
2.8	2.44832	0.46366	2.79091	2.16474	3.72278	3.60898	4.41566	4.67871	5.09485
2.9	2.57880	1.03286	2.98216	2.50457	4.01724	3.94125	4.70115	4.94611	5.38381
3.0	2.70935	1.40613	3.18900	2.83477	4.32133	4.27136	4.98963	5.21645	5.67191
3.1	2.84060	1.71841	3.41488	3.16147	4.63265	4.60020	5.28040	5.48943	5.96116
3.2	2.97326	2.00059	3.66239	3.48837	4.94910	4.92827	5.57294	5.76481	6.25152
3.3	3.10821	2.26618	3.93231	3.81766	5.26902	5.25883	5.86684	6.04237	6.53801
3.4	3.24653	2.52282	4.22321	4.15049	5.59112	5.58201	6.16179	6.32193	6.83381
3.5	3.38965	2.77560	4.53197	4.48729	5.91449	5.90988	6.45759	6.60333	7.13298
3.6	3.53946	3.02829	4.85477	4.82802	6.23448	6.23645	6.75406	6.88646	7.43819
3.7	3.69845	3.28368	5.18794	5.17232	6.56283	6.56275	7.06111	7.17119	7.74574
3.8	3.86978	3.54483	5.52842	5.51969	6.89636	6.88879	7.34864	7.45744	8.05419
3.9	4.05725	3.81312	5.87387	5.86960	7.20970	7.21461	7.64660	7.74513	8.36344
4.0	4.26480	4.09027	6.22251	6.22156	7.53236	7.54026	7.94494	8.03419	8.67228
4.1		4.37727		6.57517		7.86579		8.24888	8.91945
4.2		4.67455		6.93011		8.19126		8.55006	9.21292
4.3		4.98204		7.28617		8.51673		8.85193	9.50766
4.4		5.29920		7.64320		8.84228		9.15453	9.80363
4.5		5.62518		8.00111		9.16798		9.45793	10.10085
4.6		5.95897		8.35986		9.49382		9.76215	10.39923
4.7		6.29954		8.71939		9.81991		10.06723	10.69875
4.8		6.64697		9.07968		10.14628		10.37320	10.99939
4.9		6.99746		9.44068		10.47294		10.68008	11.30110
5.0		7.35344		9.80231		10.79990		10.98789	11.60386
5.1		7.71347		10.16448		11.12718		11.29664	11.90764
5.2		8.07720		10.52708		11.45477		11.50633	12.21240
5.3		8.44472		10.88598		11.78264		11.81696	12.51812
5.4		8.81864		11.25304		12.11079		12.12929	12.82478
5.5		9.18994		11.61611		12.43919		12.54097	13.13236
5.6		9.56750		11.97903		12.76779		12.85431	13.44080
5.7		9.94815		12.34165		13.09658		13.16851	13.75012
5.8		10.33166		12.70382		13.42550		13.48354	14.06028
5.9		10.71774		13.06542		13.75485		13.79936	14.37126
6.0		11.10608		13.42633		14.08364		14.11593	14.68305
6.1		11.49621		13.78644		14.41278		14.43321	14.99561
6.2		11.88780		14.14565		14.74191		14.75115	
6.3		12.28039		14.50389		15.07101		15.06971	

$\Delta = 0.15$   
 $0.700$   $0.925$   
 $TE_{10}$   $TE_{10}$   
 $0.91190$   $0.82955$   
 $0.71684$   $1.10663$   
 $1.07761$   $1.44471$   
 $1.37075$   $1.75121$   
 $1.63385$   $2.04260$   
 $1.88028$   $2.32659$   
 $2.11657$   $2.60736$   
 $2.34648$   $2.88744$   
 $2.57240$   $3.16833$   
 $2.79595$   $3.45099$   
 $3.01835$   $3.73594$   
 $3.24047$   $4.02346$   
 $3.46302$   $4.31364$   
 $3.68658$   $4.60646$   
 $3.91159$   $4.90182$   
 $4.13845$   $5.19961$   
 $4.36748$   $5.49969$   
 $4.59895$   $5.80193$   
 $4.83308$   $6.10620$   
 $5.07006$   $6.41239$   
 $5.31003$   $6.72040$   
 $5.55311$   $7.03014$   
 $5.79939$   $7.34152$   
 $6.04892$   $7.65477$   
 $6.30173$   $7.96988$   
 $6.55783$   $8.28468$   
 $6.81720$   $8.60175$   
 $7.07980$   $8.92000$   
 $7.34558$   $9.23932$   
 $7.61448$   $9.55957$   
 $7.88643$   $9.88066$   
 $8.16133$   $10.20245$

$2W$  GUIDE WIDTH  
 $DK$  RELATIVE DIELEC.  
 $TRC$  CONSTANT  
 $\lambda_0$  FREE SPACE WAVE  
 $\lambda_g$  GUIDE WAVELENGTH  
 $B = 2\pi W/\lambda_0$   
 $N$  NORMALIZED FREQ.  
 $BC$  CUTOFF-FREQ.  
 $K = 2\pi W/\lambda_g$   
 $N$  NORMALIZED  
 $PROPAGATION$   
 $CONSTANT$



$$a = \frac{a}{W}, \delta = \frac{d}{W}, \gamma = \frac{c}{W}$$

NORMALIZED PROPAGATION CONSTANTS FOR TE MODES IN WAVEGUIDES WHICH CONTAIN DIELECTRIC SLABS									
MODE	$\alpha \cdot \delta/2$		TE 10		TE 20		TE 10		TE 20
	0.125	0.125	0.200	0.200	0.300	0.300	0.400	0.400	
TE 10	0.125	0.125	0.200	0.200	0.300	0.300	0.400	0.400	TE 10
TE 20	0.125	0.125	0.200	0.200	0.300	0.300	0.400	0.400	TE 20
TE 30	0.125	0.125	0.200	0.200	0.300	0.300	0.400	0.400	TE 30
TE 40	0.125	0.125	0.200	0.200	0.300	0.300	0.400	0.400	TE 40
TE 50	0.125	0.125	0.200	0.200	0.300	0.300	0.400	0.400	TE 50
TE 60	0.125	0.125	0.200	0.200	0.300	0.300	0.400	0.400	TE 60
TE 70	0.125	0.125	0.200	0.200	0.300	0.300	0.400	0.400	TE 70
TE 80	0.125	0.125	0.200	0.200	0.300	0.300	0.400	0.400	TE 80
TE 90	0.125	0.125	0.200	0.200	0.300	0.300	0.400	0.400	TE 90
TE 100	0.125	0.125	0.200	0.200	0.300	0.300	0.400	0.400	TE 100
TE 110	0.125	0.125	0.200	0.200	0.300	0.300	0.400	0.400	TE 110
TE 120	0.125	0.125	0.200	0.200	0.300	0.300	0.400	0.400	TE 120
TE 130	0.125	0.125	0.200	0.200	0.300	0.300	0.400	0.400	TE 130
TE 140	0.125	0.125	0.200	0.200	0.300	0.300	0.400	0.400	TE 140
TE 150	0.125	0.125	0.200	0.200	0.300	0.300	0.400	0.400	TE 150
TE 160	0.125	0.125	0.200	0.200	0.300	0.300	0.400	0.400	TE 160
TE 170	0.125	0.125	0.200	0.200	0.300	0.300	0.400	0.400	TE 170
TE 180	0.125	0.125	0.200	0.200	0.300	0.300	0.400	0.400	TE 180
TE 190	0.125	0.125	0.200	0.200	0.300	0.300	0.400	0.400	TE 190
TE 200	0.125	0.125	0.200	0.200	0.300	0.300	0.400	0.400	TE 200
TE 210	0.125	0.125	0.200	0.200	0.300	0.300	0.400	0.400	TE 210
TE 220	0.125	0.125	0.200	0.200	0.300	0.300	0.400	0.400	TE 220
TE 230	0.125	0.125	0.200	0.200	0.300	0.300	0.400	0.400	TE 230
TE 240	0.125	0.125	0.200	0.200	0.300	0.300	0.400	0.400	TE 240
TE 250	0.125	0.125	0.200	0.200	0.300	0.300	0.400	0.400	TE 250
TE 260	0.125	0.125	0.200	0.200	0.300	0.300	0.400	0.400	TE 260
TE 270	0.125	0.125	0.200	0.200	0.300	0.300	0.400	0.400	TE 270
TE 280	0.125	0.125	0.200	0.200	0.300	0.300	0.400	0.400	TE 280
TE 290	0.125	0.125	0.200	0.200	0.300	0.300	0.400	0.400	TE 290
TE 300	0.125	0.125	0.200	0.200	0.300	0.300	0.400	0.400	TE 300
TE 310	0.125	0.125	0.200	0.200	0.300	0.300	0.400	0.400	TE 310
TE 320	0.125	0.125	0.200	0.200	0.300	0.300	0.400	0.400	TE 320
TE 330	0.125	0.125	0.200	0.200	0.300	0.300	0.400	0.400	TE 330
TE 340	0.125	0.125	0.200	0.200	0.300	0.300	0.400	0.400	TE 340
TE 350	0.125	0.125	0.200	0.200	0.300	0.300	0.400	0.400	TE 350



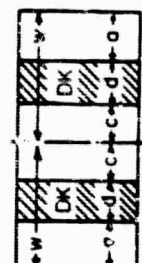


NORMALIZED PROPAGATION CONSTANTS FOR TE MODES IN WAVEGUIDES WHICH CONTAIN DIELECTRIC SLABS										
MODE	0.025		0.125		0.200		0.300		0.400	
	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20
BC	1.56894	3.12561	1.53380	2.80126	1.47967	2.80249	1.39287	2.25480	1.30941	2.13222
0										
1.1	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.2	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.3	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.4	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.5	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.6	0.31413	-0.00000	0.46922	-0.00000	0.65384	-0.00000	0.90949	-0.00000	1.13543	-0.00000
1.7	0.65535	-0.00000	0.75461	-0.00000	0.90065	-0.00000	1.12996	-0.00000	1.34433	-0.00000
1.8	0.88337	-0.00000	0.96990	-0.00000	1.10516	-0.00000	1.32719	-0.00000	1.53845	-0.00000
1.9	1.07297	-0.00000	1.15507	-0.00000	1.28803	-0.00000	1.51099	-0.00000	1.72332	-0.00000
2.0	1.24187	-0.00000	1.32292	-0.00000	1.48779	-0.00000	1.68647	-0.00000	1.90236	-0.00000
2.1	1.39760	-0.00000	1.47946	-0.00000	1.61898	-0.00000	1.85673	-0.00000	2.07776	-0.00000
2.2	1.54414	-0.00000	1.62805	-0.00000	1.77438	-0.00000	2.02391	-0.00000	2.25110	-0.00000
2.3	1.68391	-0.00000	1.77081	-0.00000	1.92587	-0.00000	2.18959	-0.00000	2.42355	-0.00000
2.4	1.81845	-0.00000	1.90917	-0.00000	2.07489	-0.00000	2.35505	-0.00000	2.59603	-0.00000
2.5	1.94886	-0.00000	2.04415	-0.00000	2.22257	-0.00000	2.52134	-0.00000	2.76931	-0.00000
2.6	2.07591	-0.00000	2.17555	-0.00000	2.36991	-0.00000	2.68941	-0.00000	2.94401	-0.00000
2.7	2.20019	-0.00000	2.30599	-0.00000	2.51785	-0.00000	2.86011	-0.00000	3.12067	-0.00000
2.8	2.32214	-0.00000	2.43899	-0.00000	2.66731	-0.00000	3.03419	-0.00000	3.29974	-0.00000
2.9	2.44211	-0.00000	2.56401	-0.00000	2.81925	-0.00000	3.21233	-0.00000	3.47675	-0.00000
3.0	2.56038	-0.00000	2.69148	-0.00000	2.97470	-0.00000	3.39511	-0.00000	3.65561	-0.00000
3.1	2.67718	-0.00000	2.81381	-0.00000	3.13474	-0.00000	3.58295	-0.00000	3.83659	-0.00000
3.2	2.79269	-0.00000	2.94645	-0.00000	3.30056	-0.00000	3.77617	-0.00000	4.01990	-0.00000
3.3	2.90706	-0.00000	3.07487	-0.00000	3.47336	-0.00000	3.97491	-0.00000	4.20579	-0.00000
3.4	3.02043	-0.00000	3.20463	-0.00000	3.65428	-0.00000	4.17913	-0.00000	4.39440	-0.00000
3.5	3.13291	-0.00000	3.33637	-0.00000	3.84432	-0.00000	4.38868	-0.00000	4.58586	-0.00000
3.6	3.24488	-0.00000	3.47089	-0.00000	4.04416	-0.00000	4.60328	-0.00000	4.78031	-0.00000
3.7	3.35583	-0.00000	3.60919	-0.00000	4.25407	-0.00000	4.82258	-0.00000	4.97784	-0.00000
3.8	3.46583	-0.00000	3.75254	-0.00000	4.47383	-0.00000	5.04619	-0.00000	5.17852	-0.00000
3.9	3.57583	-0.00000	3.90249	-0.00000	4.70283	-0.00000	5.27371	-0.00000	5.38240	-0.00000
4.0	3.68471	-0.00000	4.06097	-0.00000	4.94013	-0.00000	5.50476	-0.00000	5.58951	-0.00000
4.1	2.66986	-0.00000	3.97566	-0.00000	5.14885	-0.00000	5.71410	-0.00000	5.65457	-0.00000
4.2	2.82303	-0.00000	4.21801	-0.00000	5.41002	-0.00000	5.95734	-0.00000	5.88614	-0.00000
4.3	2.97177	-0.00000	4.46676	-0.00000	5.67364	-0.00000	6.20188	-0.00000	6.11899	-0.00000
4.4	3.11671	-0.00000	4.72240	-0.00000	5.93934	-0.00000	6.44782	-0.00000	6.35329	-0.00000
4.5	3.25836	-0.00000	4.98513	-0.00000	6.20706	-0.00000	6.69528	-0.00000	6.58920	-0.00000
4.6	3.39714	-0.00000	5.25492	-0.00000	6.47676	-0.00000	6.94434	-0.00000	6.82685	-0.00000
4.7	3.53339	-0.00000	5.53160	-0.00000	6.74837	-0.00000	7.19508	-0.00000	7.06636	-0.00000
4.8	3.66740	-0.00000	5.81447	-0.00000	7.02188	-0.00000	7.44759	-0.00000	7.30783	-0.00000
4.9	3.79942	-0.00000	6.10333	-0.00000	7.29728	-0.00000	7.70194	-0.00000	7.55135	-0.00000
5.0	3.92964	-0.00000	6.39754	-0.00000	7.57458	-0.00000	7.95821	-0.00000	7.79700	-0.00000
5.1	4.05824	-0.00000	6.69662	-0.00000	7.85383	-0.00000	8.21648	-0.00000	8.04484	-0.00000
5.2	4.18539	-0.00000	7.00012	-0.00000	8.13510	-0.00000	8.47684	-0.00000	8.29495	-0.00000
5.3	4.31122	-0.00000	7.30772	-0.00000	8.41848	-0.00000	8.73938	-0.00000	8.54738	-0.00000
5.4	4.43884	-0.00000	7.61919	-0.00000	8.70409	-0.00000	9.00418	-0.00000	8.80218	-0.00000
5.5	4.56936	-0.00000	7.93440	-0.00000	8.99204	-0.00000	9.27135	-0.00000	9.05940	-0.00000
5.6	4.68188	-0.00000	8.25334	-0.00000	9.28249	-0.00000	9.54099	-0.00000	9.31910	-0.00000
5.7	4.80349	-0.00000	8.57604	-0.00000	9.57887	-0.00000	9.81320	-0.00000	9.58656	-0.00000
5.8	4.92425	-0.00000	8.90263	-0.00000	9.87143	-0.00000	10.08809	-0.00000	9.84953	-0.00000
5.9	5.04424	-0.00000	9.23325	-0.00000	10.17023	-0.00000	10.36877	-0.00000	10.11807	-0.00000
6.0	5.16353	-0.00000	9.56806	-0.00000	10.47210	-0.00000	10.64638	-0.00000	10.38320	-0.00000
6.1	5.28216	-0.00000	9.90722	-0.00000	10.77718	-0.00000	10.92394	-0.00000	10.65396	-0.00000
6.2	5.40021	-0.00000	10.25086	-0.00000	11.08857	-0.00000	11.21664	-0.00000	10.92735	-0.00000
6.3	5.51771	-0.00000	10.59906	-0.00000	11.39736	-0.00000	11.50656	-0.00000	11.21028	-0.00000
6.4	5.63521	-0.00000	10.92735	-0.00000	11.70607	-0.00000	11.79547	-0.00000	11.49155	-0.00000
6.5	5.75271	-0.00000	11.25476	-0.00000	12.01478	-0.00000	12.08438	-0.00000	11.78046	-0.00000
6.6	5.87021	-0.00000	11.58307	-0.00000	12.32349	-0.00000	12.37309	-0.00000	12.06937	-0.00000
6.7	5.98771	-0.00000	11.91138	-0.00000	12.63220	-0.00000	12.66170	-0.00000	12.35828	-0.00000
6.8	6.10521	-0.00000	12.23969	-0.00000	12.94091	-0.00000	12.95031	-0.00000	12.64719	-0.00000
6.9	6.22271	-0.00000	12.56800	-0.00000	13.24962	-0.00000	13.24962	-0.00000	12.93610	-0.00000
7.0	6.34021	-0.00000	12.89631	-0.00000	13.55833	-0.00000	13.54851	-0.00000	13.22501	-0.00000
7.1	6.45771	-0.00000	13.22462	-0.00000	13.86704	-0.00000	13.83740	-0.00000	13.51390	-0.00000
7.2	6.57521	-0.00000	13.55293	-0.00000	14.17575	-0.00000	14.12628	-0.00000	13.80039	-0.00000
7.3	6.69271	-0.00000	13.88124	-0.00000	14.48446	-0.00000	14.41487	-0.00000	14.08688	-0.00000
7.4	6.81021	-0.00000	14.20955	-0.00000	14.79317	-0.00000	14.69338	-0.00000	14.37337	-0.00000
7.5	6.92771	-0.00000	14.53786	-0.00000	15.10188	-0.00000	14.97189	-0.00000	14.65986	-0.00000
7.6	7.04521	-0.00000	14.86617	-0.00000	15.41059	-0.00000	15.25040	-0.00000	14.94635	-0.00000
7.7	7.16271	-0.00000	15.19448	-0.00000	15.71930	-0.00000	15.52891	-0.00000	15.23284	-0.00000
7.8	7.28021	-0.00000	15.52279	-0.00000	16.02801	-0.00000	15.80742	-0.00000	15.51933	-0.00000
7.9	7.39771	-0.00000	15.85110	-0.00000	16.33672	-0.00000	16.08593	-0.00000	15.80084	-0.00000
8.0	7.51521	-0.00000	16.17941	-0.00000	16.64543	-0.00000	16.36444	-0.00000	16.08235	-0.00000
8.1	7.63271	-0.00000	16.50772	-0.00000	16.95414	-0.00000	16.64295	-0.00000	16.36386	-0.00000
8.2	7.75021	-0.00000	16.83603	-0.00000	17.26285	-0.00000	16.92146	-0.00000	16.64537	-0.00000
8.3	7.86771	-0.00000	17.16434	-0.00000	17.57156	-0.00000	17.20000	-0.00000	16.92688	-0.00000
8.4	7.98521	-0.00000	17.49265	-0.00000	17.88027	-0.00000	17.47851	-0.00000	17.20839	-0.00000
8.5	8.10271	-0.00000	17.82096	-0.00000	18.18898	-0.00000	17.75702	-0.00000	17.48990	-0.00000
8.6	8.22021	-0.00000	18.14927	-0.00000	18.49769	-0.00000	18.03553	-0.00000	17.77141	-0.00000
8.7	8.33771	-0.00000	18.47758	-0.00000	18.80640	-0.00000	18.31404	-0.00000	18.05292	-0.00000
8.8	8.45521	-0.00000	18.80589	-0.00000	19.11511	-0.00000	18.59255	-0.00000	18.33443	-0.00000
8.9	8.57271	-0.00000	19.13420	-0.00000	19.42382	-0.00000	18.87106	-0.00000	18.61594	-0.00000
9.0	8.69021	-0.00000	19.46251	-0.00000	19.73253	-0.00000	19.14957	-0.00000	18.89745	-0.00000
9.1	8.80771	-0.00000	19.79082	-0.00000	20.04124	-0.00000	19.42808	-0.00000	19.17896	-0.00000
9.2	8.92521	-0.00000	20.11913	-0.00000	20.34995	-0.00000	19.70659	-0.00000	19.46047	-0.00000
9.3	9.04271	-0.00000	20.44744	-0.00000	20.65866	-0.00000	19.98510	-0.00000	19.74198	-0.00000
9.4	9.16021	-0.00000	20.77575	-0.00000	20.96737	-0.00000	20.26361	-0.00000	20.02349	-0.00000
9.5	9.27771	-0.00000	21.10406	-0.00000	21.27608	-0.00000	20.54212	-0.00000	20.30500	-0.00000
9.6	9.39521	-0.00000	21.43237	-0.00000	21.58479	-0.00000	20.82063	-0.00000	20.58651	-0.00000
9.7	9.51271	-0.00000	21.76068	-0.00000	21.89350</					

6.2  
6.3

NORMALIZED PROPAGATION CONSTANTS FOR TE MODES IN WAVEGUIDES WHICH CONTAIN DIELECTRIC SLABS

MODE	CONSTANTS FOR TE MODES IN			WAVEGUIDES WHICH CONTAIN			DIELECTRIC SLABS			DK = 12.25			DELTA = 0.10		
	0.050	0.125	0.200	0.300	0.400	0.500	0.600	0.700	0.800	0.900	1.000	1.100	1.200	1.300	1.400
BC	1.55507	2.98319	1.48893	2.44914	1.38395	2.08003	1.24438	1.82928	1.13271	1.71444	1.04717	1.68055	0.92907	0.84027	0.75000
0	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
0.9	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.0	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.1	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.2	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.3	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.4	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.5	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.6	0.38075	-0.00000	0.62747	-0.00000	0.96704	-0.00000	1.38131	-0.00000	1.70778	-0.00000	1.94365	-0.00000	2.29183	-0.00000	2.61135
1.7	0.69468	-0.00000	0.88118	-0.00000	1.16718	-0.00000	1.51295	-0.00000	1.93965	-0.00000	2.17151	-0.00000	2.51751	-0.00000	2.81914
1.8	0.91700	-0.00000	1.08965	-0.00000	1.39819	-0.00000	1.73933	-0.00000	2.16916	-0.00000	2.39721	-0.00000	2.62273	-0.00000	2.84947
1.9	1.10482	-0.00000	1.27588	-0.00000	1.60626	-0.00000	2.06548	-0.00000	2.39912	-0.00000	2.62273	-0.00000	2.84947	-0.00000	3.07850
2.0	1.27274	-0.00000	1.44919	-0.00000	1.79942	-0.00000	2.29488	-0.00000	2.63149	-0.00000	2.84947	-0.00000	3.07850	-0.00000	3.31062
2.1	1.42853	-0.00000	1.61468	-0.00000	2.00013	-0.00000	2.53008	-0.00000	2.86764	-0.00000	3.07850	-0.00000	3.31062	-0.00000	3.54643
2.2	1.57559	-0.00000	1.77527	-0.00000	2.20621	-0.00000	2.72286	-0.00000	3.06117	-0.00000	3.29803	-0.00000	3.54643	-0.00000	3.78637
2.3	1.71620	-0.00000	1.93366	-0.00000	2.42121	-0.00000	3.02432	-0.00000	3.25158	-0.00000	3.47999	-0.00000	3.78637	-0.00000	4.03076
2.4	1.85187	-0.00000	2.09184	-0.00000	2.64847	-0.00000	3.28490	-0.00000	3.46619	-0.00000	3.68443	-0.00000	3.98702	-0.00000	4.27980
2.5	1.98364	-0.00000	2.26185	-0.00000	2.89087	-0.00000	3.55438	-0.00000	3.66594	-0.00000	3.88443	-0.00000	4.27980	-0.00000	4.53361
2.6	2.11228	-0.00000	2.41594	-0.00000	3.15034	-0.00000	3.83208	-0.00000	3.84799	-0.00000	4.03076	-0.00000	4.53361	-0.00000	4.79221
2.7	2.23336	-0.00000	2.58681	-0.00000	3.42746	-0.00000	4.11703	-0.00000	4.03233	-0.00000	4.27980	-0.00000	4.79221	-0.00000	5.05560
2.8	2.36234	-0.00000	2.76779	-0.00000	3.72123	-0.00000	4.40812	-0.00000	4.28626	-0.00000	4.53361	-0.00000	5.05560	-0.00000	5.32373
2.9	2.48459	-0.00000	2.96301	-0.00000	4.02940	-0.00000	4.70430	-0.00000	4.58352	-0.00000	4.79221	-0.00000	5.32373	-0.00000	5.59441
3.0	2.60540	-0.00000	3.17731	-0.00000	4.34911	-0.00000	5.00464	-0.00000	4.87980	-0.00000	5.05560	-0.00000	5.59441	-0.00000	5.86741
3.1	2.72504	-0.00000	3.41562	-0.00000	4.67753	-0.00000	5.30838	-0.00000	5.28779	-0.00000	5.59441	-0.00000	5.86741	-0.00000	6.14198
3.2	2.84372	-0.00000	3.68164	-0.00000	5.01221	-0.00000	5.61494	-0.00000	5.79719	-0.00000	6.14198	-0.00000	6.14198	-0.00000	6.41738
3.3	2.96166	-0.00000	3.97629	-0.00000	5.36122	-0.00000	5.92385	-0.00000	6.08855	-0.00000	6.41738	-0.00000	6.41738	-0.00000	6.69272
3.4	3.07904	-0.00000	4.29716	-0.00000	5.69316	-0.00000	6.23480	-0.00000	6.36861	-0.00000	6.69272	-0.00000	6.69272	-0.00000	6.96815
3.5	3.19605	-0.00000	4.63948	-0.00000	6.03698	-0.00000	6.54753	-0.00000	6.68171	-0.00000	6.96815	-0.00000	6.96815	-0.00000	7.24328
3.6	3.31290	-0.00000	4.99790	-0.00000	6.38193	-0.00000	6.86188	-0.00000	6.97447	-0.00000	7.24328	-0.00000	7.24328	-0.00000	7.51859
3.7	3.42978	-0.00000	5.36768	-0.00000	6.72747	-0.00000	7.17770	-0.00000	7.28586	-0.00000	7.51859	-0.00000	7.51859	-0.00000	7.79359
3.8	3.54694	-0.00000	5.74513	-0.00000	7.07318	-0.00000	7.49489	-0.00000	7.59190	-0.00000	7.79359	-0.00000	7.79359	-0.00000	8.06844
3.9	3.66466	-0.00000	6.12750	-0.00000	7.41872	-0.00000	7.81335	-0.00000	7.90044	-0.00000	8.06844	-0.00000	8.06844	-0.00000	8.34301
4.0	3.78330	-0.00000	6.51274	-0.00000	7.76385	-0.00000	8.14197	-0.00000	8.21140	-0.00000	8.34301	-0.00000	8.34301	-0.00000	8.61858
4.1	3.90215	-0.00000	6.90914	-0.00000	8.13426	-0.00000	8.46660	-0.00000	8.52694	-0.00000	8.61858	-0.00000	8.61858	-0.00000	8.89415
4.2	4.02110	-0.00000	7.30371	-0.00000	8.48790	-0.00000	8.79318	-0.00000	8.84445	-0.00000	8.94029	-0.00000	8.94029	-0.00000	9.21979
4.3	4.14015	-0.00000	7.70060	-0.00000	8.84296	-0.00000	9.12181	-0.00000	9.16437	-0.00000	9.21979	-0.00000	9.21979	-0.00000	9.49539
4.4	4.25930	-0.00000	8.09983	-0.00000	9.19959	-0.00000	9.45258	-0.00000	9.48672	-0.00000	9.53559	-0.00000	9.53559	-0.00000	9.81114
4.5	4.37855	-0.00000	8.50143	-0.00000	9.55789	-0.00000	9.78558	-0.00000	9.81150	-0.00000	9.86068	-0.00000	9.86068	-0.00000	10.13641
4.6	4.49780	-0.00000	8.90547	-0.00000	9.91796	-0.00000	10.12089	-0.00000	10.13873	-0.00000	10.18849	-0.00000	10.18849	-0.00000	10.46608
4.7	4.61705	-0.00000	9.31195	-0.00000	10.27989	-0.00000	10.45857	-0.00000	10.46839	-0.00000	10.51814	-0.00000	10.51814	-0.00000	10.79612
4.8	4.73630	-0.00000	9.72084	-0.00000	10.64371	-0.00000	10.79868	-0.00000	10.80350	-0.00000	10.85324	-0.00000	10.85324	-0.00000	11.12849
4.9	4.85555	-0.00000	10.13203	-0.00000	11.00945	-0.00000	11.14126	-0.00000	11.13504	-0.00000	11.18479	-0.00000	11.18479	-0.00000	11.46319
5.0	4.97480	-0.00000	10.54633	-0.00000	11.37708	-0.00000	11.48632	-0.00000	11.47202	-0.00000	11.52177	-0.00000	11.52177	-0.00000	11.80016
5.1	5.09405	-0.00000	10.96053	-0.00000	11.74659	-0.00000	11.83388	-0.00000	11.81142	-0.00000	11.86117	-0.00000	11.86117	-0.00000	12.13943
5.2	5.21330	-0.00000	11.37732	-0.00000	12.11790	-0.00000	12.18393	-0.00000	12.15323	-0.00000	12.20298	-0.00000	12.20298	-0.00000	12.48097
5.3	5.33255	-0.00000	11.79639	-0.00000	12.49094	-0.00000	12.53644	-0.00000	12.49744	-0.00000	12.54719	-0.00000	12.54719	-0.00000	12.82476
5.4	5.45180	-0.00000	12.21443	-0.00000	12.86562	-0.00000	12.89137	-0.00000	12.84403	-0.00000	12.89378	-0.00000	12.89378	-0.00000	13.17079
5.5	5.57105	-0.00000	12.63408	-0.00000	13.24162	-0.00000	13.24865	-0.00000	13.19297	-0.00000	13.24263	-0.00000	13.24263	-0.00000	13.51814
5.6	5.69030	-0.00000	13.05405	-0.00000	13.61945	-0.00000	13.60823	-0.00000	13.54423	-0.00000	13.59398	-0.00000	13.59398	-0.00000	13.86958
5.7	5.80955	-0.00000	13.47402	-0.00000	13.99837	-0.00000	13.97001	-0.00000	13.89778	-0.00000	13.94753	-0.00000	13.94753	-0.00000	14.22230
5.8	5.92880	-0.00000	13.89375	-0.00000	14.37848	-0.00000	14.33390	-0.00000	14.25356	-0.00000	14.30331	-0.00000	14.30331	-0.00000	14.57723
5.9	6.04805	-0.00000	14.31299	-0.00000	14.75964	-0.00000	14.69979	-0.00000	14.61155	-0.00000	14.66130	-0.00000	14.66130	-0.00000	14.93434
6.0	6.16730	-0.00000	14.73185	-0.00000	15.14175	-0.00000	15.06759	-0.00000	14.97167	-0.00000	15.02142	-0.00000	15.02142	-0.00000	15.29363
6.1	6.28655	-0.00000	15.14926	-0.00000	15.52469	-0.00000	15.43717	-0.00000	15.33388	-0.00000	15.38363	-0.00000	15.38363	-0.00000	15.65505
6.2	6.40580	-0.00000	15.56600	-0.00000	15.90835	-0.00000	15.80842	-0.00000	15.69811	-0.00000	15.74786	-0.00000	15.74786	-0.00000	16.01858
6.3	6.52505	-0.00000	15.98165	-0.00000	16.29264	-0.00000	16.18122	-0.00000	16.06429	-0.00000	16.11404	-0.00000	16.11404	-0.00000	16.38000



$\alpha = \frac{c}{W}, \delta = \frac{c}{W}, \gamma = \frac{c}{W}$   
 $\alpha = \frac{c}{W}, \delta = \frac{c}{W}, \gamma = \frac{c}{W}$

2W GUIDE WIDTH  
 DK RELATIVE DIELEC-  
 TRIC CONSTANT  
 $\lambda_0$  FREE SPACE WAVE  
 $\lambda_g$  GUIDE WAVELENGTH  
 $B = 2\pi w/\lambda_0$   
 NORMALIZED FREQ.  
 BC CUTOFF-FREQ.  
 $K = 2\pi w/\lambda_g$   
 NORMALIZED  
 PROPAGATION  
 CONSTANT

α-δ/λ	NORMA-LIZED PROPAGATION CONSTANTS FOR TE				MODES IN				WAVEGUIDES WHICH				DIELECTRIC SLABS				CONTAIN				DETA = 0.15				
	0.075	0.075	0.125	0.125	0.075	0.125	0.200	0.200	0.075	0.125	0.200	0.200	0.075	0.125	0.200	0.200	0.075	0.125	0.200	0.200	0.075	0.125	0.200	0.200	
BC	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20	
0.8	1.51450	2.59606	1.43813	2.19155	1.29469	1.62737	1.13032	1.54156	1.01146	1.48686	0.92476	1.45489	0.82222	1.27245	0.75140	1.20871	0.70000	1.15869	0.65000	1.10871	0.60000	1.05871	0.55000	1.00871	0.50000
0.9	1.51450	2.59606	1.43813	2.19155	1.29469	1.62737	1.13032	1.54156	1.01146	1.48686	0.92476	1.45489	0.82222	1.27245	0.75140	1.20871	0.70000	1.15869	0.65000	1.10871	0.60000	1.05871	0.55000	1.00871	0.50000
1.0	1.51450	2.59606	1.43813	2.19155	1.29469	1.62737	1.13032	1.54156	1.01146	1.48686	0.92476	1.45489	0.82222	1.27245	0.75140	1.20871	0.70000	1.15869	0.65000	1.10871	0.60000	1.05871	0.55000	1.00871	0.50000
1.1	1.51450	2.59606	1.43813	2.19155	1.29469	1.62737	1.13032	1.54156	1.01146	1.48686	0.92476	1.45489	0.82222	1.27245	0.75140	1.20871	0.70000	1.15869	0.65000	1.10871	0.60000	1.05871	0.55000	1.00871	0.50000
1.2	1.51450	2.59606	1.43813	2.19155	1.29469	1.62737	1.13032	1.54156	1.01146	1.48686	0.92476	1.45489	0.82222	1.27245	0.75140	1.20871	0.70000	1.15869	0.65000	1.10871	0.60000	1.05871	0.55000	1.00871	0.50000
1.3	1.51450	2.59606	1.43813	2.19155	1.29469	1.62737	1.13032	1.54156	1.01146	1.48686	0.92476	1.45489	0.82222	1.27245	0.75140	1.20871	0.70000	1.15869	0.65000	1.10871	0.60000	1.05871	0.55000	1.00871	0.50000
1.4	1.51450	2.59606	1.43813	2.19155	1.29469	1.62737	1.13032	1.54156	1.01146	1.48686	0.92476	1.45489	0.82222	1.27245	0.75140	1.20871	0.70000	1.15869	0.65000	1.10871	0.60000	1.05871	0.55000	1.00871	0.50000
1.5	1.51450	2.59606	1.43813	2.19155	1.29469	1.62737	1.13032	1.54156	1.01146	1.48686	0.92476	1.45489	0.82222	1.27245	0.75140	1.20871	0.70000	1.15869	0.65000	1.10871	0.60000	1.05871	0.55000	1.00871	0.50000
1.6	1.51450	2.59606	1.43813	2.19155	1.29469	1.62737	1.13032	1.54156	1.01146	1.48686	0.92476	1.45489	0.82222	1.27245	0.75140	1.20871	0.70000	1.15869	0.65000	1.10871	0.60000	1.05871	0.55000	1.00871	0.50000
1.7	1.51450	2.59606	1.43813	2.19155	1.29469	1.62737	1.13032	1.54156	1.01146	1.48686	0.92476	1.45489	0.82222	1.27245	0.75140	1.20871	0.70000	1.15869	0.65000	1.10871	0.60000	1.05871	0.55000	1.00871	0.50000
1.8	1.51450	2.59606</																							



[illegible]

NORMALIZED PROPAGATION CONSTANTS FOR TE MODES IN

WAVEGUIDES WHICH CONTAIN DIELECTRIC SLABS

RELATIVE DIELECTRIC CONSTANT OF SLAB

RELATIVE DIELECTRIC CONSTANT OF GUIDE

RELATIVE DIELECTRIC CONSTANT OF SURROUNDING MEDIUM

RELATIVE DIELECTRIC CONSTANT OF SURROUNDING MEDIUM

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RELATIVE DIELECTRIC CONSTANT OF SURROUNDING MEDIUM



$$V = \frac{\pi}{\lambda_0} \sqrt{\epsilon_r - \epsilon_0} d$$

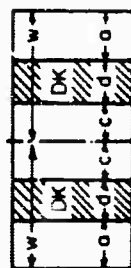
$$\delta = \frac{d}{W}, \gamma = \frac{c}{W}$$







NORMALIZED PROPAGATION CONSTANTS FOR TE MODES IN WAVEGUIDES WHICH CONTAIN DIELECTRIC SLABS	TE 10				TE 20				TE 30				TE 40				TE 50				TE 60				TE 70				TE 80				TE 90				TE 100				TE 110				TE 120				TE 130				TE 140				TE 150				TE 160				TE 170				TE 180				TE 190				TE 200				TE 210				TE 220				TE 230				TE 240				TE 250				TE 260				TE 270				TE 280				TE 290				TE 300				TE 310				TE 320				TE 330				TE 340				TE 350				TE 360				TE 370				TE 380				TE 390				TE 400				TE 410				TE 420				TE 430				TE 440				TE 450				TE 460				TE 470				TE 480				TE 490				TE 500				TE 510				TE 520				TE 530				TE 540				TE 550				TE 560				TE 570				TE 580				TE 590				TE 600				TE 610				TE 620				TE 630				TE 640				TE 650				TE 660				TE 670				TE 680				TE 690				TE 700				TE 710				TE 720				TE 730				TE 740				TE 750				TE 760				TE 770				TE 780				TE 790				TE 800				TE 810				TE 820				TE 830				TE 840				TE 850				TE 860				TE 870				TE 880				TE 890				TE 900				TE 910				TE 920																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																																		
0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075	0.075





[illegible]

WAVELENGTH PROPAGATION CONSTANTS FOR TE MODES IN WAVEGUIDES WHICH CONTAIN DIELECTRIC SLABS									
MODE	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20	TE 10
0.025	0.025	0.125	0.125	0.200	0.200	0.300	0.300	0.400	0.400
0.050	0.050	0.250	0.250	0.400	0.400	0.600	0.600	0.800	0.800
0.075	0.075	0.375	0.375	0.600	0.600	0.900	0.900	1.200	1.200
0.100	0.100	0.500	0.500	0.800	0.800	1.200	1.200	1.600	1.600
0.125	0.125	0.625	0.625	1.000	1.000	1.500	1.500	2.000	2.000
0.150	0.150	0.750	0.750	1.200	1.200	1.800	1.800	2.400	2.400
0.175	0.175	0.875	0.875	1.400	1.400	2.100	2.100	2.800	2.800
0.200	0.200	1.000	1.000	1.600	1.600	2.400	2.400	3.200	3.200
0.225	0.225	1.125	1.125	1.800	1.800	2.700	2.700	3.600	3.600
0.250	0.250	1.250	1.250	2.000	2.000	3.000	3.000	4.000	4.000
0.275	0.275	1.375	1.375	2.200	2.200	3.300	3.300	4.400	4.400
0.300	0.300	1.500	1.500	2.400	2.400	3.600	3.600	4.800	4.800
0.325	0.325	1.625	1.625	2.600	2.600	3.900	3.900	5.200	5.200
0.350	0.350	1.750	1.750	2.800	2.800	4.200	4.200	5.600	5.600
0.375	0.375	1.875	1.875	3.000	3.000	4.500	4.500	6.000	6.000
0.400	0.400	2.000	2.000	3.200	3.200	4.800	4.800	6.400	6.400
0.425	0.425	2.125	2.125	3.400	3.400	5.100	5.100	6.800	6.800
0.450	0.450	2.250	2.250	3.600	3.600	5.400	5.400	7.200	7.200
0.475	0.475	2.375	2.375	3.800	3.800	5.700	5.700	7.600	7.600
0.500	0.500	2.500	2.500	4.000	4.000	6.000	6.000	8.000	8.000
0.525	0.525	2.625	2.625	4.200	4.200	6.300	6.300	8.400	8.400
0.550	0.550	2.750	2.750	4.400	4.400	6.600	6.600	8.800	8.800
0.575	0.575	2.875	2.875	4.600	4.600	6.900	6.900	9.200	9.200
0.600	0.600	3.000	3.000	4.800	4.800	7.200	7.200	9.600	9.600
0.625	0.625	3.125	3.125	5.000	5.000	7.500	7.500	10.000	10.000
0.650	0.650	3.250	3.250	5.200	5.200	7.800	7.800	10.400	10.400
0.675	0.675	3.375	3.375	5.400	5.400	8.100	8.100	10.800	10.800
0.700	0.700	3.500	3.500	5.600	5.600	8.400	8.400	11.200	11.200
0.725	0.725	3.625	3.625	5.800	5.800	8.700	8.700	11.600	11.600
0.750	0.750	3.750	3.750	6.000	6.000	9.000	9.000	12.000	12.000
0.775	0.775	3.875	3.875	6.200	6.200	9.300	9.300	12.400	12.400
0.800	0.800	4.000	4.000	6.400	6.400	9.600	9.600	12.800	12.800
0.825	0.825	4.125	4.125	6.600	6.600	9.900	9.900	13.200	13.200
0.850	0.850	4.250	4.250	6.800	6.800	10.200	10.200	13.600	13.600
0.875	0.875	4.375	4.375	7.000	7.000	10.500	10.500	14.000	14.000
0.900	0.900	4.500	4.500	7.200	7.200	10.800	10.800	14.400	14.400
0.925	0.925	4.625	4.625	7.400	7.400	11.100	11.100	14.800	14.800
0.950	0.950	4.750	4.750	7.600	7.600	11.400	11.400	15.200	15.200
0.975	0.975	4.875	4.875	7.800	7.800	11.700	11.700	15.600	15.600
1.000	1.000	5.000	5.000	8.000	8.000	12.000	12.000	16.000	16.000
1.025	1.025	5.125	5.125	8.200	8.200	12.300	12.300	16.400	16.400
1.050	1.050	5.250	5.250	8.400	8.400	12.600	12.600	16.800	16.800
1.075	1.075	5.375	5.375	8.600	8.600	12.900	12.900	17.200	17.200
1.100	1.100	5.500	5.500	8.800	8.800	13.200	13.200	17.600	17.600
1.125	1.125	5.625	5.625	9.000	9.000	13.500	13.500	18.000	18.000
1.150	1.150	5.750	5.750	9.200	9.200	13.800	13.800	18.400	18.400
1.175	1.175	5.875	5.875	9.400	9.400	14.100	14.100	18.800	18.800
1.200	1.200	6.000	6.000	9.600	9.600	14.400	14.400	19.200	19.200
1.225	1.225	6.125	6.125	9.800	9.800	14.700	14.700	19.600	19.600
1.250	1.250	6.250	6.250	10.000	10.000	15.000	15.000	20.000	20.000
1.275	1.275	6.375	6.375	10.200	10.200	15.300	15.300	20.400	20.400
1.300	1.300	6.500	6.500	10.400	10.400	15.600	15.600	20.800	20.800
1.325	1.325	6.625	6.625	10.600	10.600	15.900	15.900	21.200	21.200
1.350	1.350	6.750	6.750	10.800	10.800	16.200	16.200	21.600	21.600
1.375	1.375	6.875	6.875	11.000	11.000	16.500	16.500	22.000	22.000
1.400	1.400	7.000	7.000	11.200	11.200	16.800	16.800	22.400	22.400
1.425	1.425	7.125	7.125	11.400	11.400	17.100	17.100	22.800	22.800
1.450	1.450	7.250	7.250	11.600	11.600	17.400	17.400	23.200	23.200
1.475	1.475	7.375	7.375	11.800	11.800	17.700	17.700	23.600	23.600
1.500	1.500	7.500	7.500	12.000	12.000	18.000	18.000	24.000	24.000
1.525	1.525	7.625	7.625	12.200	12.200	18.300	18.300	24.400	24.400
1.550	1.550	7.750	7.750	12.400	12.400	18.600	18.600	24.800	24.800
1.575	1.575	7.875	7.875	12.600	12.600	18.900	18.900	25.200	25.200
1.600	1.600	8.000	8.000	12.800	12.800	19.200	19.200	25.600	25.600
1.625	1.625	8.125	8.125	13.000	13.000	19.500	19.500	26.000	26.000
1.650	1.650	8.250	8.250	13.200	13.200	19.800	19.800	26.400	26.400
1.675	1.675	8.375	8.375	13.400	13.400	20.100	20.100	26.800	26.800
1.700	1.700	8.500	8.500	13.600	13.600	20.400	20.400	27.200	27.200
1.725	1.725	8.625	8.625	13.800	13.800	20.700	20.700	27.600	27.600
1.750	1.750	8.750	8.750	14.000	14.000	21.000	21.000	28.000	28.000
1.775	1.775	8.875	8.875	14.200	14.200	21.300	21.300	28.400	28.400
1.800	1.800	9.000	9.000	14.400	14.400	21.600	21.600	28.800	28.800
1.825	1.825	9.125	9.125	14.600	14.600	21.900	21.900	29.200	29.200
1.850	1.850	9.250	9.250	14.800	14.800	22.200	22.200	29.600	29.600
1.875	1.875	9.375	9.375	15.000	15.000	22.500	22.500	30.000	30.000
1.900	1.900	9.500	9.500	15.200	15.200	22.800	22.800	30.400	30.400
1.925	1.925	9.625	9.625	15.400	15.400	23.100	23.100	30.800	30.800
1.950	1.950	9.750	9.750	15.600	15.600	23.400	23.400	31.200	31.200
1.975	1.975	9.875	9.875	15.800	15.800	23.700	23.700	31.600	31.600
2.000	2.000	10.000	10.000	16.000	16.000	24.000	24.000	32.000	32.000
2.025	2.025	10.125	10.125	16.200	16.200	24.300	24.300	32.400	32.400
2.050	2.050	10.250	10.250	16.400	16.400	24.600	24.600	32.800	32.800
2.075	2.075	10.375	10.375	16.600	16.600	24.900	24.900	33.200	33.200
2.100	2.100	10.500	10.500	16.800	16.800	25.200	25.200	33.600	33.600
2.125	2.125	10.625	10.625	17.000	17.000	25.500	25.500	34.000	34.000
2.150	2.150	10.750	10.750	17.200	17.200	25.800	25.800	34.400	34.400
2.175	2.175	10.875	10.875	17.400	17.400	26.100	26.100	34.800	34.800
2.200	2.200	11.000	11.000	17.600	17.600	26.400	26.400	35.200	35.200
2.225	2.225	11.125	11.125	17.800	17.800	26.700	26.700	35.600	35.600
2.250	2.250	11.250	11.250	18.000	18.000	27.000	27.000	36.000	36.000
2.275	2.275	11.375	11.375	18.200	18.200	27.300	27.300	36.400	36.400
2.300	2.300	11.500	11.500	18.400	18.400	27.600	27.600	36.800	36.800
2.325	2.325	11.625	11.625	18.600	18.600	27.900	27.900	37.200	37.200
2.350	2.350	11.750	11.750	18.800	18.800	28.200	28.200	37.600	37.600
2.375	2.375	11.875	11.875	19.000	19.000	28.500	28.500	38.000	38.000
2.400	2.400	12.000	12.000	19.200	19.200	28.800	28.800	38.400	38.400
2.425	2.425	12.125	12.125	19.400	19.400	29.100	29.100	38.800	38.800
2.450	2.450	12.250	12.250	19.600	19.600	29.400	29.400	39.200	39.200
2.475	2.475	12.375	12.375	19.800	19.800	29.700	29.700	39.600	39.600
2.500	2.500	12.500	12.500	20.000	20.000	30.000	30.000	40.000	40.000
2.525	2.525	12.625	12.625	20.200	20.200	30.300	30.300	40.400	40.400
2.550	2.550	12.750	12.750	20.400	20.400	30.600	30.600	40.800	40.800
2.575	2.575	12.875	12.875	20.600	20.600	30.900	30.900	41.200	41.200
2.600	2.600	13.000	13.000	20.800	20.800	31.200	31.200	41.600	41.600
2.625	2.625	13.125	13.125	21.000	21.000	31.500	31.500	42.000	42.000
2.650	2.650	13.250	13.250	21.200	21.200	31.800	31.800	42.400	42.400
2.675	2.675	13.375	13.375	21.400	21.400	32.100	32.100	42.800	42.800
2.700	2.700	13.500	13.500	21.600	21.600	32.400	32.400	43.200	43.200
2.725	2.								

[illegible]



[illegible]

NORMALIZED PROPAGATION CONSTANTS FOR TE MODES IN WAVEGUIDES WHICH CONTAIN DIELECTRIC SLABS									
MODE	TE 10		TE 20		TE 30		TE 40		DELTA = 0.28
	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20	TE 10	TE 20	TE 10
BC	1.00000	1.35520	0.88973	1.10990	0.72919	0.95141	0.63131	0.88176	0.43077
0	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
0.6	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
0.7	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
0.8	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
0.9	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.0	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000	-0.00000
1.1	0.26786	-0.00000	1.53623	-0.00000	2.42137	2.04826	2.84779	3.08420	3.47944
1.2	0.90791	-0.00000	2.08878	1.88811	2.98141	2.76804	3.33980	3.54338	3.91129
1.3	1.35986	-0.00000	2.61202	2.39267	3.49568	3.39458	3.84099	3.71633	4.34844
1.4	1.81622	-0.00000	3.19759	3.09911	4.04813	3.99838	4.34928	4.23734	4.79236
1.5	2.31981	1.99735	3.80232	3.75893	4.60331	4.57586	4.86248	4.81844	5.24386
1.6	2.88112	2.74249	4.41349	4.39467	5.15779	5.14354	5.37888	5.38097	5.70732
1.7	3.48676	3.42990	5.02332	5.01647	5.70976	5.70254	5.89734	5.89047	6.17172
1.8	4.11550	4.09311	5.62813	5.62557	6.25867	6.25532	6.41718	6.40707	6.64583
1.9	4.78160	4.74341	6.22649	6.22723	6.80438	6.80345	6.93800	6.93207	7.12819
2.0	5.38663	5.38525	6.81802	6.82191	7.34707	7.34802	7.45957	7.45630	7.61726
2.1	6.01661	6.01991	7.40278	7.41067	7.88698	7.88985	7.98175	7.98031	8.11243
2.2	6.63972	6.64911	7.98105	7.99436	8.42439	8.42954	8.50442	8.50440	8.61307
2.3	7.25516	7.27354	8.55320	8.57366	8.95954	8.96760	9.02780	9.02876	9.11855
2.4	7.86267	7.89394	9.11964	9.14914	9.49268	9.50441	9.55089	9.55358	9.66198
2.5	8.46233	8.51086	9.68081	9.72122	10.02392	10.04027	10.07445	10.07877	10.20823
2.6	9.05450	9.12466	10.23717	10.29024	10.58352	10.58543	10.59819	10.60448	10.75839
2.7	9.63970	9.73849	10.78918	10.85846	11.11007	11.12193	11.13058	11.13011	11.28120
2.8	10.21857	10.34336	11.33729	11.42005	11.60830	11.64430	11.65713	11.65414	11.81422
2.9	10.79181	10.94816	11.88193	11.98116	12.13375	12.17821	12.16911	12.16407	12.32911
3.0	11.36012	11.54074	12.42352	12.53989	12.65809	12.71186	12.69241	12.71138	12.87438
3.1	11.92415	12.14795	12.98242	13.09632	13.18144	13.24826	13.21844	13.23593	13.40051
3.2	12.48451	12.74267	13.49898	13.65055	13.70390	13.77840	13.73816	13.76678	13.93287
3.3	13.04175	13.33380	14.03348	14.20263	14.22858	14.31127	14.26082	14.29478	14.46047
3.4	13.59634	13.92132	14.56622	14.75266	14.74659	14.84338	14.78253	14.82288	15.00084
3.5	14.14866	14.50524	15.09740	15.30069	15.26701	15.37609	15.30416	15.35109	15.53269
3.6	14.69907	15.08559	15.62724	15.84682	15.78692	15.90796	15.82841	15.87931	16.06555
3.7	15.24784	15.66246	16.15591	16.39111	16.30641	16.43943	16.34630	16.40749	16.59353
3.8	15.79519	16.23600	16.68353	16.93365	16.82552	16.97046	16.86634	16.93560	17.12829
3.9	16.34133	16.80629	17.21029	17.47451	17.34439	17.50101	17.38704	17.46357	17.65642
4.0	16.88632	17.37347	17.73623	18.01378	17.86287	18.03107	17.90693	17.99137	18.18439
4.1	17.43770	18.49910	18.55154	18.83661	18.66281	18.83661	18.68195	18.75827	18.94681
4.2	17.98785	19.05784	19.08785	19.37334	19.19861	19.37334	19.14628	19.22347	19.41212
4.3	18.53799	19.61404	19.62279	19.90861	19.61805	19.80324	19.57334	19.64947	19.83791
4.4	19.08785	20.16785	20.15644	20.44279	20.14893	20.33324	20.09008	20.16624	20.35477
4.5	19.63799	20.71938	20.68885	20.97469	20.67324	20.85749	20.60701	20.68294	20.86947
4.6	20.18785	21.26878	21.25023	21.54556	21.19998	21.38424	21.13294	21.20887	21.39537
4.7	20.73799	21.81615	21.78931	22.08074	21.72614	21.91040	21.65861	21.73454	21.92107
4.8	21.28785	22.36161	22.33454	22.62706	22.27675	22.46101	22.20385	22.27978	22.46631
4.9	21.83799	22.90826	22.87679	23.17017	22.81921	23.00347	22.74646	22.82234	23.00887
5.0	22.38785	23.45720	23.42469	23.68744	23.33604	23.52030	23.26324	23.33917	23.52567
5.1	22.93799	24.00615	23.97324	24.23744	23.88447	24.06970	23.81266	23.88854	24.07417
5.2	23.48785	24.55510	24.52169	24.75744	24.66805	24.85230	24.59524	24.67117	24.85767
5.3	24.03799	25.10405	25.06970	25.29869	25.20924	25.39349	25.13643	25.21236	25.40000
5.4	24.58785	25.65300	25.61865	25.84200	25.75255	25.93680	25.67974	25.75567	25.94221
5.5	25.13799	26.20195	26.16760	26.39100	26.30155	26.48580	26.22874	26.30467	26.49127
5.6	25.68785	26.75090	26.71655	26.94000	26.85055	27.03480	26.77349	26.84942	27.03600
5.7	26.23799	27.29985	27.26550	27.48900	27.39955	27.58380	27.32674	27.40267	27.58927
5.8	26.78785	27.84880	27.81445	28.03800	27.94855	28.13280	27.87574	27.95167	28.13827
5.9	27.33799	28.39775	28.36340	28.59200	28.50255	28.68680	28.42974	28.50567	28.69227
6.0	27.88785	28.94670	28.91235	29.14100	29.05155	29.23580	28.97469	29.05062	29.23721
6.1	28.43799	29.49565	29.46130	29.69000	29.59955	29.78380	29.53249	29.60842	29.79497
6.2	28.98785	30.04460	29.99925	30.24900	30.15855	30.34280	30.09149	30.16742	30.35357
6.3	29.53799	30.59355	30.54820	30.74800	30.65765	30.84190	30.59069	30.66662	30.85417

2W GUIDE WIDTH  
DK RELATIVE DIELEC-  
TRIC CONSTANT  
 $\lambda_0$  FREE SPACE WAVE  
 $\lambda_g$  GUIDE WAVELENGTH  
 $B = 2\pi w/\lambda_0$   
NORMALIZED FREQ.  
BC CUTOFF-FREQ.  
 $K = 2\pi w/\lambda_g$   
NORMALIZED  
PROPAGATION  
CONSTANT



$$\alpha = \frac{\pi}{w}, \delta = \frac{\pi}{w}, \gamma = \frac{\pi}{w}$$